



Studi sperimentali delle Risonanze Giganti nei nuclei atomici: studio della struttura nucleare e applicazioni

Outline

Giant resonances : their classifications and physical issues addressed through their investigations

Compression modes – Monopole and Dipole

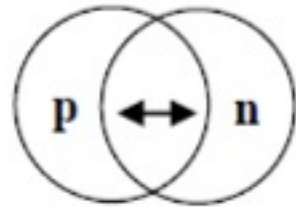
E1 response : why do we study it ?
what do we need to know?

E1 response in the low energy region :
Pygmy Dipole Resonance (in neutron rich nuclei)
stable nuclei and radioactive beams

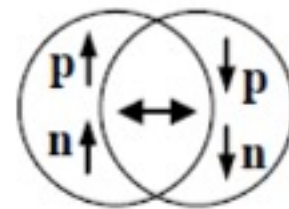
Giant resonances: Fundamental modes of nuclear excitation

$\Delta L = 1$

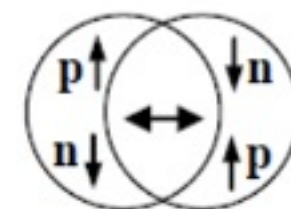
ISGDR
??



IVGDR

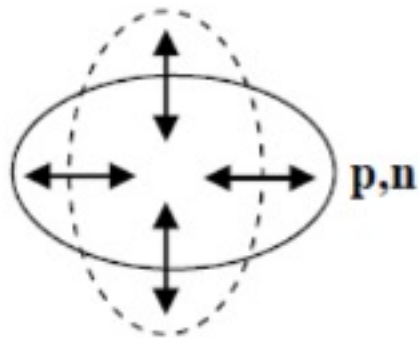


ISSGDR



IVSGDR

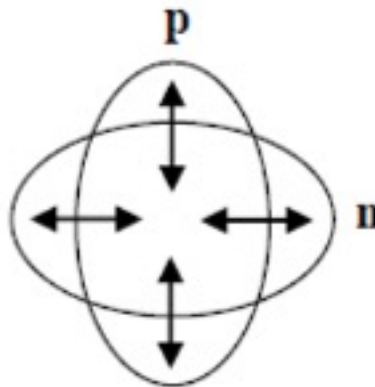
$\Delta L = 2$



ISGQR

$\Delta T = 0$

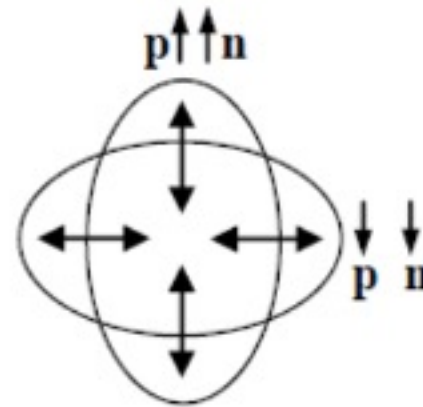
$\Delta S = 0$



IVGQR

$\Delta T = 1$

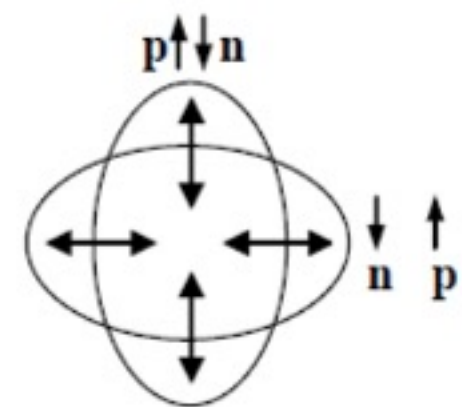
$\Delta S = 0$



ISSGQR

$\Delta T = 0$

$\Delta S = 1$



IVSGQR

$\Delta T = 1$

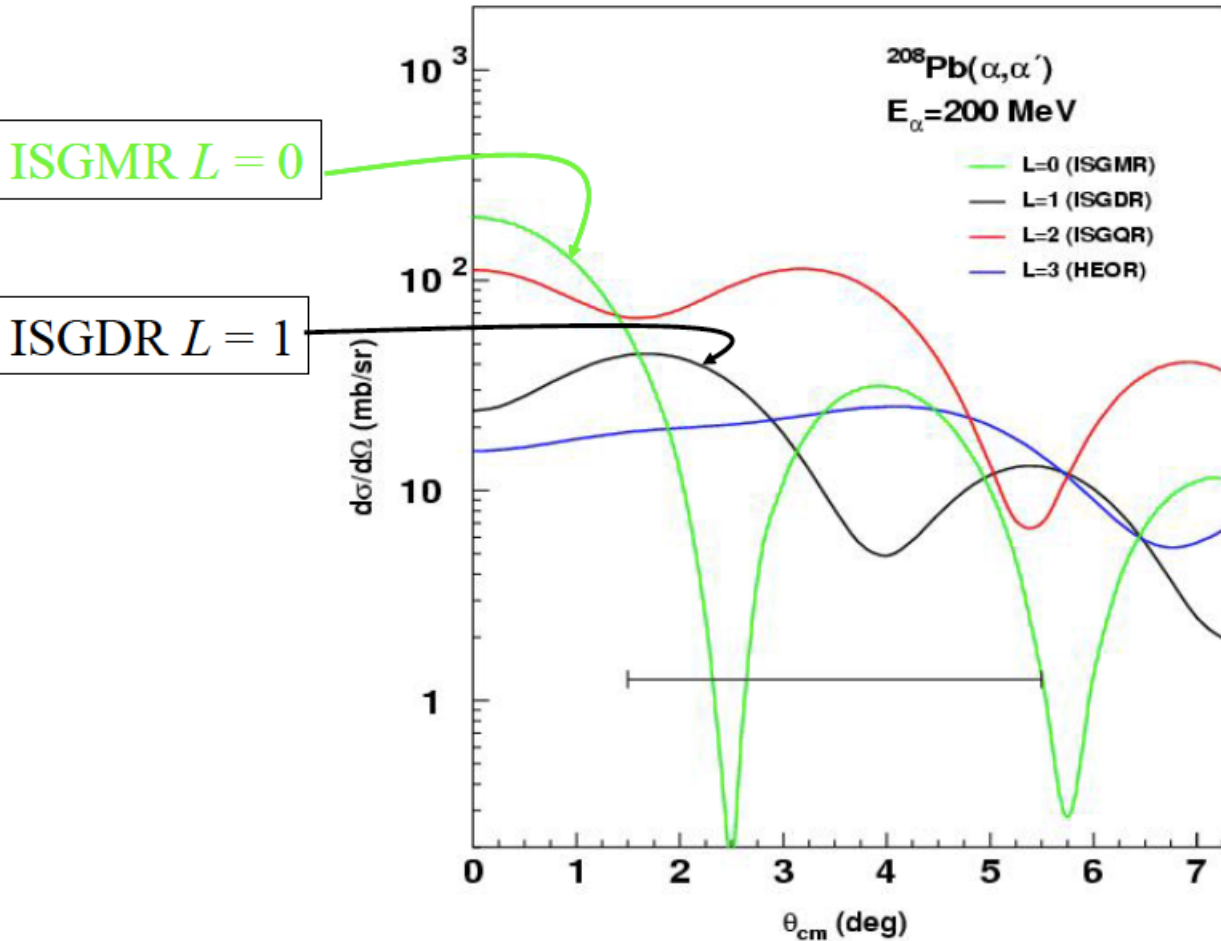
$\Delta S = 1$

Microscopic picture: GRs are coherent (1p-1h) excitations induced by single-particle operators.

- **Excitation energy depends on**
 - i) multipole L ($L\hbar\omega$, since radial operator $\propto r^L$; except for ISGMR and ISGDR, $2\hbar\omega$ & $3\hbar\omega$, respectively),*
 - ii) strength of effective interaction and*
 - iii) collectivity.*
- **Exhaust appreciable % of EWSR**
- **Acquire a width due to coupling to continuum and to underlying 2p-2h configurations.**

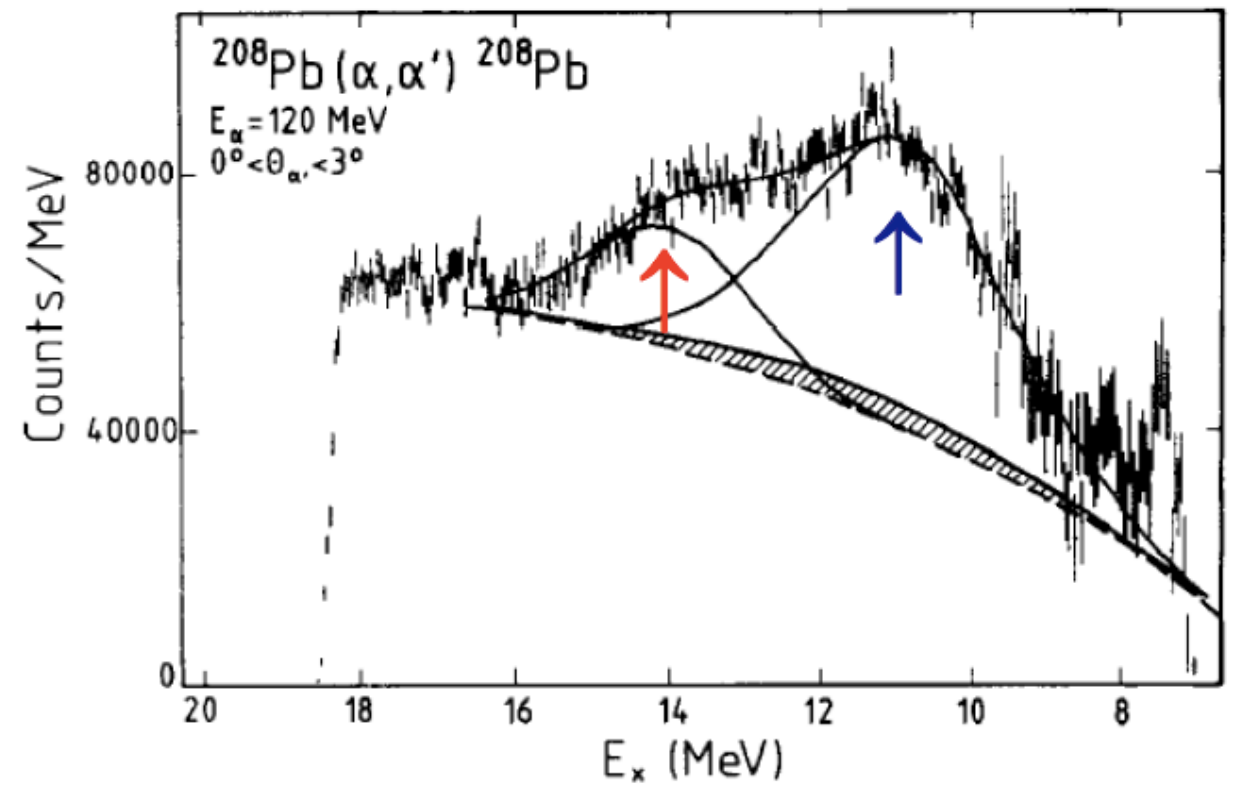
features

Angular distribution of
the population cross sections

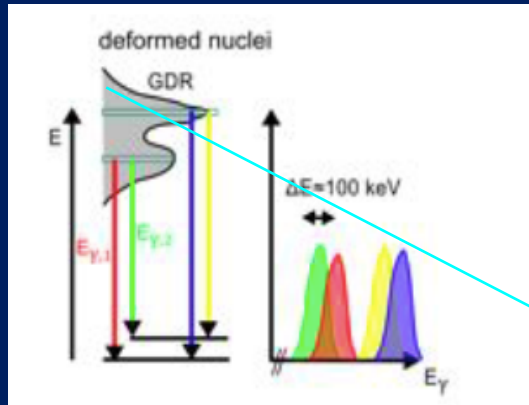


ISGQR at 10.9 MeV

ISGMR at 13.9 MeV



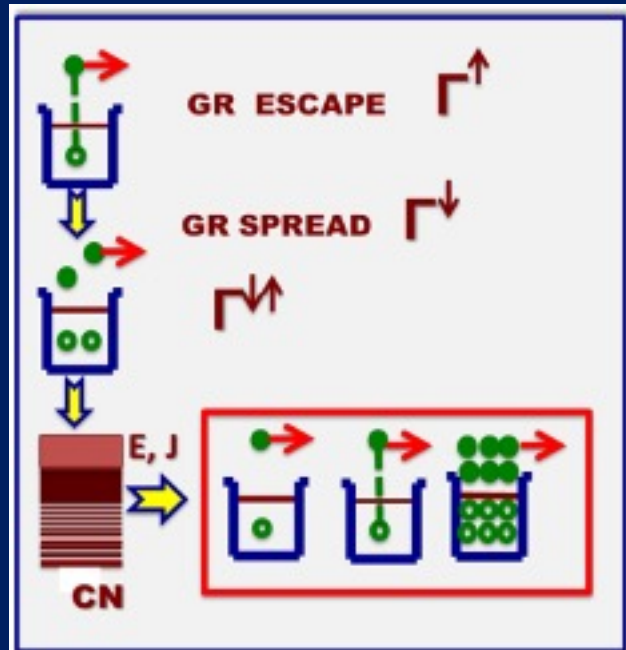
Damping of the giant dipole vibration in the nucleus



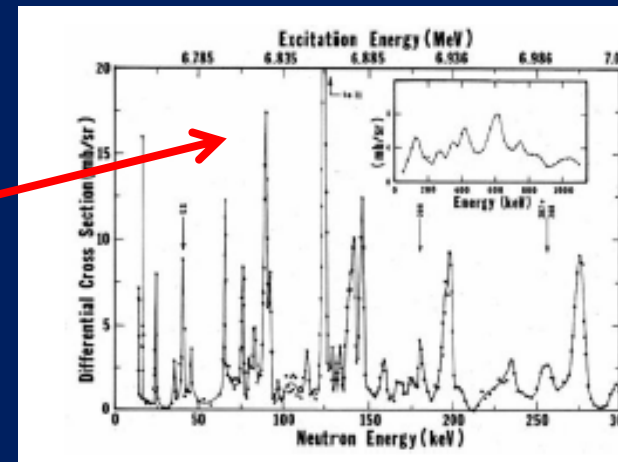
Exp. on gamma decay from GDR :
why?



Damping
mechanisms



- Ordered motion : dipole oscillation
- Coupling to complex states :
Stringent test to theory
- Regime at which the
nucleus is characterized by chaos :
Compound nuclear system



Compound
nuclear states
n-resonances

Microscopic structure of ISGMR & ISGDR

Transition operators:

$$O^{L=0} = \sum_i \cancel{r_i^0 Y_0^0} + \frac{1}{2} \sum_i r_i^2 Y_0^0 + \dots$$

Constant **Overtone**

$2\hbar\omega$ excitation

$$O^{L=1} = \sum_i \cancel{r_i^1 Y_0^1} + \frac{1}{2} \sum_i r_i^3 Y_0^1 + \dots$$

Spurious **Overtone**
c.o.m. motion

$3\hbar\omega$ excitation (overtone of c.o.m. motion)

Breathing mode

ISGMR ($T=0, L=0$)



ISGDR ($T=0, L=1$)



Squezing mode

ISGQR ($T=0, L=2$)



Nucleus $\longrightarrow$ **Many-body system with a finite size**

Vibrations $\longrightarrow$ **Multipole expansion with r , Y_{lm} , τ , σ**

$\Delta S=0, \Delta T=0$ $\Delta S=0, \Delta T=1$ $\Delta S=0, \Delta T=1$ $\Delta S=1, \Delta T=1$ $\Delta S=1, \Delta T=1$

L=0: Monopole	ISGMR $r^2 Y_0$	IAS τY_0	IVGMR $\tau r^2 Y_0$	GTR $\tau \sigma Y_0$	IVSGMR $\tau \sigma r^2 Y_0$
L=1: Dipole	ISGDR $r^3 Y_1$		IVGDR $\tau r Y_1$		IVSGDR $\tau \sigma r Y_1$
L=2: Quadrupole	ISGQR $r^2 Y_2$		IVGQR $\tau r^2 Y_2$		IVSGQR $\tau \sigma r^2 Y_2$
L=3: Octupole	LEOR, HEOR $r^3 Y_3$				

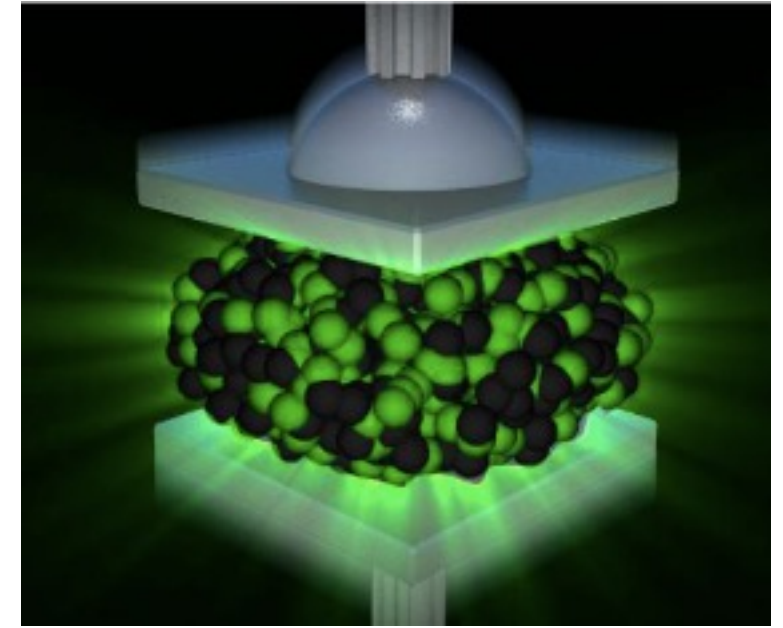
In fluid mechanics, **compressibility** is a measure of the relative volume change of a fluid as a response to a pressure change.

$$\beta = -\frac{1}{V} \frac{\partial V}{\partial P}$$

where P is pressure, V is volume.

Incompressibility or **bulk modulus** (K) is a measure of a substance's resistance to uniform compression and can be formally defined:

$$K = -V \frac{\partial P}{\partial V}$$



ISGMR ($T=0, L=0$)



ISGDR ($T=0, L=1$)



In Constrained and Scaling Models:

$$E_{ISGMR} = \hbar \sqrt{\frac{K_A}{m \langle r^2 \rangle}}$$

$$E_{ISGDR} = \hbar \sqrt{\frac{7}{3} \frac{K_A + \frac{27}{25} \varepsilon_F}{m \langle r^2 \rangle}}$$

ε_F is the Fermi energy and the nucleus incompressibility:

→ $K_A = \left[r^2 (d^2(E/A)/dr^2) \right]_{r=R_0}$

J.P. Blaizot, Phys. Rep. 64 (1980) 171

In HF+RPA calculations,

$$K_{nm} = \left[9 \rho^2 \frac{d^2(E/A)}{d\rho^2} \right]_{\rho=\rho_0}$$

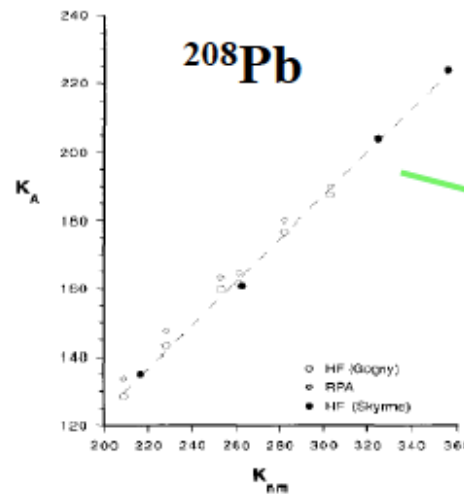
Nuclear matter

E/A : binding energy per nucleon

K_A : incompressibility

ρ : nuclear density

ρ_0 : nuclear density at saturation



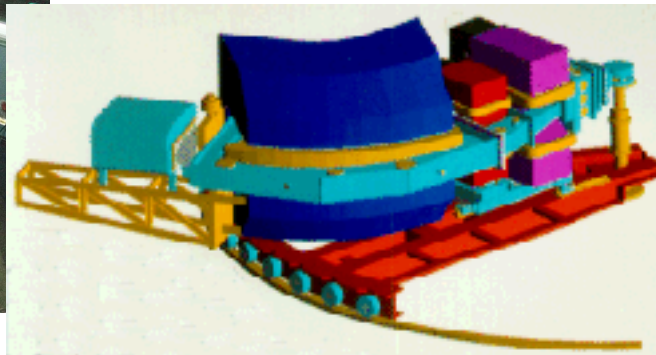
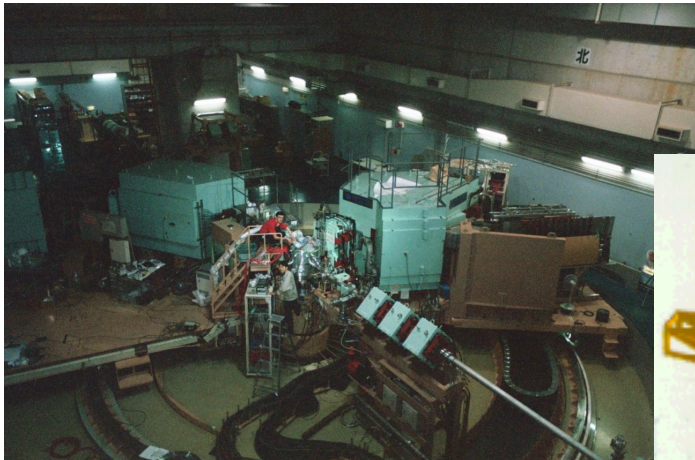
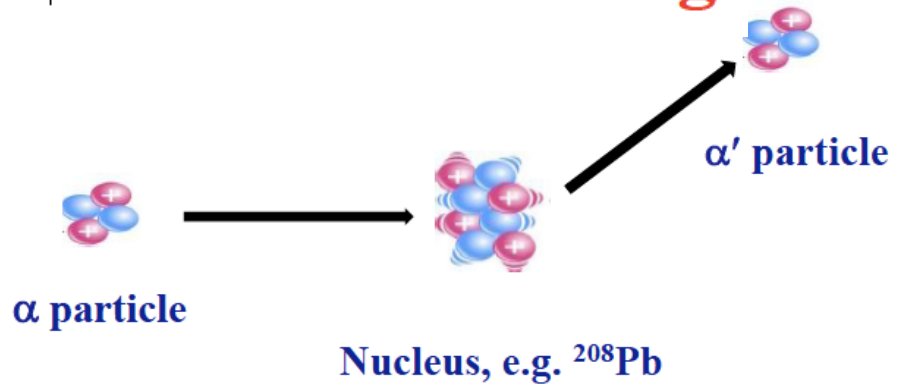
K_A is obtained from excitation energy of ISGMR & ISGDR

$K_A = 0.64 K_{nm} - 3.5$

J.P. Blaizot, NPA591 (1995) 435

Difference Technique to isolate the GMR

Inelastic α scattering



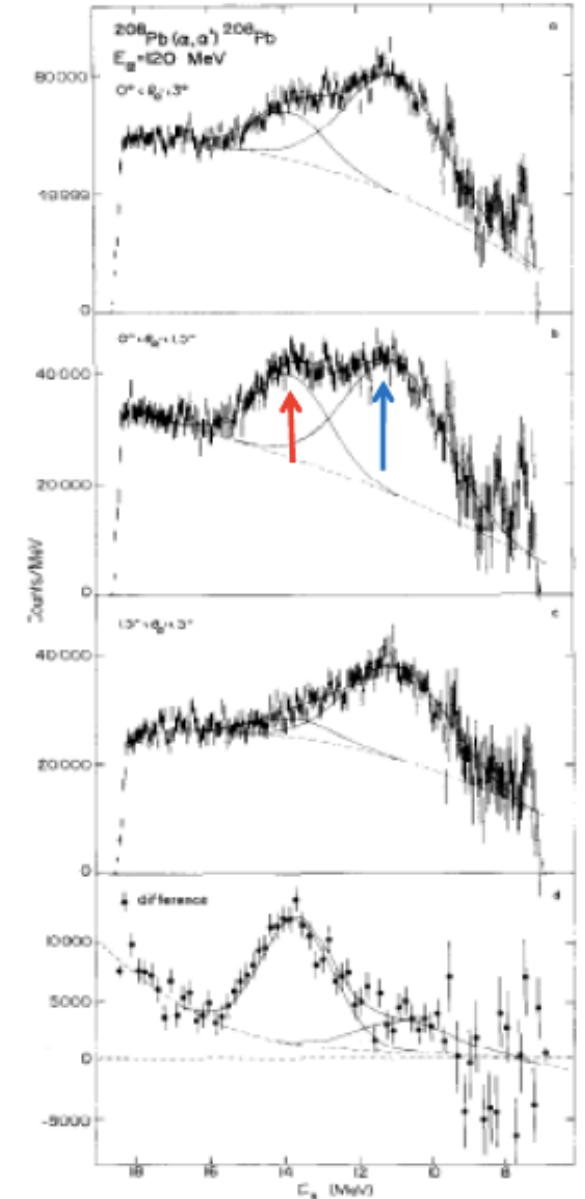
Difference of spectra

$$0^\circ < \theta_{\alpha'} < 3^\circ$$

$$0^\circ < \theta_{\alpha'} < 1.5^\circ$$

$$1.5^\circ < \theta_{\alpha'} < 3^\circ$$

Difference



Excitation of ISGDR in ^{208}Pb

From GMR data on ^{208}Pb and ^{90}Zr ,

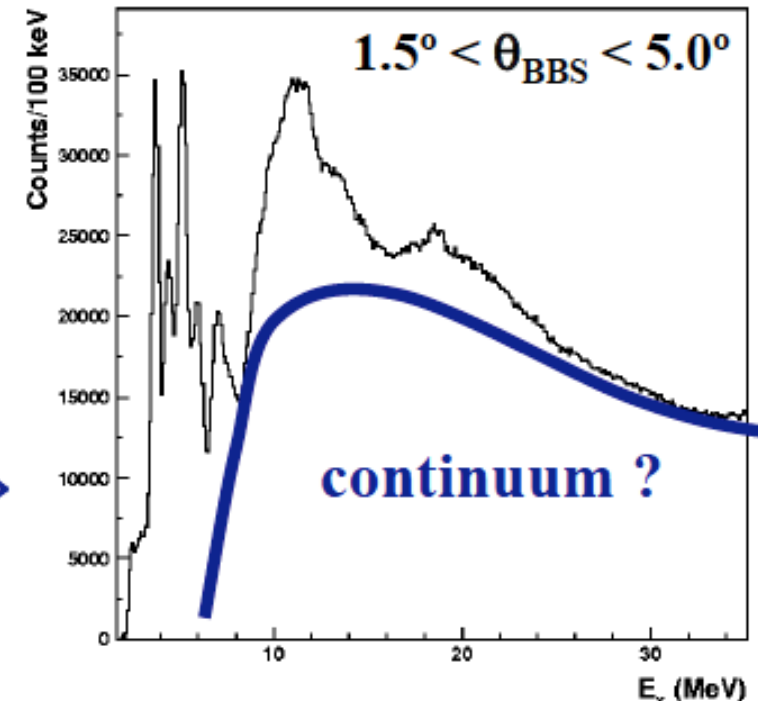
$$K_{\infty} = 240 \pm 10 \text{ MeV}$$

This number is
consistent
with both
ISGMR and
ISGDR Data

In ^{208}Pb located around 22 MeV and width of 4 MeV
 $L=1$ angular distribution peaks close to a scattering
angle of 0°

- Difficult to identify in nuclear
continuum and rides on
instrumental background

Singles $^{208}\text{Pb}(\alpha, \alpha')$ 200 MeV $\Rightarrow$



- There has been much progress in understanding ISGMR & ISGDR macroscopic properties

Systematics: E_x , Γ , %EWSR

$$\Rightarrow K_{nm} \approx 240 \text{ MeV}$$

- Sn nuclei are softer than ^{208}Pb and ^{90}Zr .
- Recently, Microscopic Structure for a few nuclei
CRPA has some success in ^{208}Pb & ^{58}Ni but fails badly in ^{116}Sn & ^{90}Zr .
- Possible observation quadrupole compression mode, i.e. overtone of ISGQR

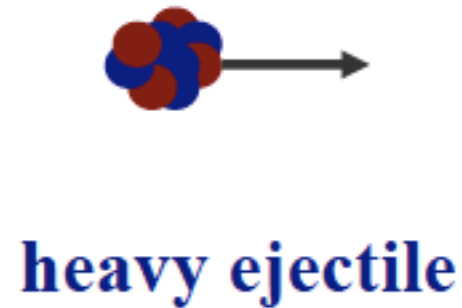
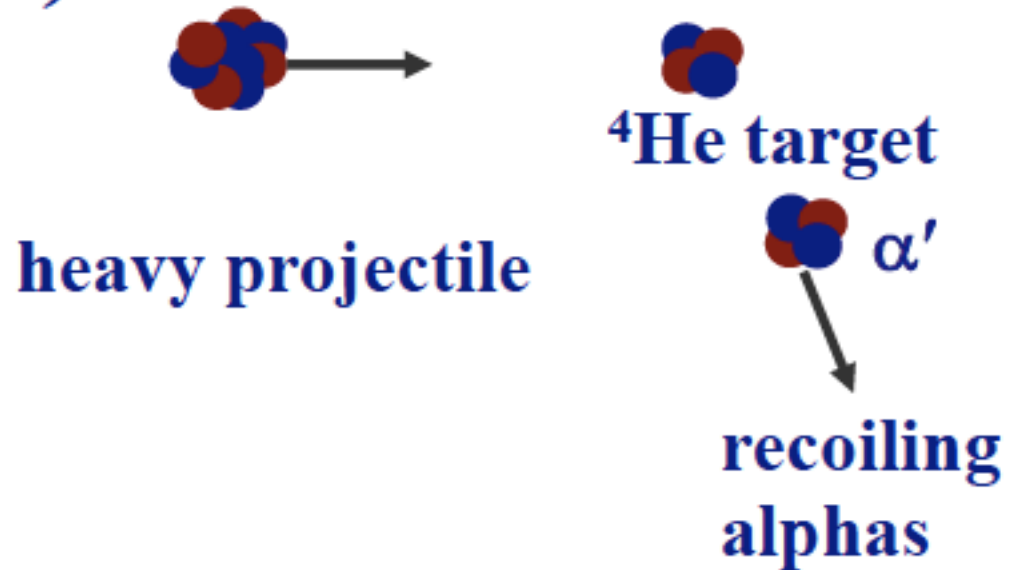
And in
nuclei
far from
Stability?

GMR studies with radioactive beams

(beam energies of 50-100 MeV/u are needed!)

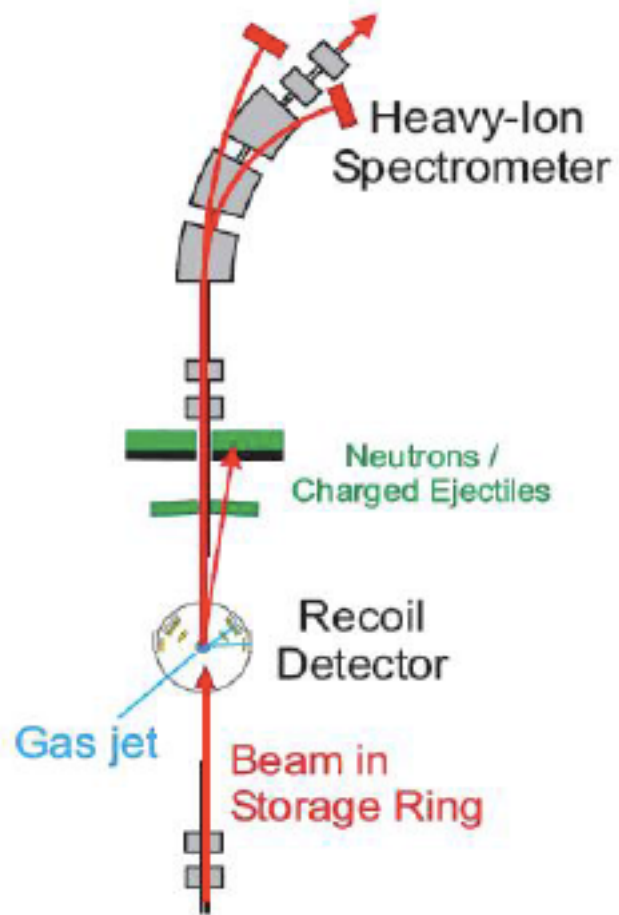
**Approach (at FAIR):
measure the recoiling alphas**

(α, α')



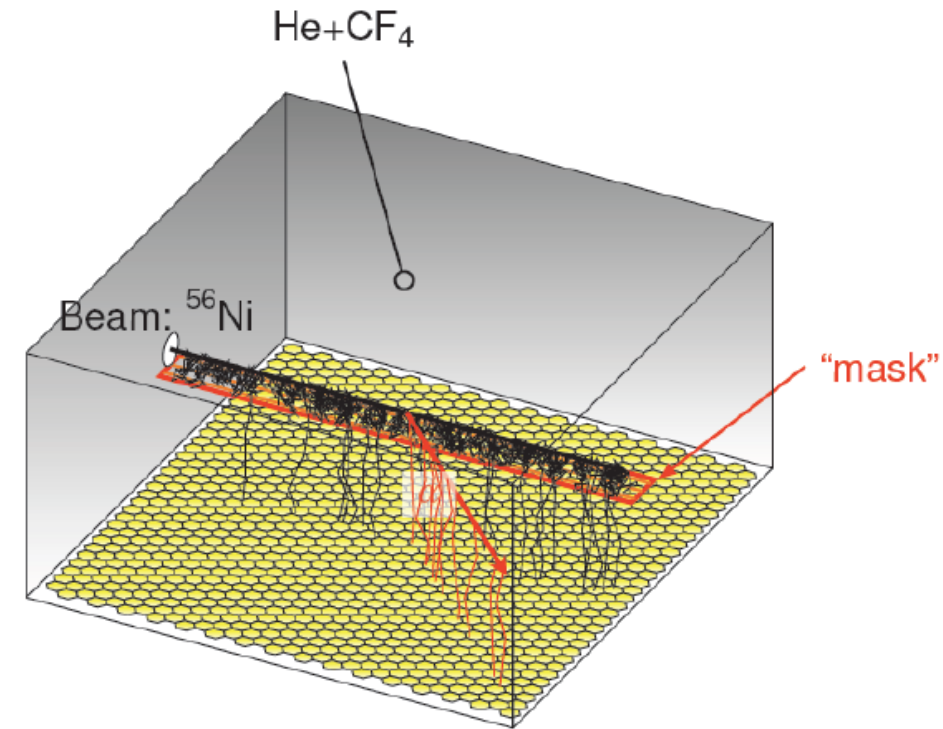
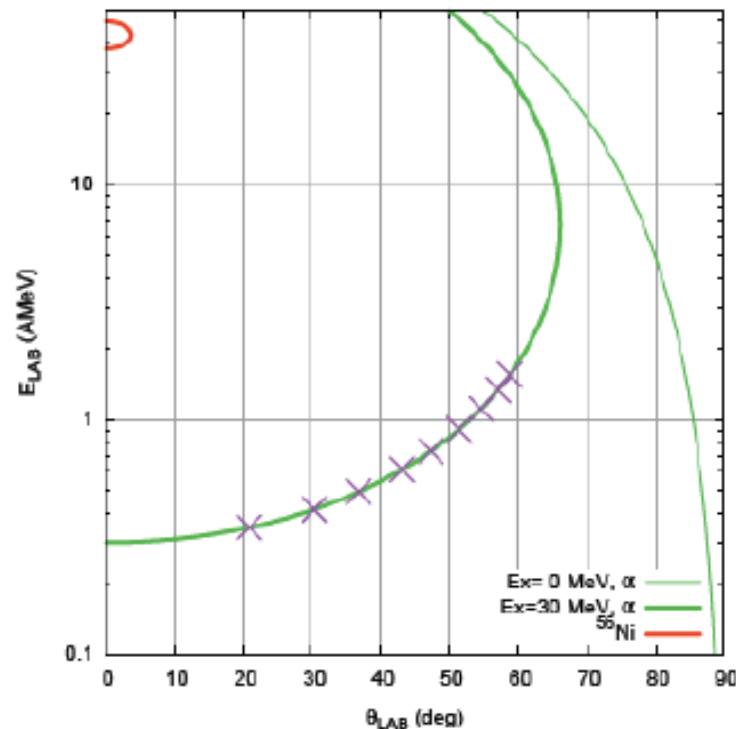
**Inconvenience:
difficulty to detect the low-
energy alphas**

Two approaches



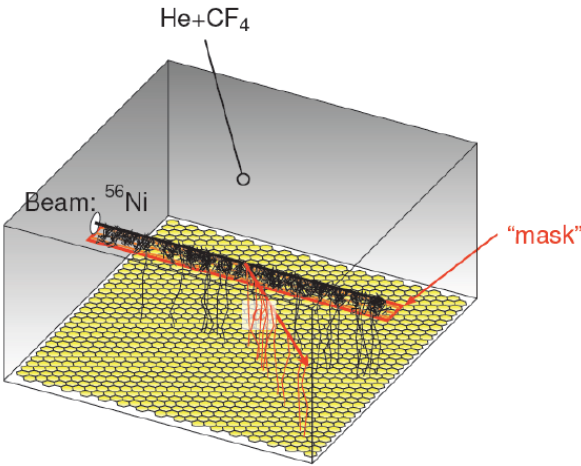
Storage ring

Measure angles and
few hundreds keV
alpha



Active target

results



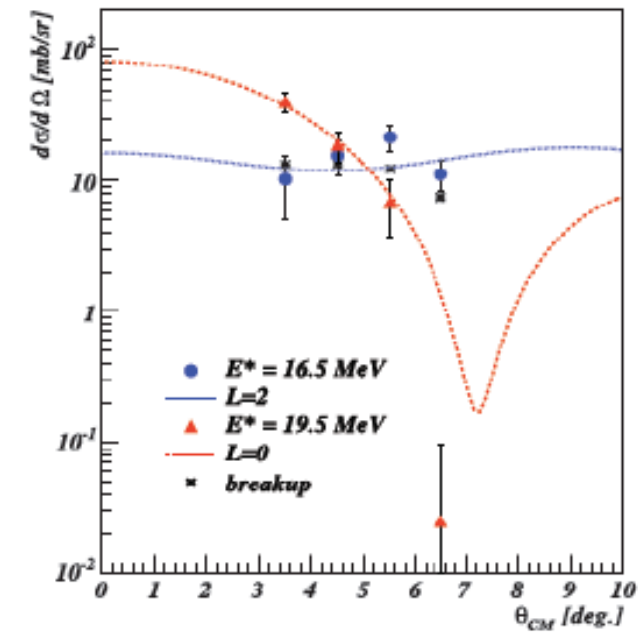
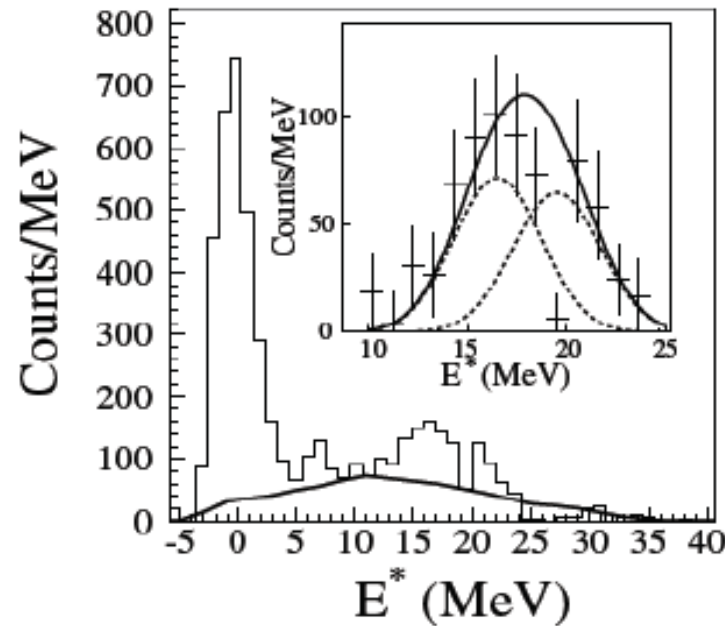
ISGMR in unstable nuclei

^{56}Ni : active target MAYA filled with deuterium gas
bombarding energy of 50A MeV

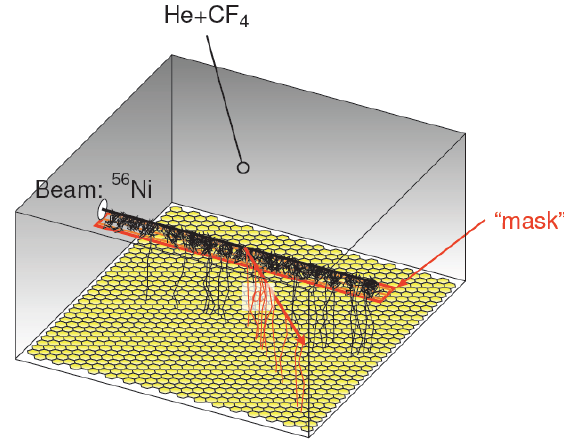
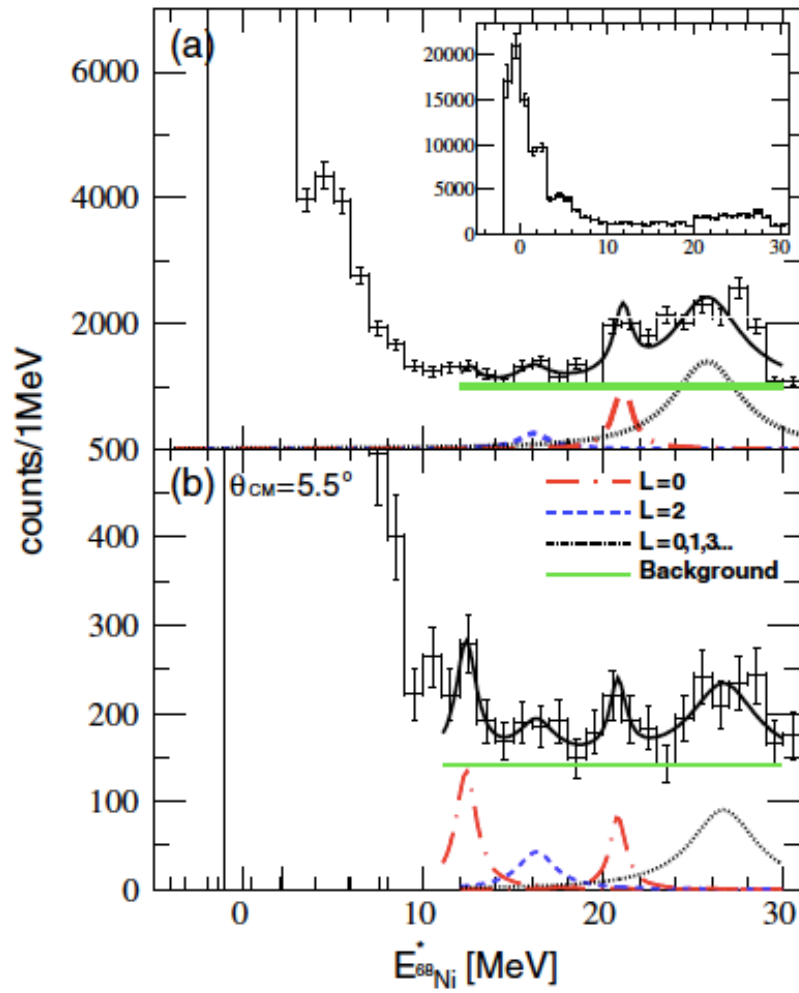
C. Monrozeau *et al.* @ GANIL

Experiments at GANIL

Phys. Rev. Lett. 100, 042501
(2008).



Results for GMR in ^{68}Ni

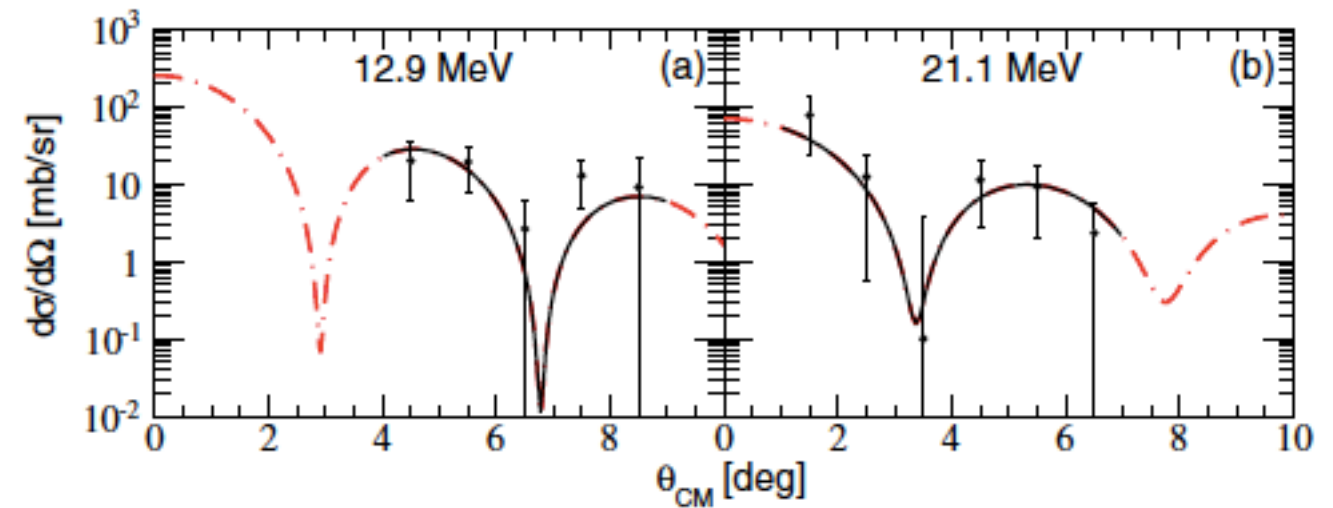


The isoscalar giant monopole resonance (ISGMR) centroid was determined to be

21.1 p/m 1.9 MeV

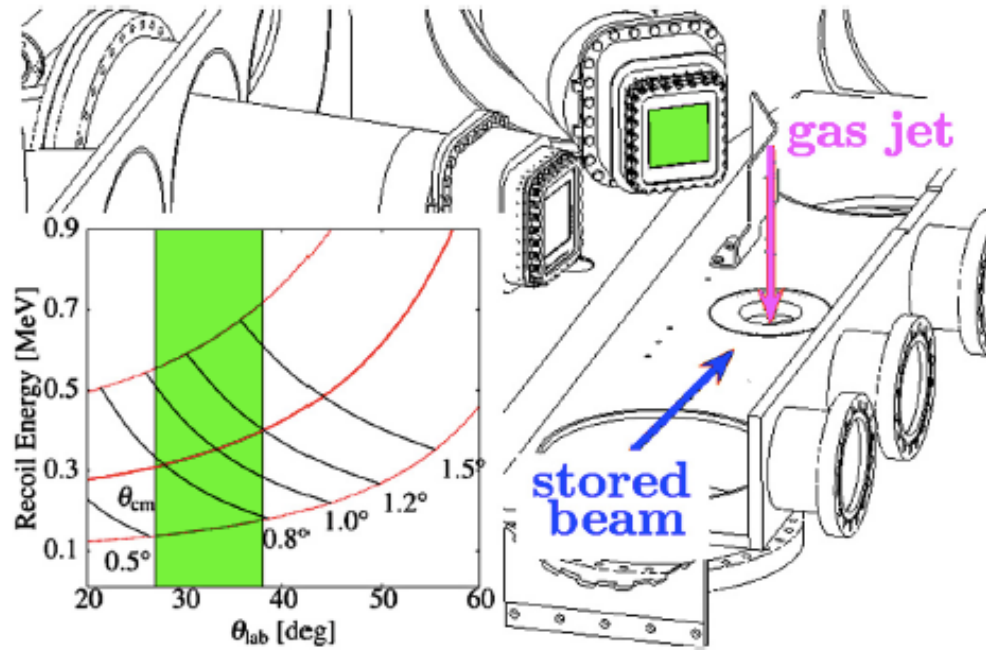
and indications for a soft monopole mode are provided for the first time at 12.9 p/m 1.0 MeV.

PRL 113, 032504 (2014)

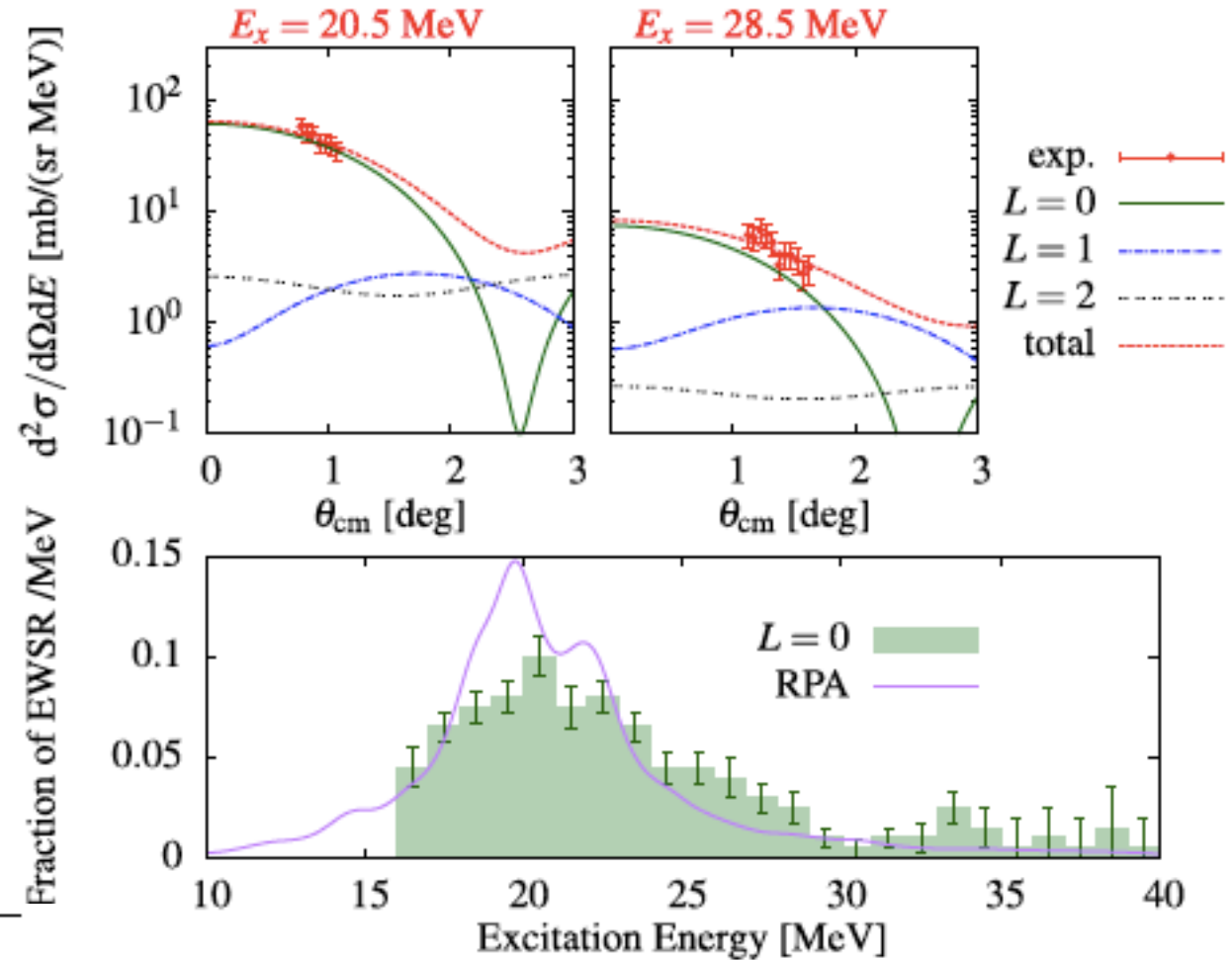


Measurements with the storage ring

The feasibility study using the ^{58}Ni stable nucleus beam



Reference	Centroid [MeV]	Width _{RMS} [MeV]	EWSR [%]
this work	20.5(6)	4.6(6)	79 ⁺¹² ₋₁₁
[9]	19.20 ^{+0.44} _{-0.19}	4.89 ^{+1.05} _{-0.31}	85 ⁺¹³ ₋₁₀
[32]	20.30 ^{+1.69} _{-0.14}	4.25 ^{+0.69} _{-0.23}	74 ⁺²² ₋₁₂
[26]	19.9 ^{+0.7} _{-0.8}	—	92 ⁺⁴ ₋₃



The collective response of nuclei

Electric Giant Resonances

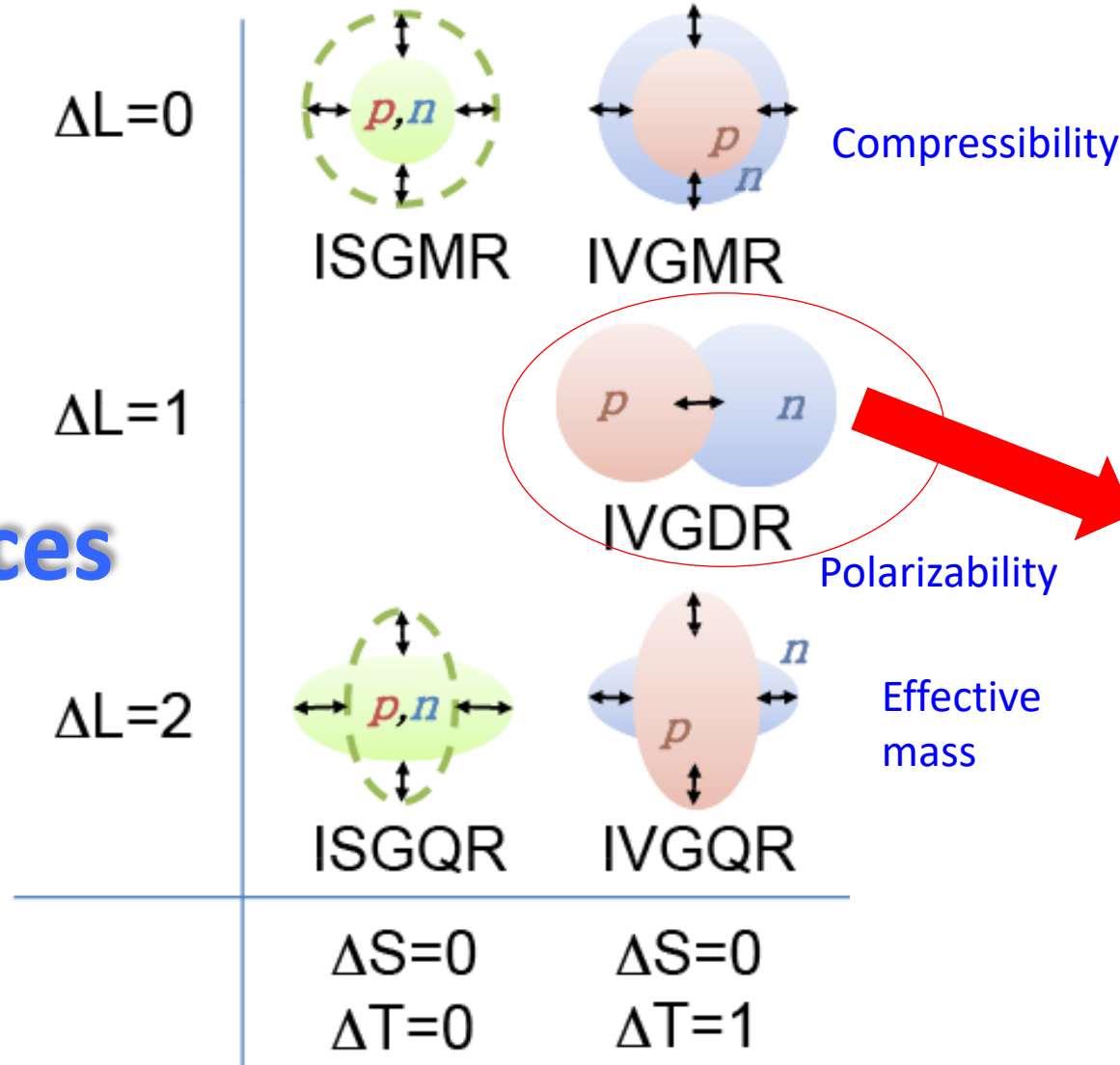
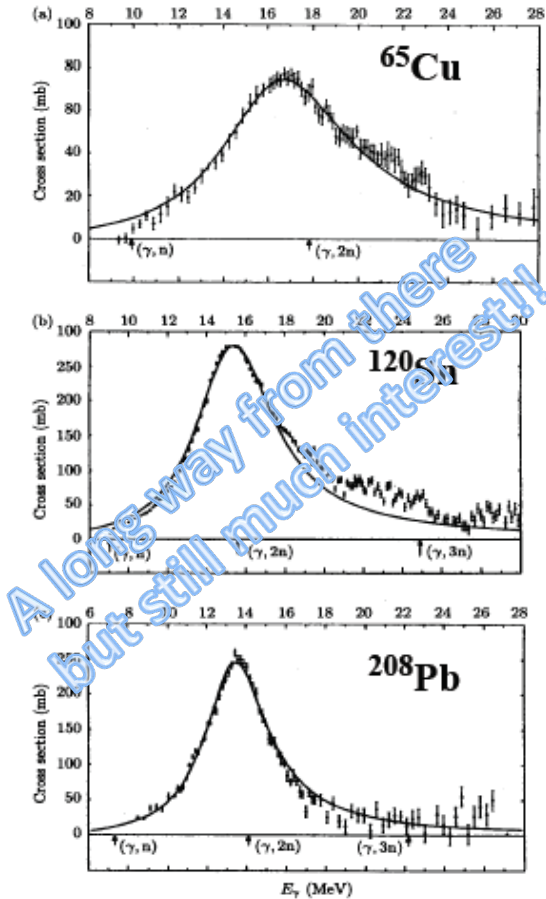


Photo-neutron cross sections



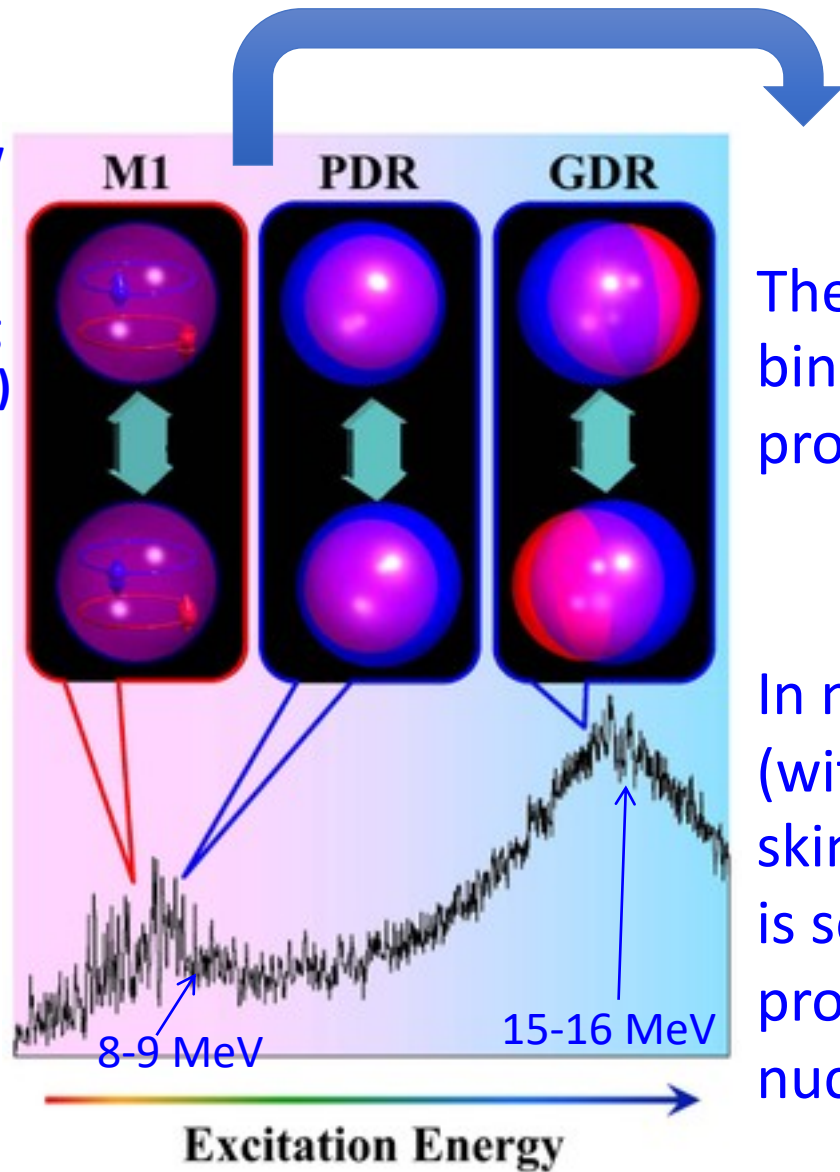
Berman and Fulz, Rev. Mod. Phys. 47 (1975) 47

A long way from there but still much interesting!!

The collective dipole response of nuclei: the details

Dipole strength at low energy below the particle binding energy (pygmy)

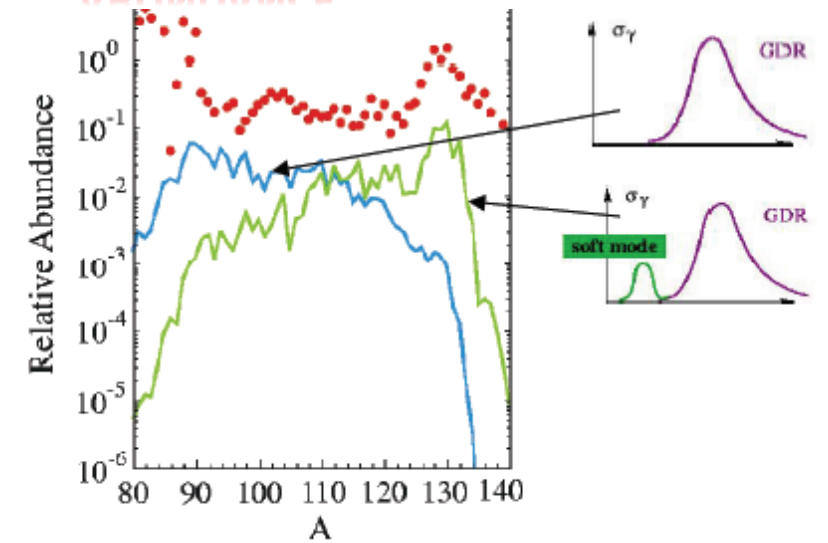
PDR (Pygmy) :
Much work is being made (exp –and theory) to understand its structure



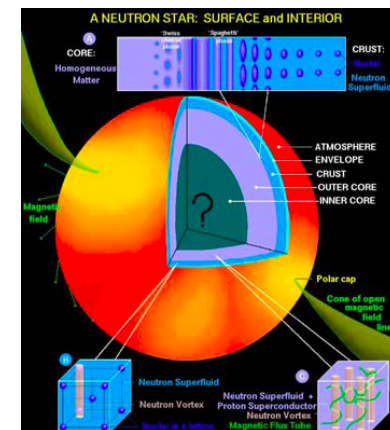
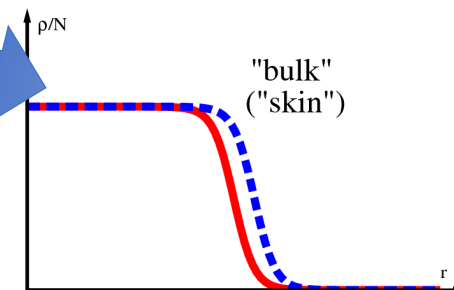
The E1 around binding affects r-process

In neutron rich (with a neutron skin) the pygmy is seen of a properties of the nucleus skin

THE R-PROCESS ABUNDANCE

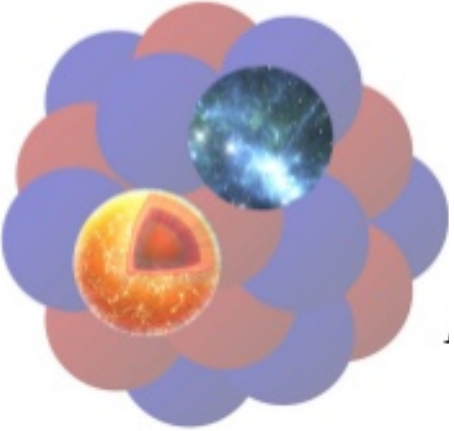


THE NUCLEAR EQUATION OF STATE



THE GDR AND THE NUCLEAR EQUATION OF STATE (EOS)

A function of nuclear density and asymmetry



EOS for Energy per nucleon

$$\frac{E}{A}(\rho, \delta) = \frac{E}{A}(\rho, 0) + S(\rho)\delta^2 + \dots$$

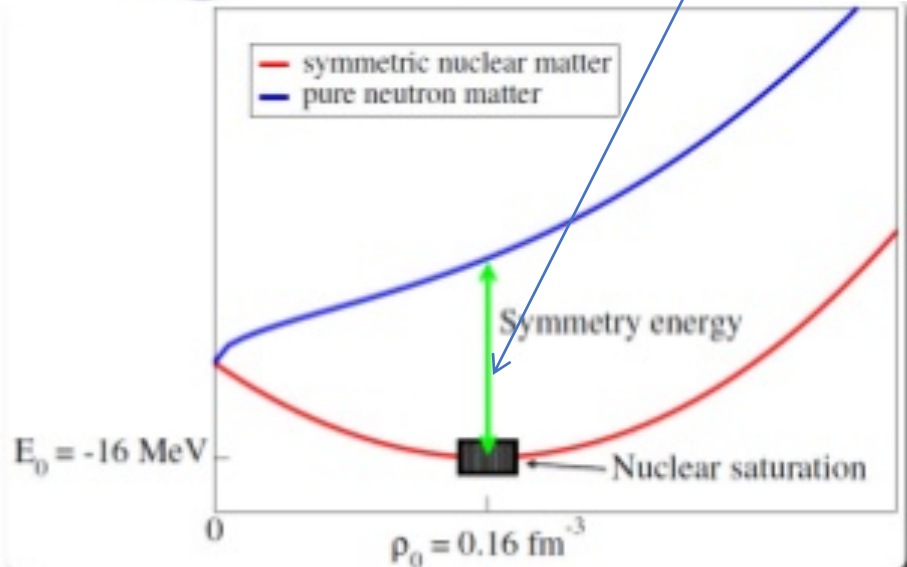
symmetry energy

$$S(\rho) = J + \frac{L}{3\rho_0}(\rho - \rho_0) + \frac{K_{sym}}{18\rho_0^2}(\rho - \rho_0)^2 + \dots$$

$$\rho(r) = \rho_n(r) + \rho_p(r)$$

$$\delta(r) = \frac{\rho_n(r) - \rho_p(r)}{\rho_n(r) + \rho_p(r)}$$

Saturation Density
 $\sim 0.16 \text{ fm}^{-3}$
 ρ_0 :



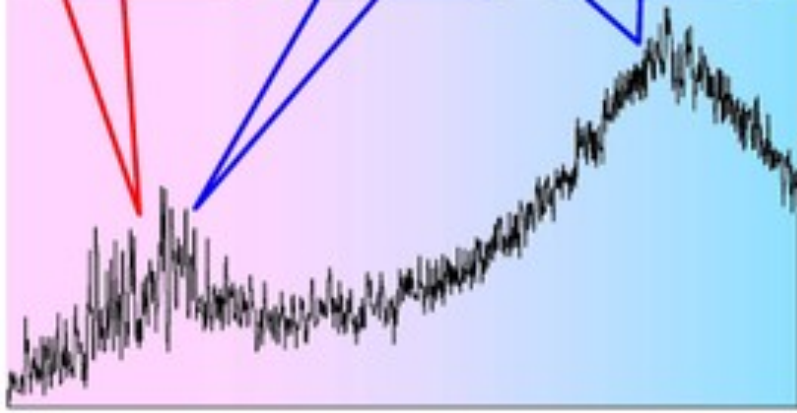
Expansion in series- L is the first derivative of an energy and this is a pressure...

$$L \propto P \propto R_{n\text{-star}}^4$$

(Baryonic Pressure)

P = pressure

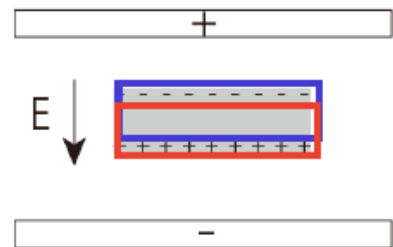
THE GDR AND THE NUCLEAR EQUATION OF STATE (EOS)



Experiments: observables related to the electric dipole response

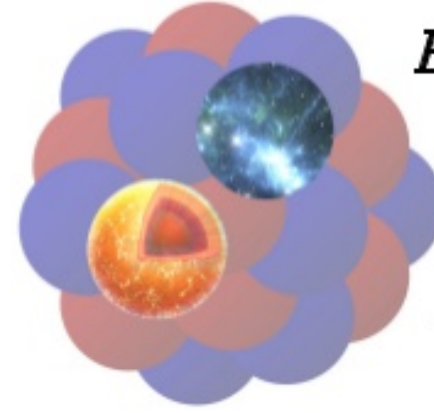
- The position and intensity of the low energy part (Pygmy)
- Polarizability (the pygmy part plays a role!)

$$\alpha_D = \frac{\hbar c}{2\pi^2} \int \frac{\sigma_{abs}}{\omega^2} d\omega$$



Electric Dipole Polarization

$$\mathbf{P} = \alpha \mathbf{E}$$



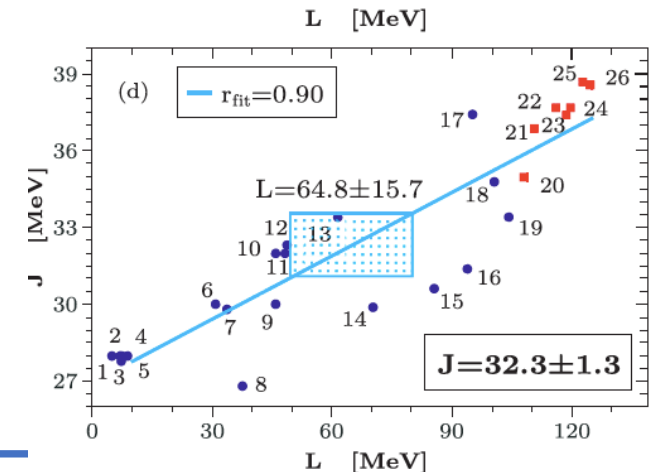
E Theoretical predictions

Correlations between different quantities

$L \rightarrow$ related to Pygmy and polarizability

$L \rightarrow$ related to the neutron skin (difference of the matter and electromagnetic radii)

L and J are related



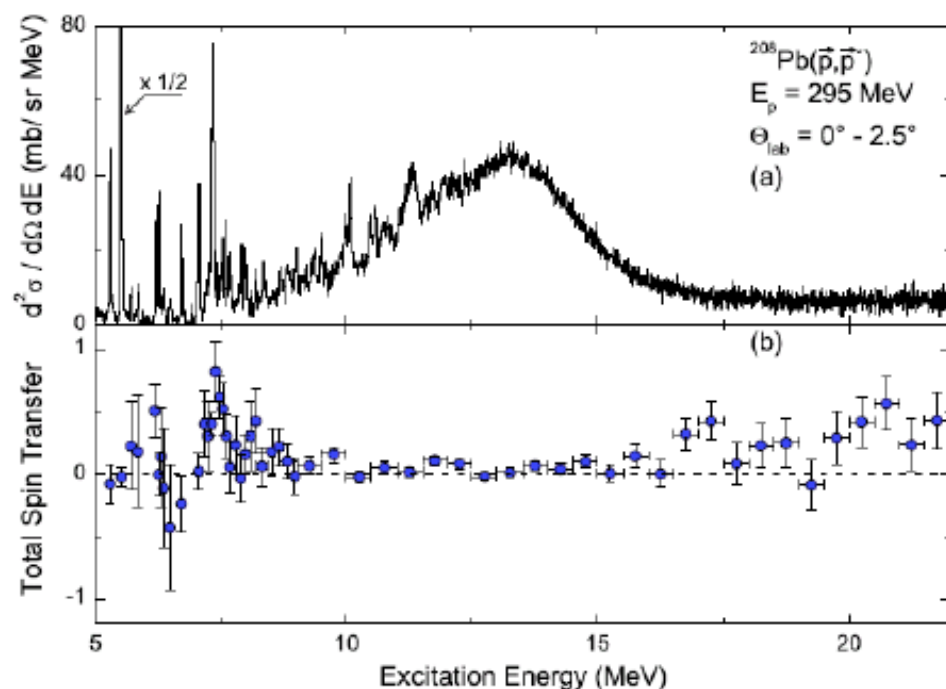
SOME RESULTS FOR ^{208}Pb

- Polarization observables at 0° $\rightarrow$ spinflip / non-spinflip separation
model-independent

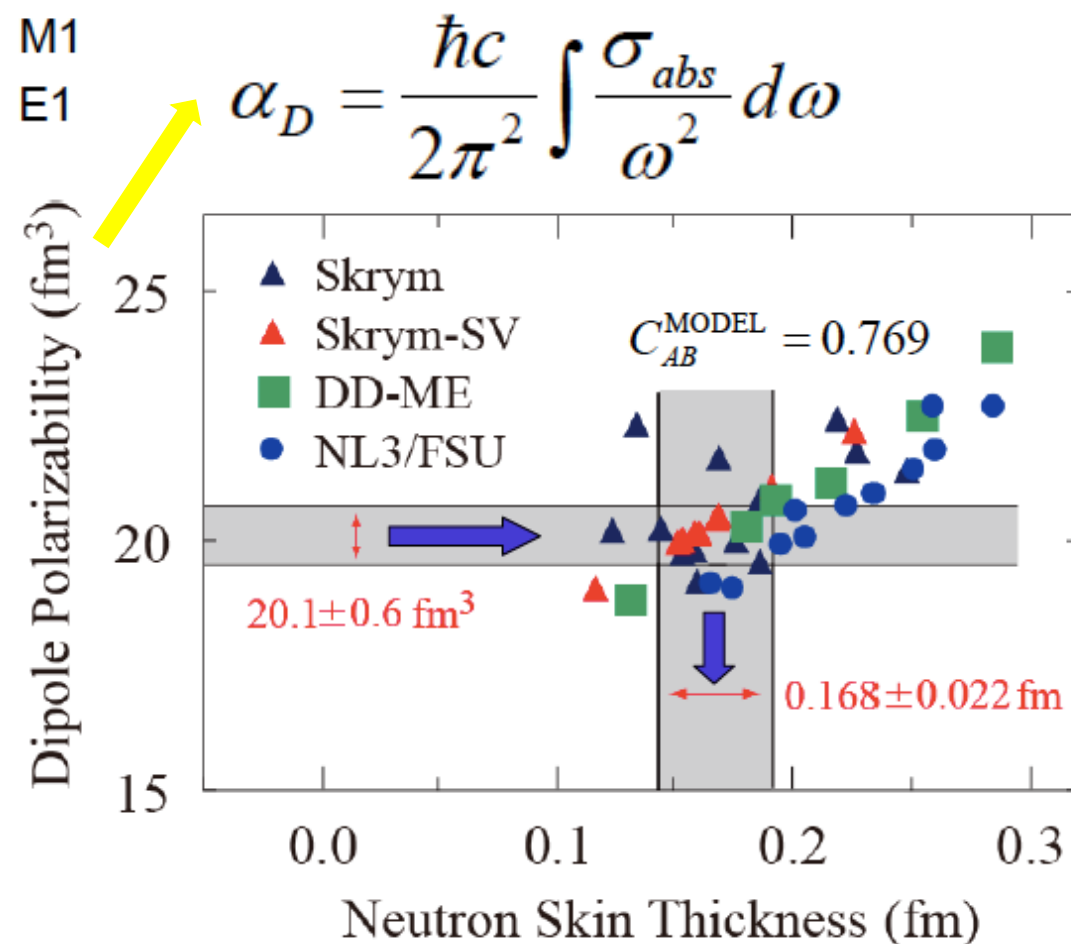
E1 and M1 decomposition

T. Suzuki, PTP 103 (2000) 859

$$\text{Total Spin Transfer } \Sigma \equiv \frac{3 - (2D_{SS} + D_{LL})}{4} = \begin{cases} 1 & \text{for } \Delta S = 1 \quad \text{M1} \\ 0 & \text{for } \Delta S = 0 \quad \text{E1} \end{cases}$$

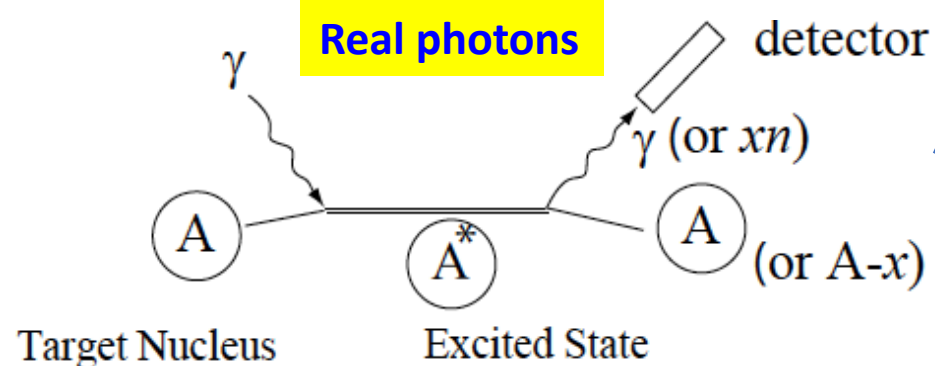


Inelastic scattering of polarized protons at 295 MeV



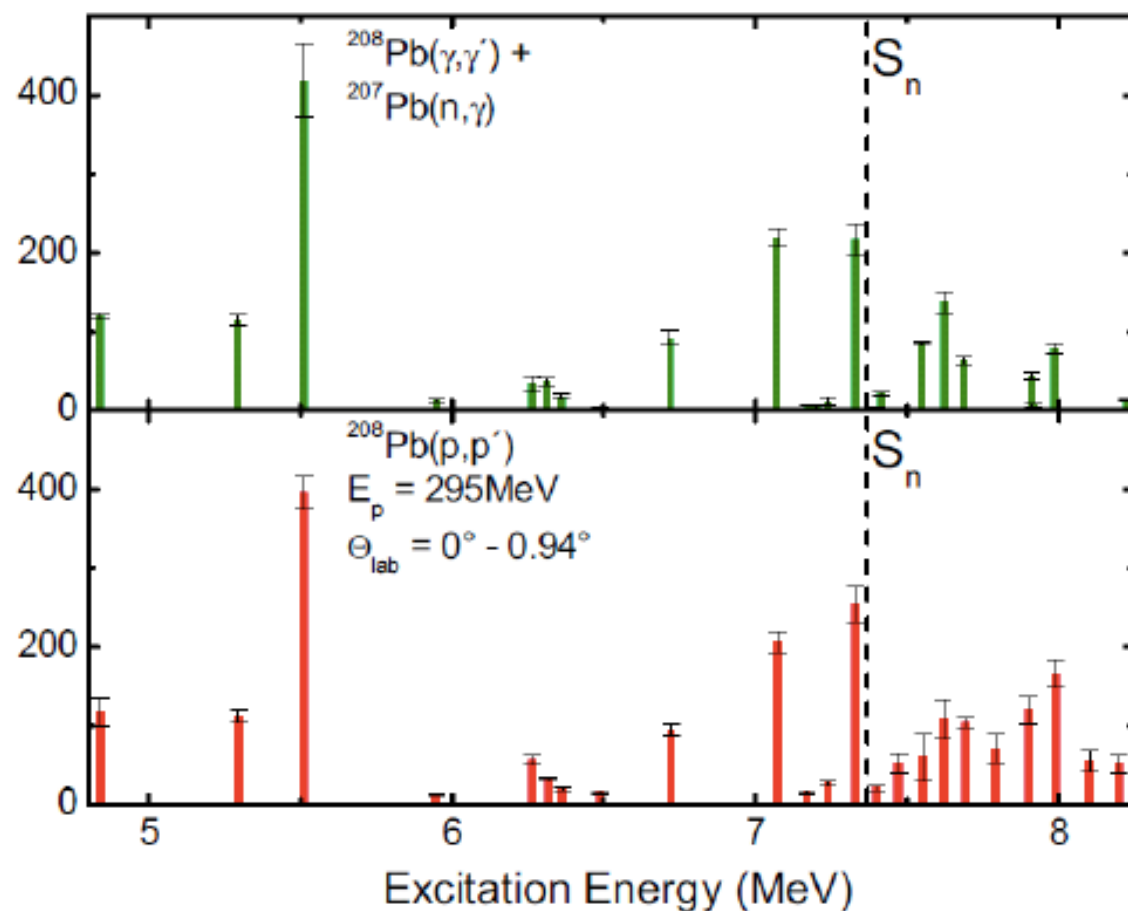
SOME RESULTS FOR ^{208}Pb

Real photons

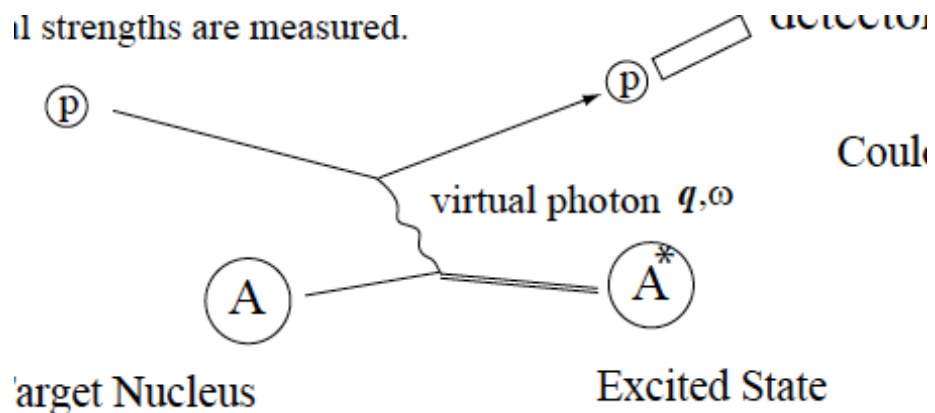


B(E1): low-lying discrete states

$B(E1) \uparrow (10^{-3} \text{ e}^2 \text{ fm}^2)$



all strengths are measured.

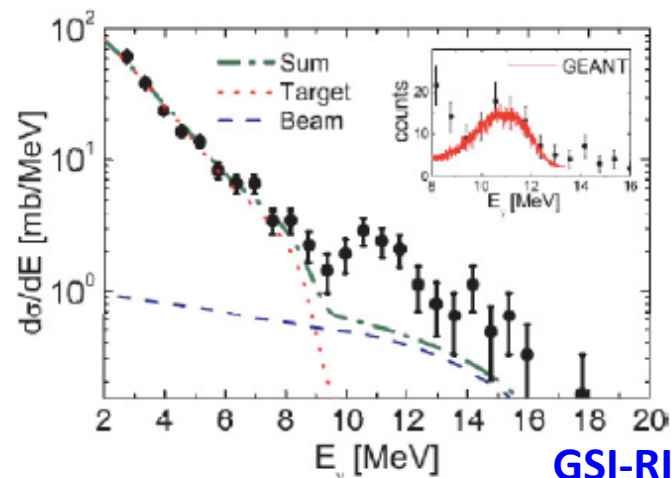
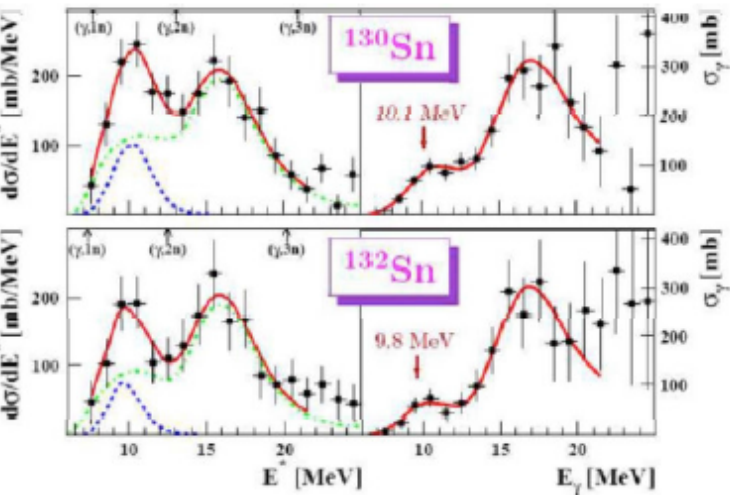


Virtual photons

Pygmy – via Coulomb excitation on exotic nuclei $^{130-132}\text{Sn}$ and ^{68}Ni

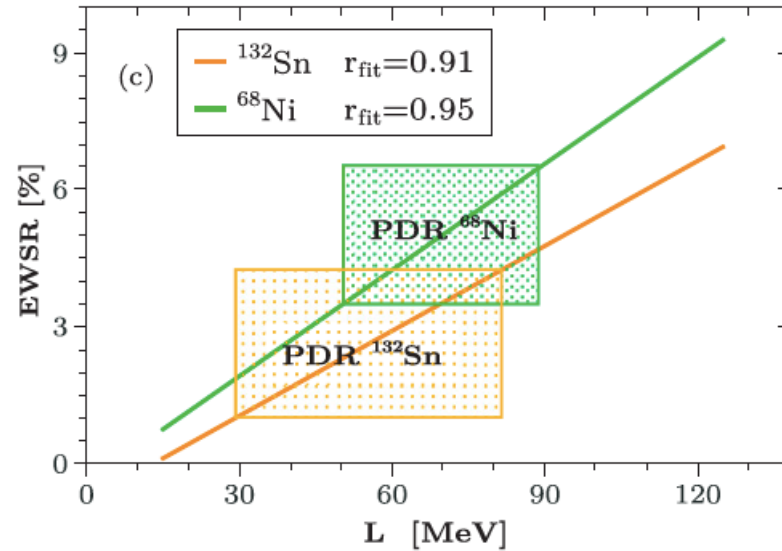
Polarizability – via Coulomb Exc. on exotic ^{68}Ni

GS- LAND



GS- RISING

Strength in the pygmy region



From the deduced L the corresponding
neutron skin is

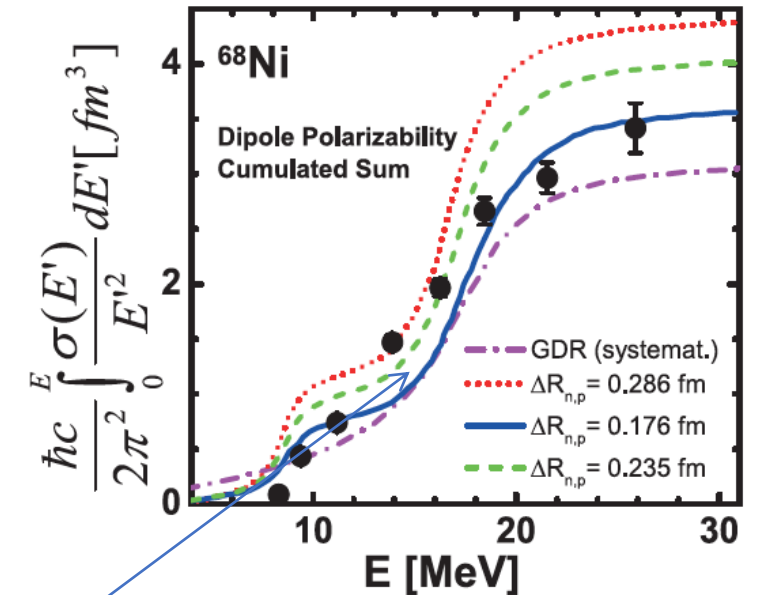
0. 200 p/m 0. 015 fm for ^{68}Ni

0. 258 p/m 0. 024 fm ^{132}Sn

0. 194 p/m 0. 024 for ^{208}Pb

Polarizability

GS- LAND

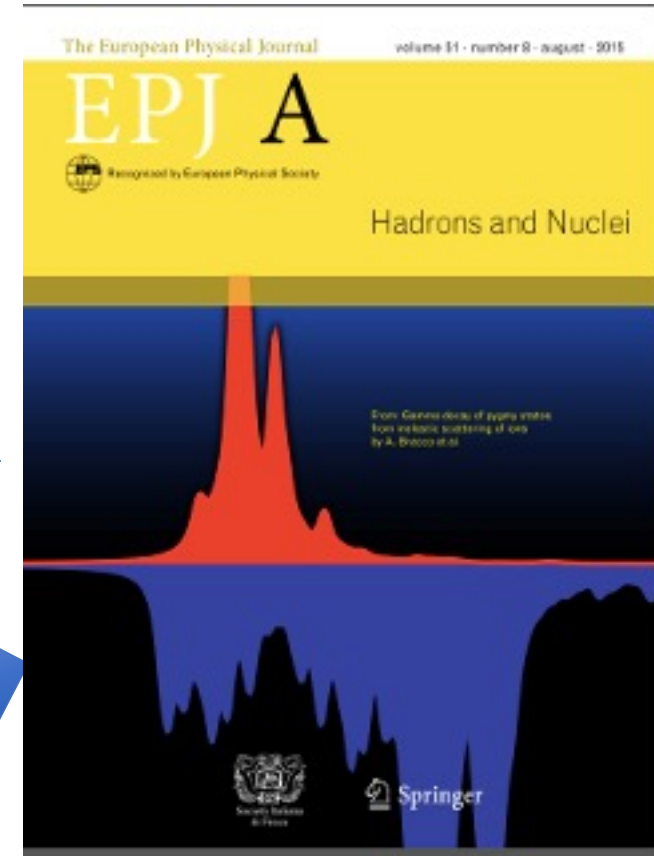


0. 176 fm for ^{68}Ni

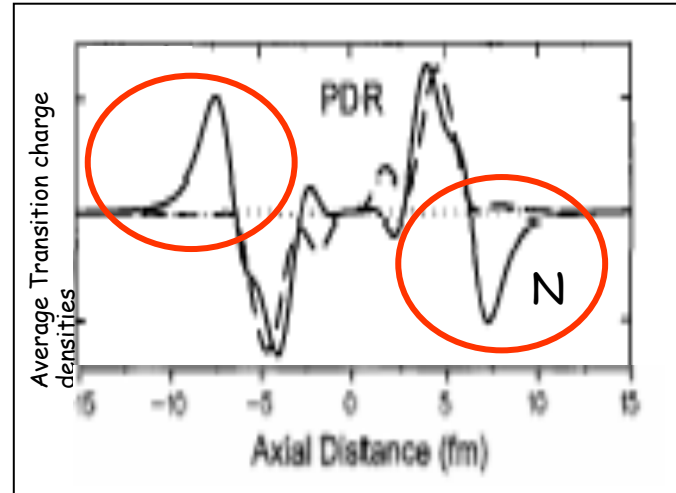
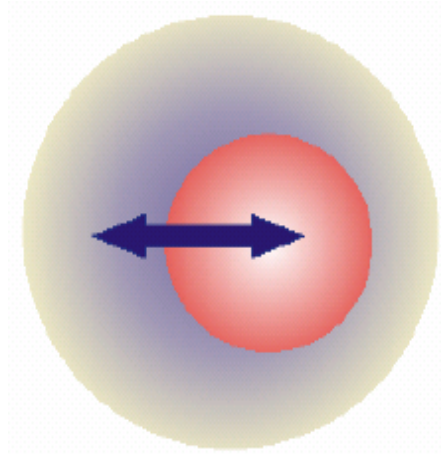
0.17 fm

STUDY OF PYGMY STATES

- **NATURE OF THE STATE**
- **CONSTRAIN BETTER THE CALCULATIONS**
- **GET INFORMATION ON THE WAVE-FUNCTION**

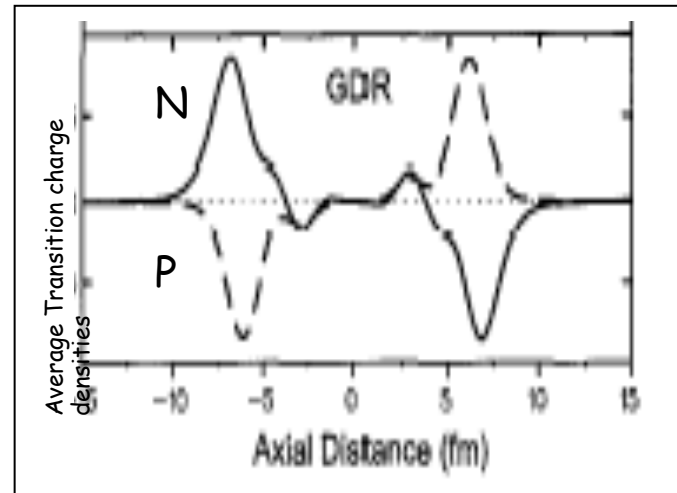
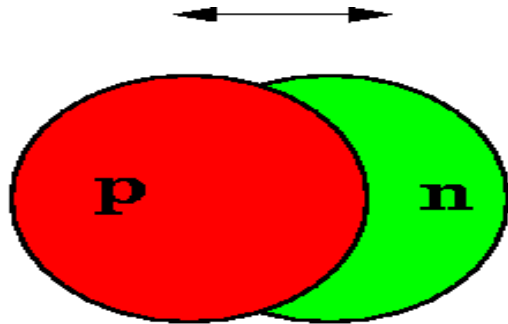


We need to understand the structure of the low lying E1 strength !



Typical transition density of protons and neutrons for 1^- states at excitation energy in the region of the tail of the GDR

There are regions in which only neutrons contribute

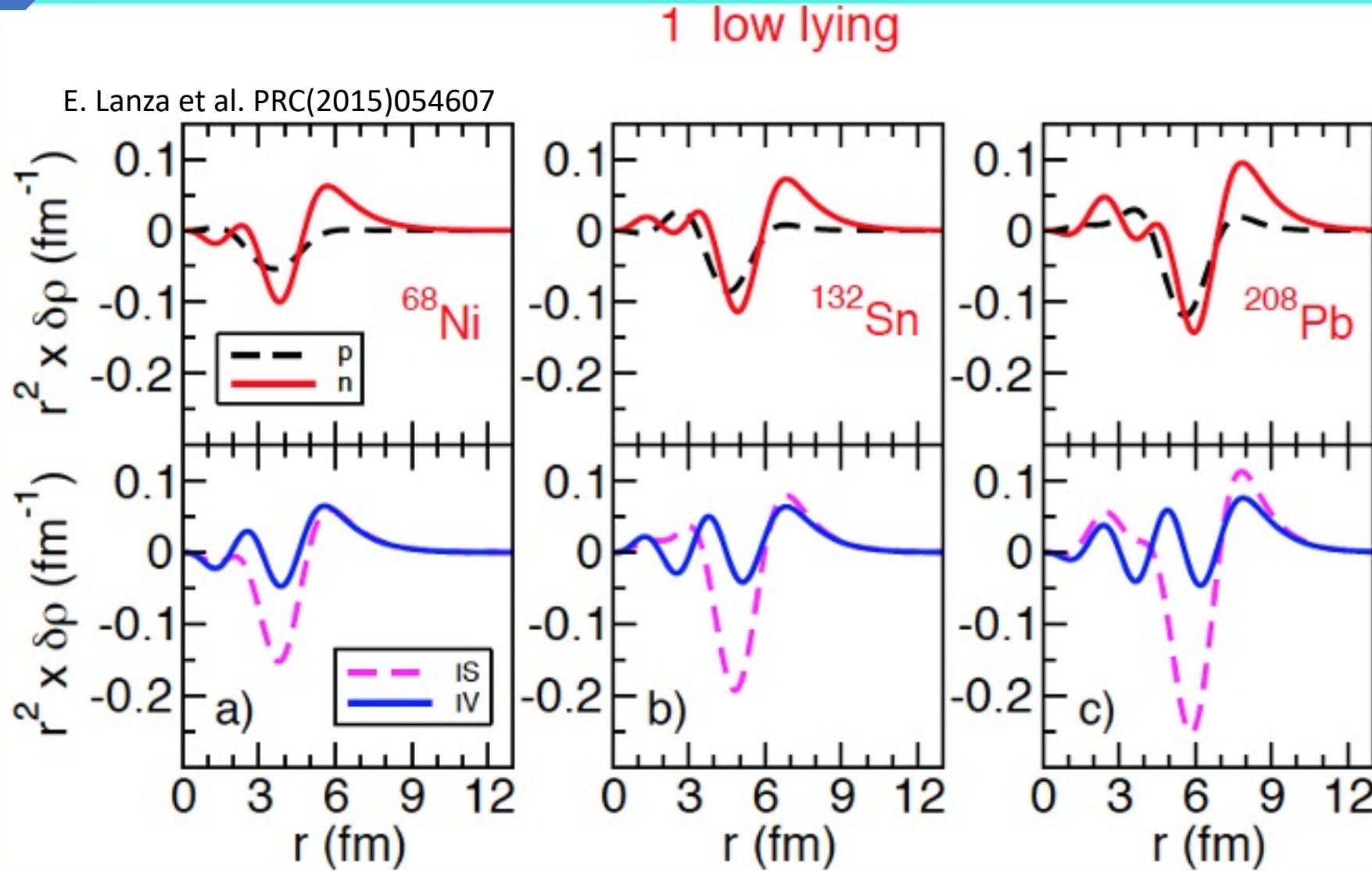


Typical transition density of protons and neutrons for 1^- states at excitation energy in the region of the GDR centroid

Neutron and proton contributions extend in the same spatial regions and proton and neutron oscillations are out of phase!

The structure of the low lying 1^- states: Isovector and Isoscalar components

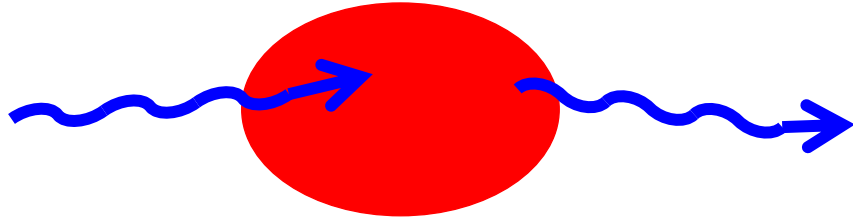
E. Lanza et al. PRC(2015)054607



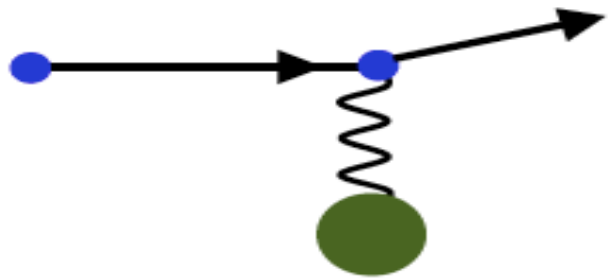
Transition densities of protons and neutrons for 1^- states in the PDR region: similar features for different

Isoscalar and isovector components of the transition densities

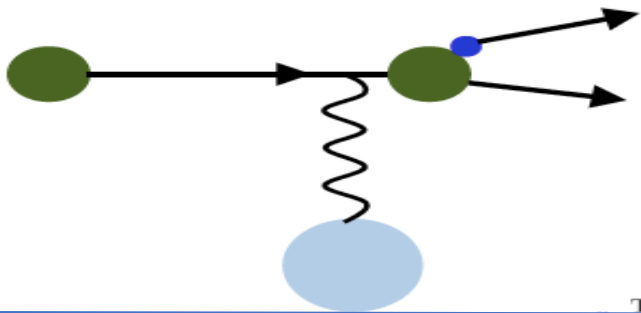
There is a need for a variety of experimental tools



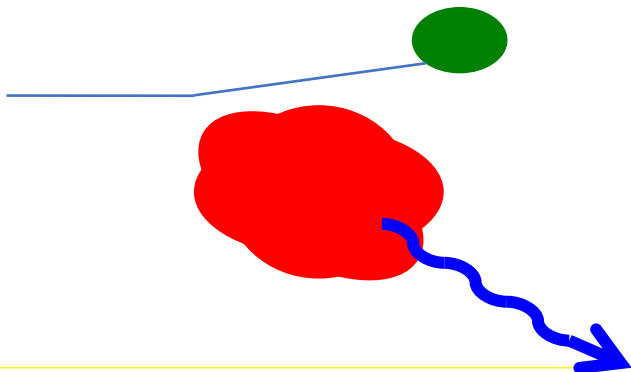
(γ, γ') real photons
on stable nuclei
Probing the entire nuclear volume



(p, p') virtual photons at $E_{\text{beam}} > 200 \text{ MeV}$
on stable nuclei
Probing the entire nuclear volume
Work at Osaka RCNP



Coulomb excitation with radioactive beams virtual photons
Exotic nuclei
E1 excitation at high energy (GSI, MSU) with particle and
gamma decay from the projectile



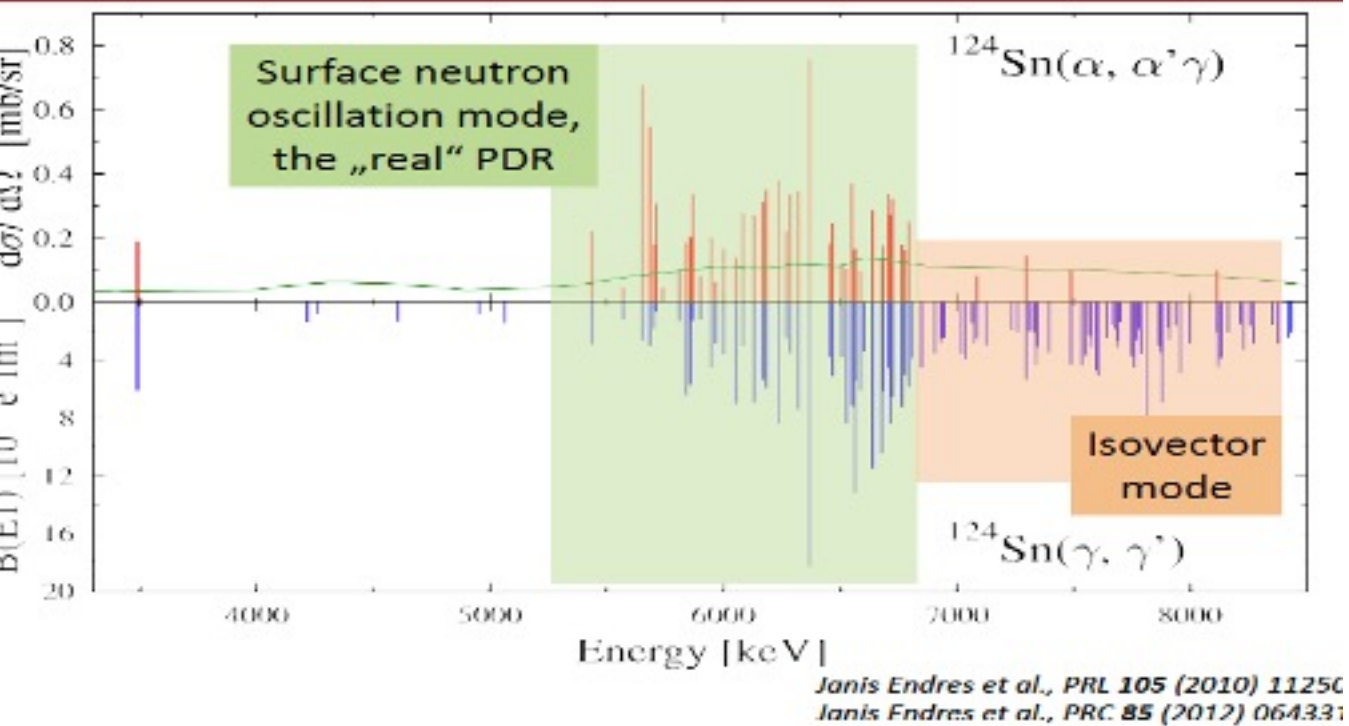
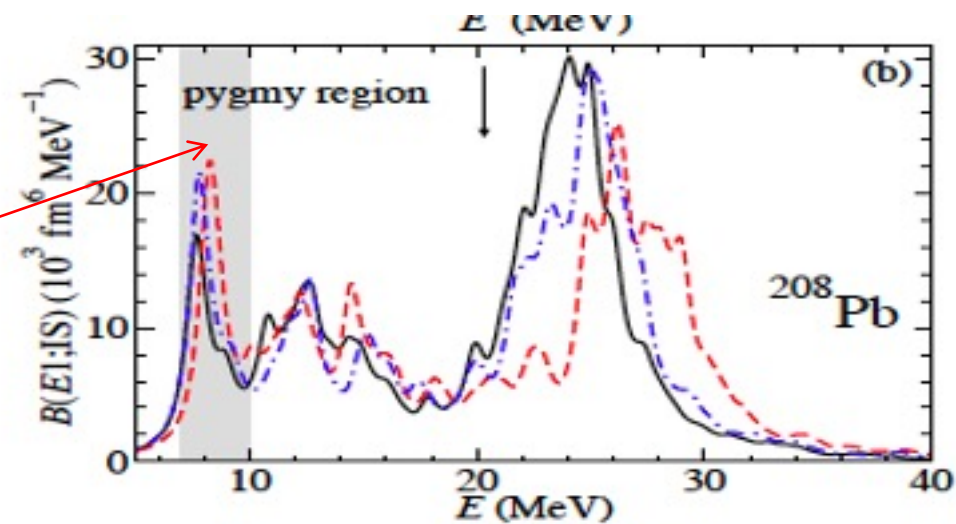
$(\alpha, \alpha' \gamma)$ or $(^{17}\text{O}, ^{17}\text{O}' \gamma)$
on stable nuclei
Probing the nuclear surface mainly

KVI and LNL

Do we need different probes ?

Isoscalar or isovector?

In ^{208}Pb
4.3-5.5 % IS EWSR strength
at 9 MeV
Theory from Rocha-Maza, Colo'



IS nature of the PDR due to outermost nucleons (neutrons in a neutron-rich nucleus).

The r_{np} is correlated with J and L.

Interesting to study the properties of the neutron skin

First work on $(\alpha, \alpha'\gamma)$ –
PLB278, 423 (1992).

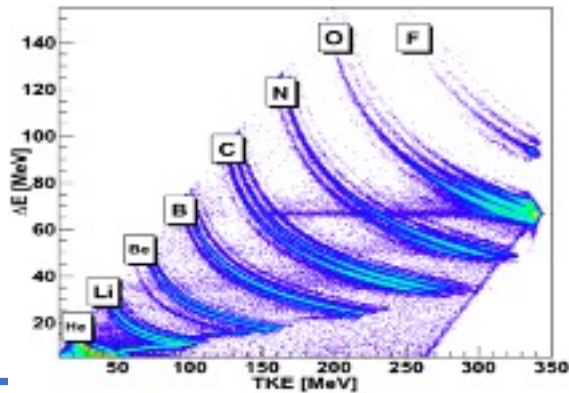
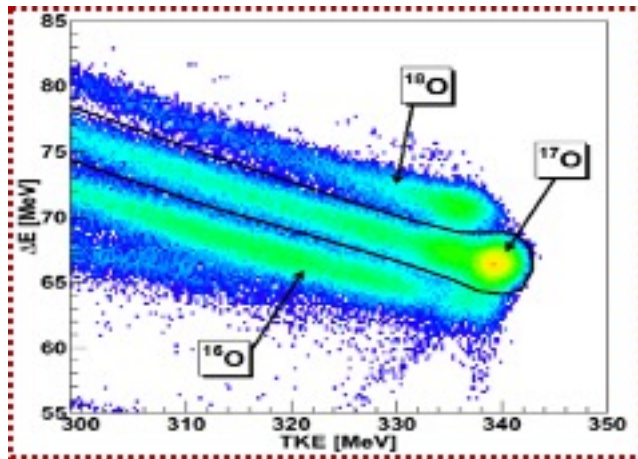
Recent measurements with ^{17}O

Using AGATA at LNL

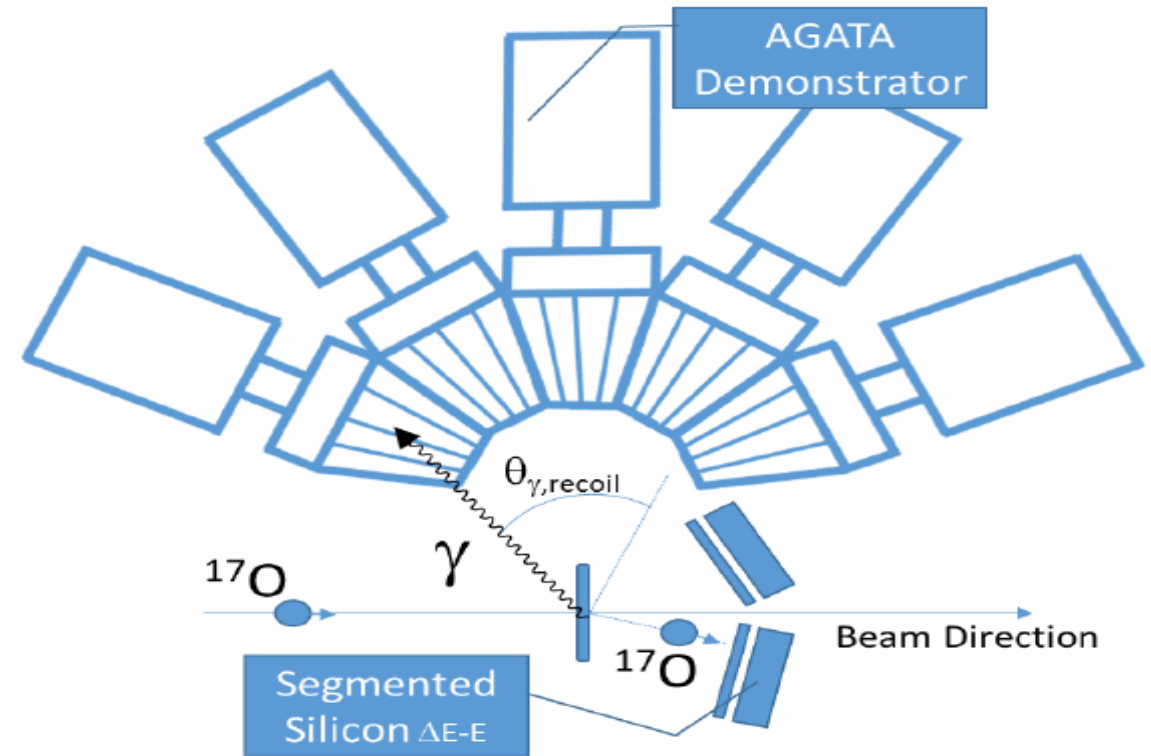
The AGATA experiments

- Large cross-section for the population of the giant resonance region
- ^{17}O is loosely bound ($S_n = 4.1$ MeV)
- Clean removal of projectile excitation

^{17}O @ 20 MeV/u on different targets
+ γ -rays in coincidence



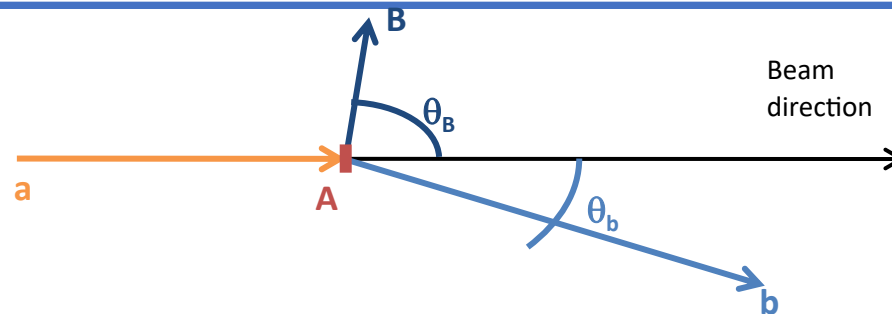
Silicon Detectors, D. Mengoni NIMA 764(2014)241



Doppler correction....AGATA performances

Projectile-like (^{16}O , $v/c \sim 20\%$)

– > 500 keV @ 5 MeV



Target-like recoils ($v/c \sim 0.5\%$)

– ~ 25 keV @ 5 MeV

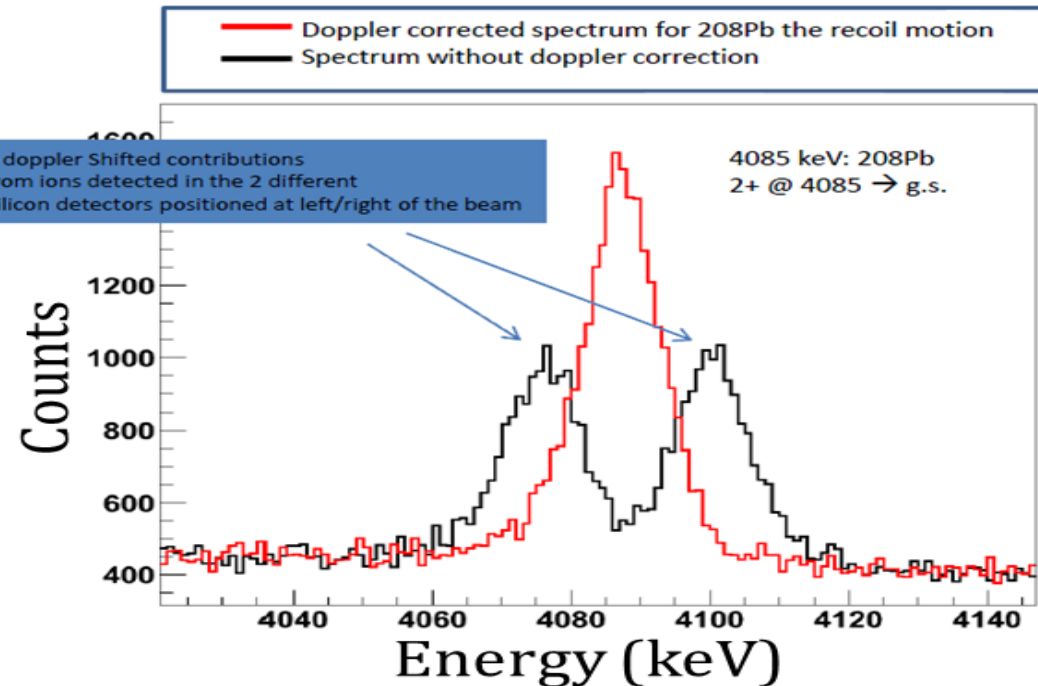
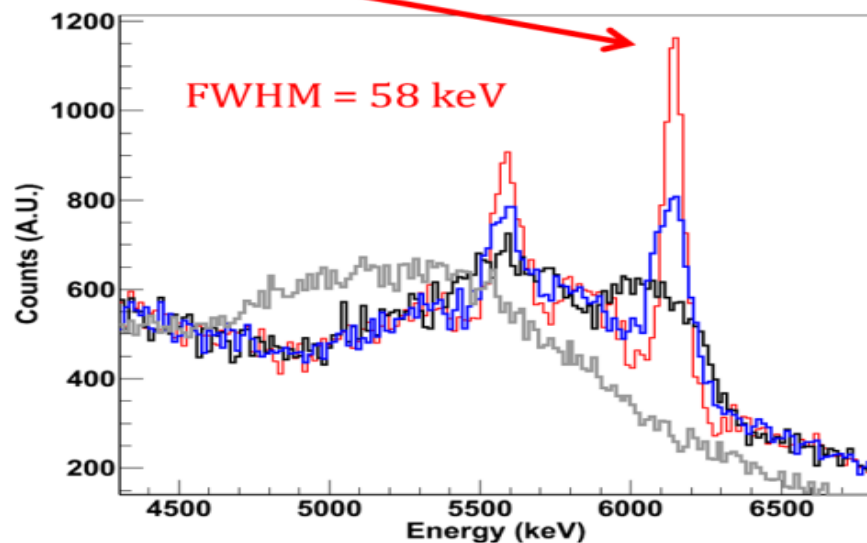
^{16}O

No Dopp Corr

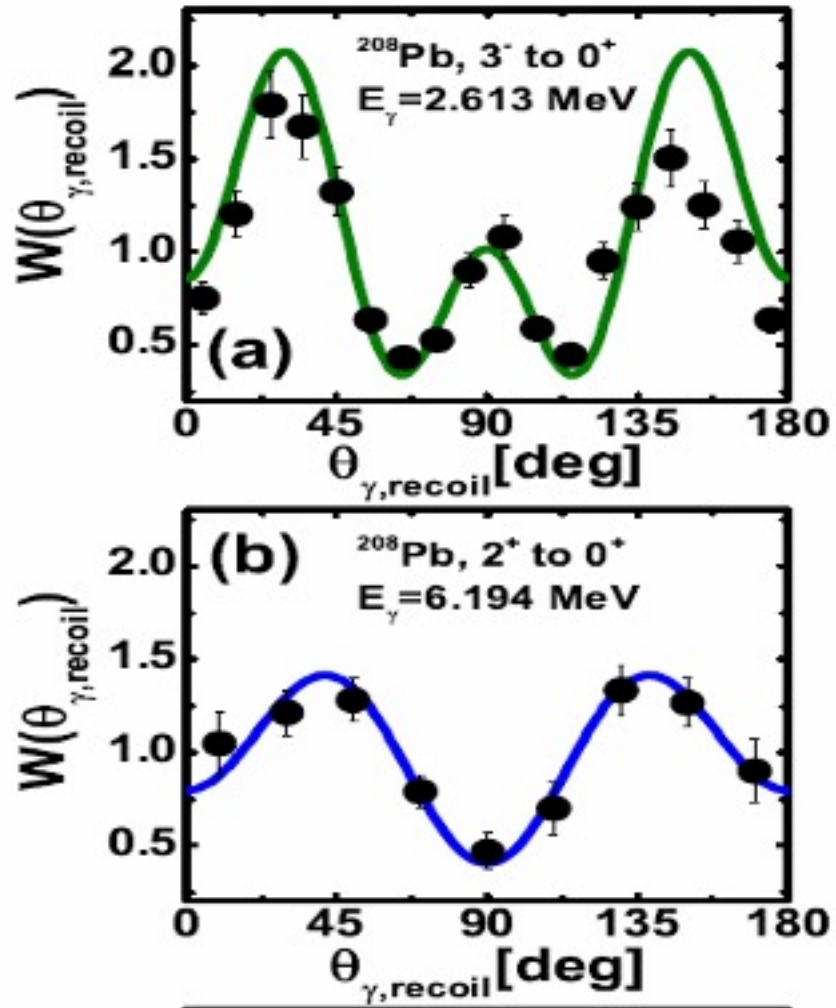
Crystal Centers

Segment Centers

PSA+Tracking

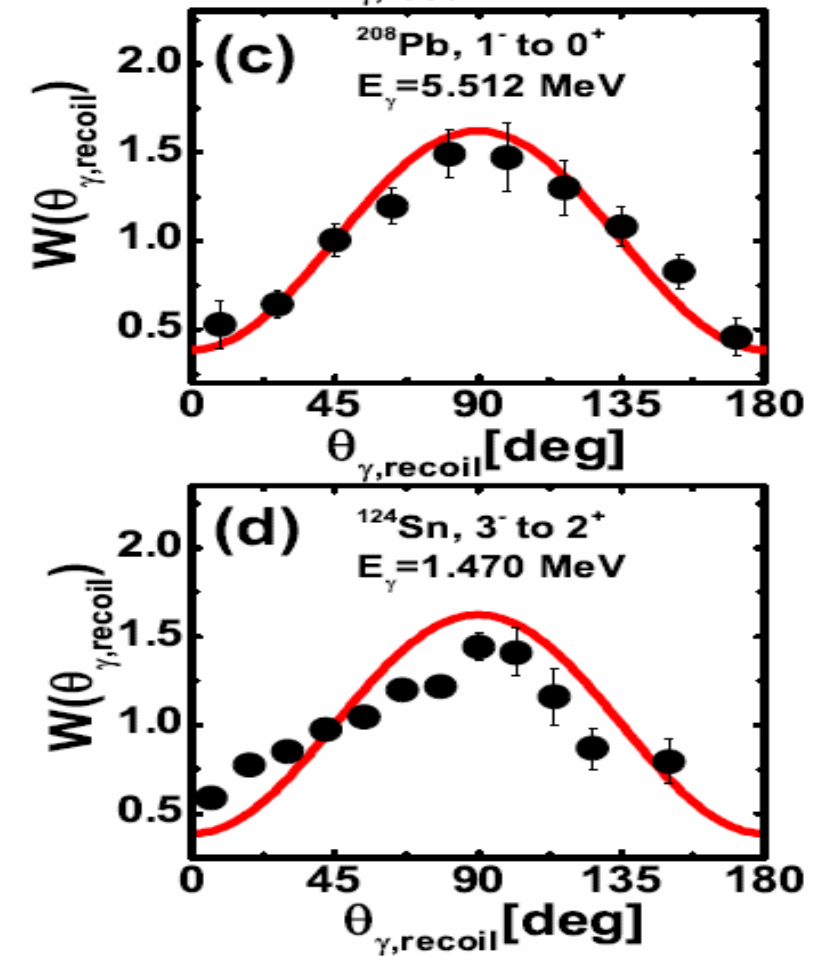


Angular distribution of γ -rays ---Some data on ^{208}Pb



*Angular Distribution
of γ 's*

*position sensitivity
of AGATA
and E- Δ E Si
telescopes (pixel type)*



Excitation of the 2^+ and 3^- states in ^{90}Zr , ^{124}Sn , ^{208}Pb

- The $B(E2)$ and $B(E3)$ known from other works*

- Calculations assumed pure isoscalar excitation namely the p and n matrix element are related by

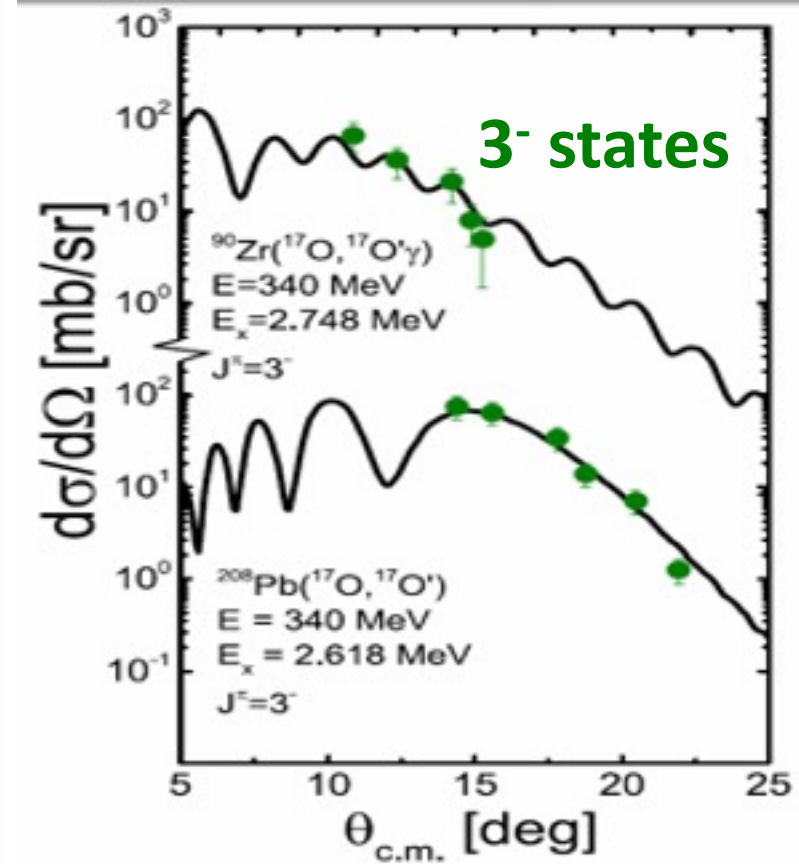
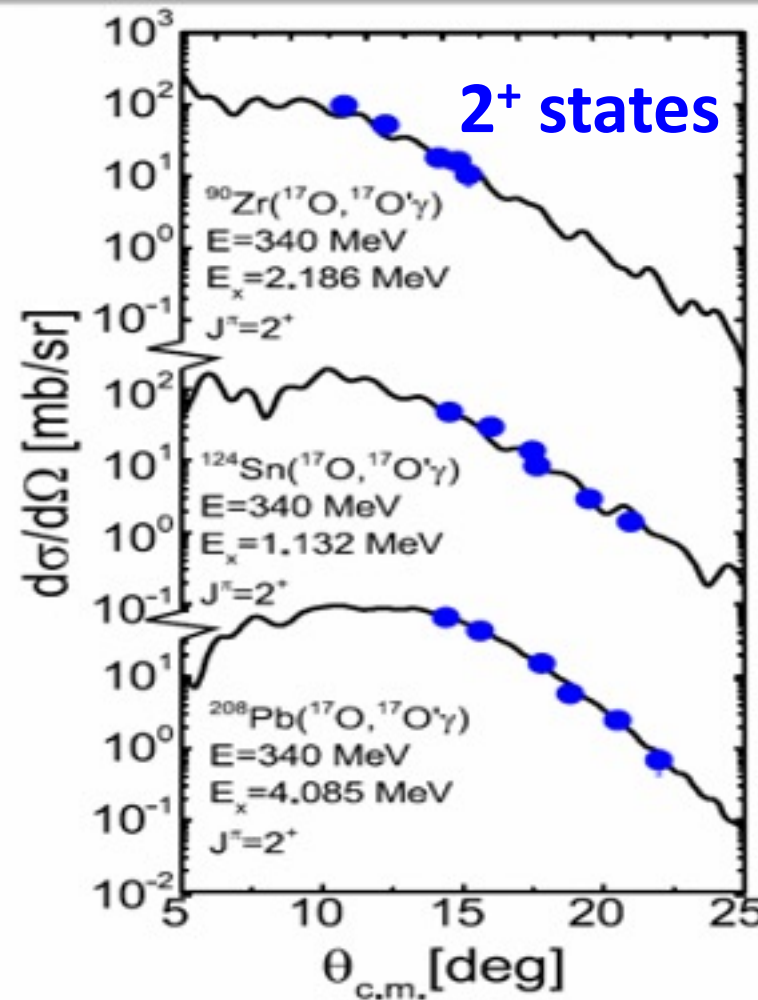
$$M_n / M_p = N/Z$$

- (e,e') and (γ,γ') experiments, see e.g.: <http://www.nndc.bnl.gov/ensdf/>

**for the case of ^{208}Pb :

D.J. Horen et al. PRC44(1991)128

Agreement with measurements at similar beam energy**

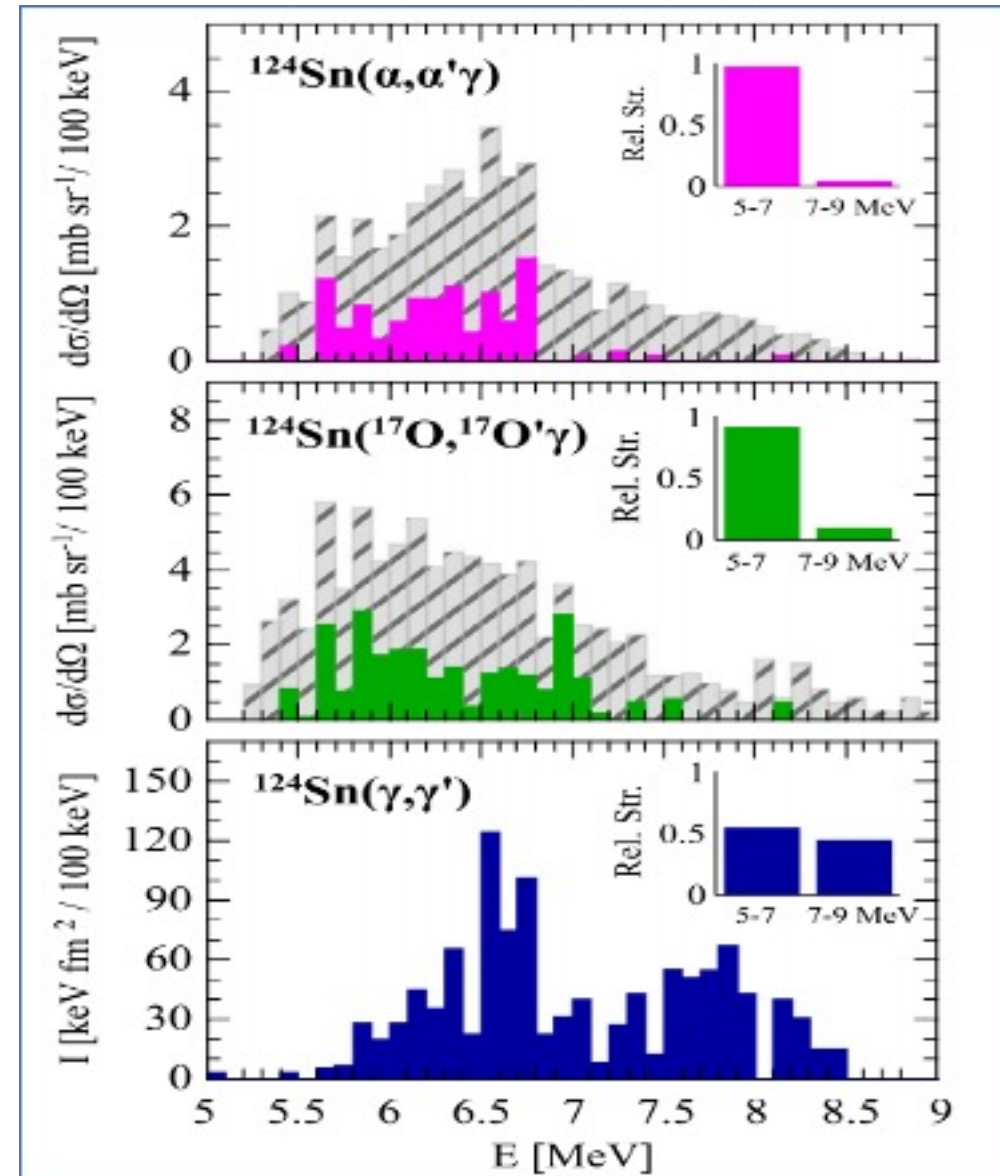


Results for ^{124}Sn and comparison with other probes

1^- excitation with ^{17}O
20 MeV/u

Comparison with alpha and
Gamma scattering

In all cases only the region below
7 MeV is populated by all
reactions



(**) J. Endres et al., PRL105, 212503 (2010)

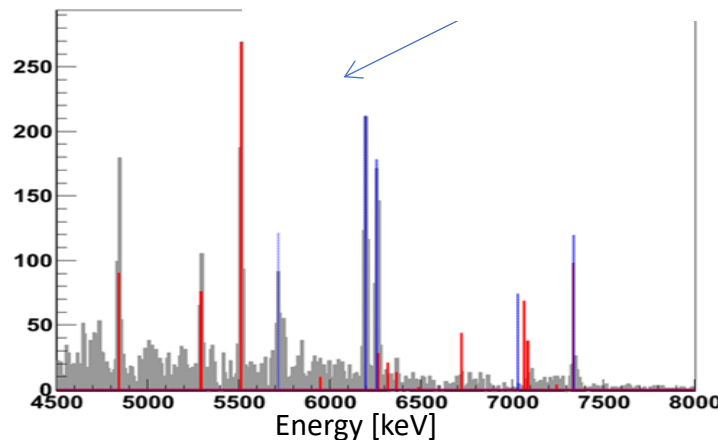
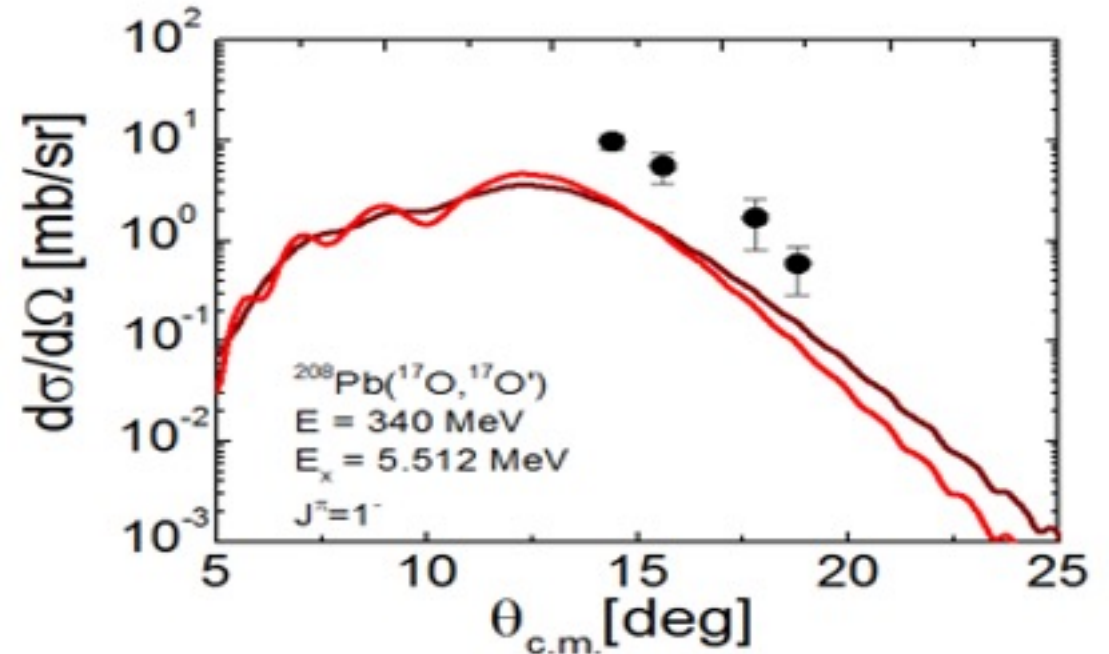
L. Pellegrini, et al., PLB738 (2014)519

1- states in ^{208}Pb

The calculation accounts only for a fraction of the measured yield

Why?

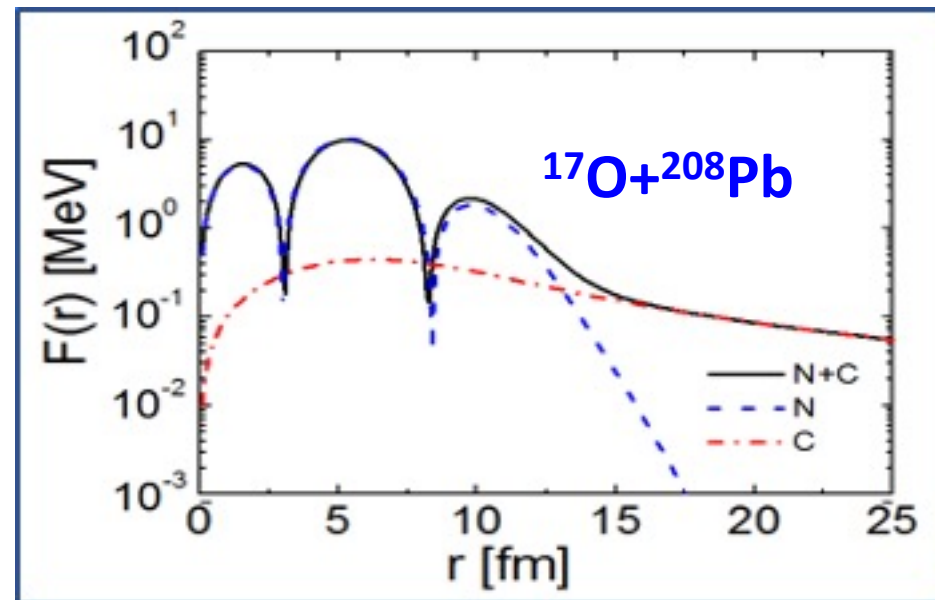
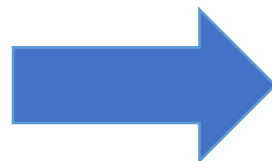
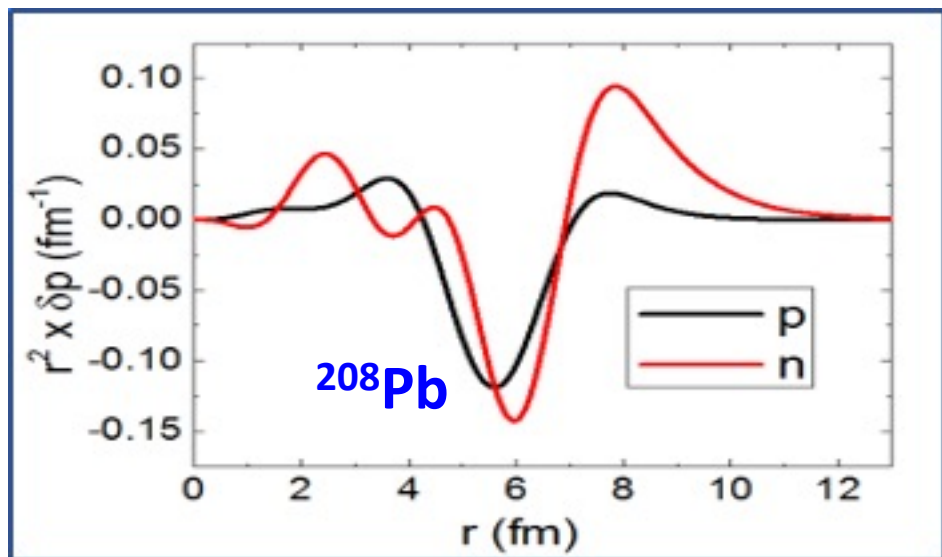
Calculations obtained using a standard form factor are found to be very similar to the Coulomb excitation alone



$B(E1) \uparrow$ values
known from $(\gamma, \gamma)^{**}$

** photon scattering experiments:
N. Ryezayeva
et. al PRL89(2002)272502,
T. Shizuma et al. PRC78(2008)061303

Use a microscopic form factor



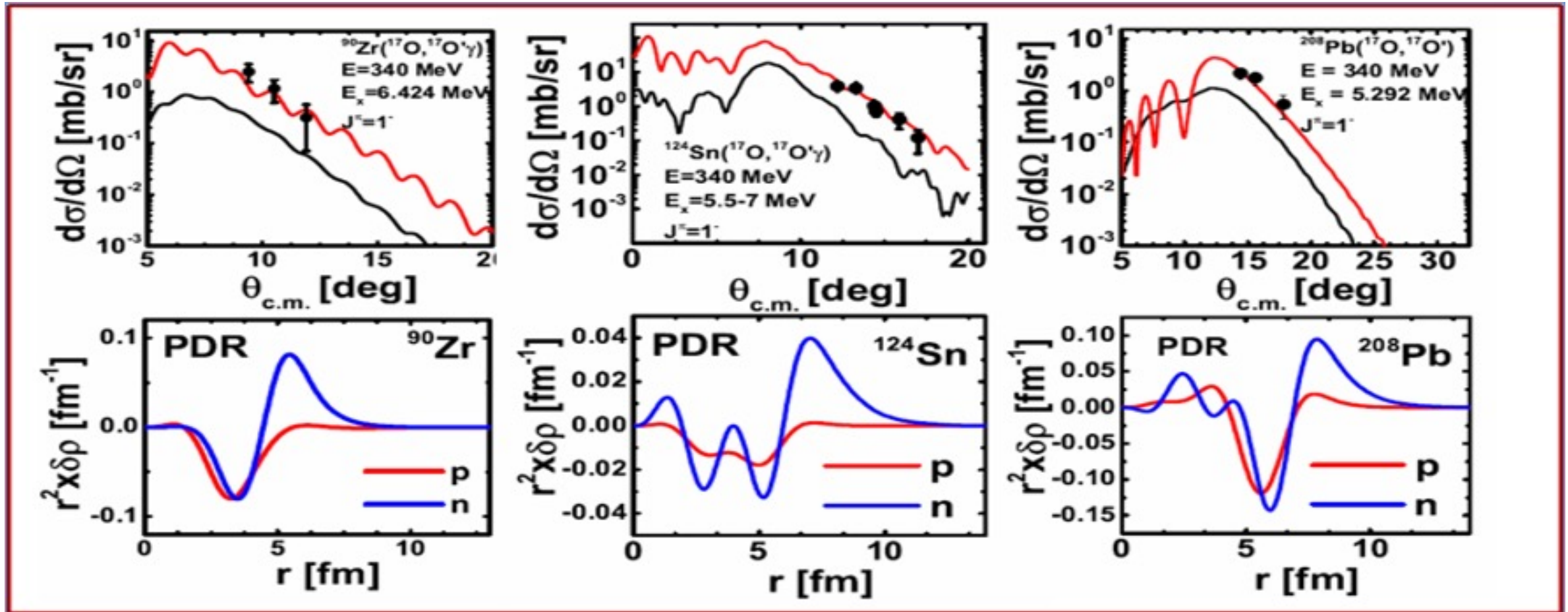
Scattering of ^{17}O at these energies is probing mainly the nuclear surface!!

*E. G. Lanza et al., Phys. Rev. C 79 (2009) 054615

**E. G. Lanza et al., Phys. Rev. C 84 (2011) 064602

***E. G. Lanza et al., PRC 89 (2014) 041601

Some results for selected 1^- in ^{90}Zr , ^{124}Sn and ^{208}Pb



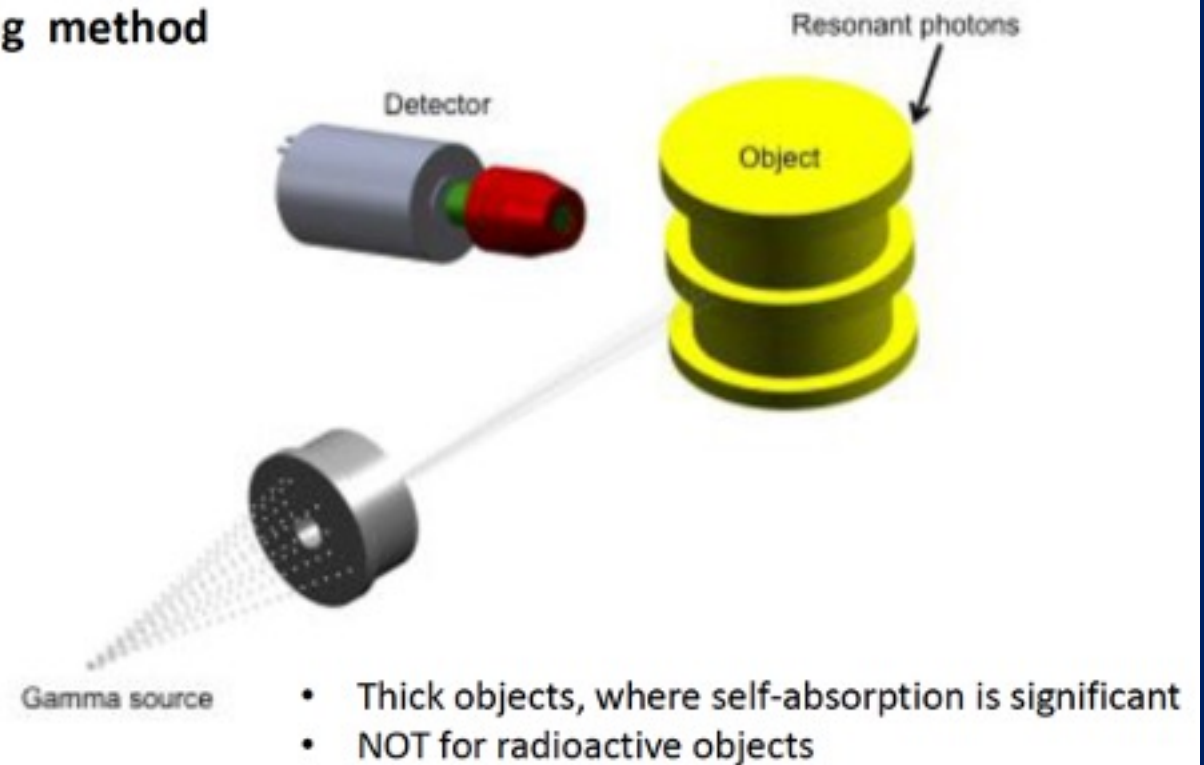
Calculated transition densities:

*(for ^{124}Sn) E. Litvinova, et al., PRC 78 (2008) 014312, **E.G. Lanza, et al., PRC 89 (2014) 041601

Industrial tomography

material inspection nuclear waste- food contaminations

Scattering method



Conclusions and perspectives

- GMR and isoscalar GDR information on nuclear compressibility
- GDR and Pygmy at $T=0$ Information on the neutron Skin- related to the Radius of the neutron stars- Information for the r-process
more work on deformed nuclei and on exotic nuclei with radioactive beams
- Finite temperature : extreme deformations
Isospin mixing measurements in other mass regions