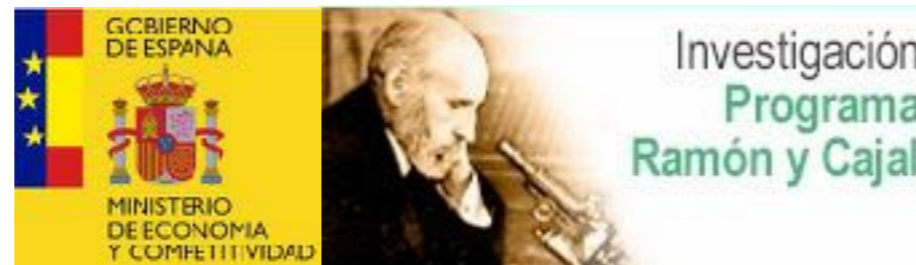


Collective and single-particle motion in beyond-mean-field approaches

Tomás R. Rodríguez

Nuclear Structure and Astrophysical Applications

Milano, September 20th, 2017



Acknowledgments



Introduction

SCCM with cranking

Collective and single-particle states in ^{44}S

Summary and Outlook

M. Borrajo (UAM-Madrid)

J. L. Egido (UAM-Madrid)

Acknowledgments



Introduction

SCCM with cranking

Collective and single-particle states in ^{44}S

Summary and Outlook

Congratulations, Xavier!

Acknowledgments

Congratulations, Xavier!



Acknowledgments

Congratulations, Xavier!



1. Introduction

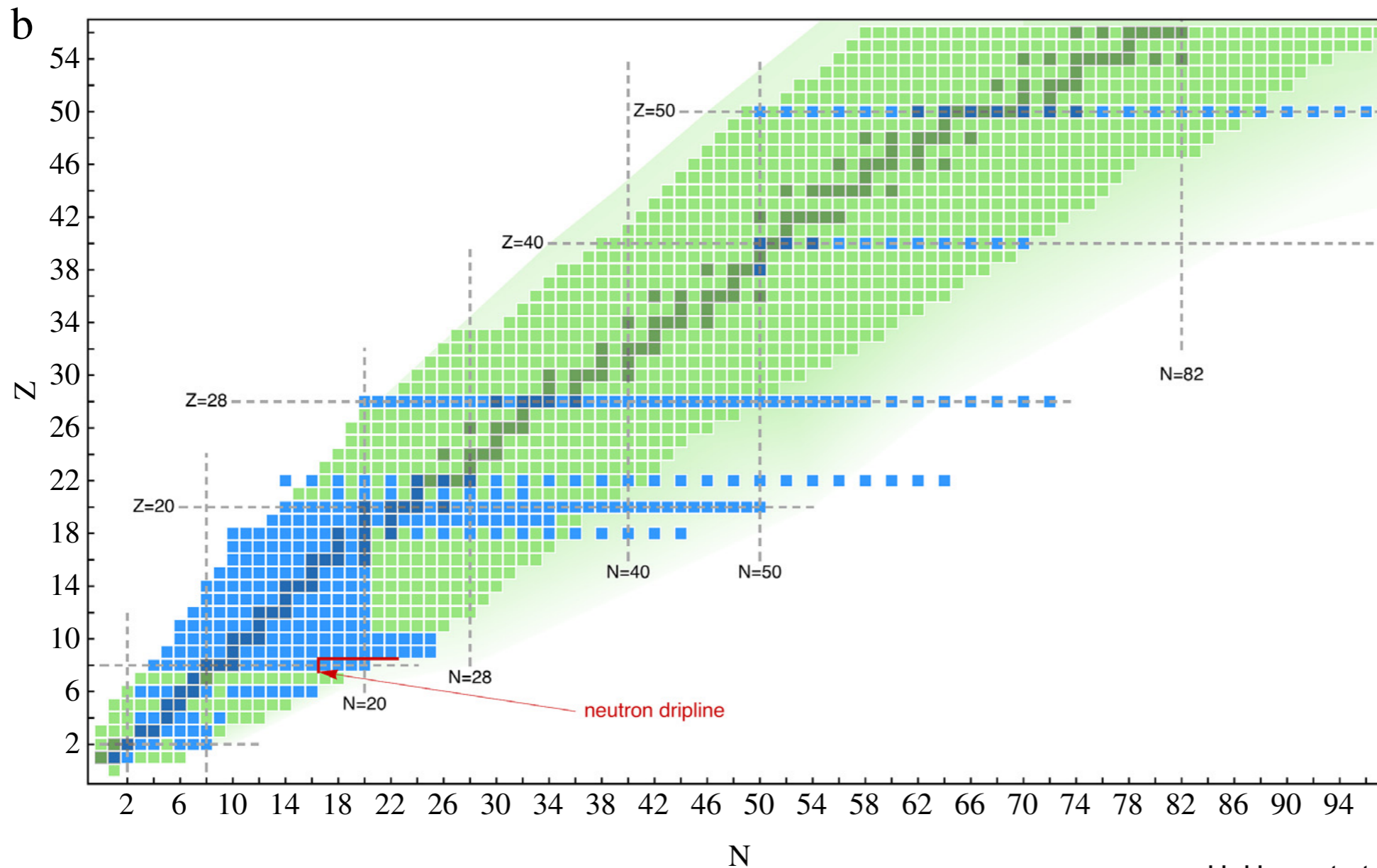
2. SCCM with cranking

3. Collective and Single-particle states in ^{44}S

4. Summary and Outlook

Ab initio world

► First-principle calculations based on EFT+SRG interactions and different many-body methods can be nowadays performed in a large amount of nuclei.

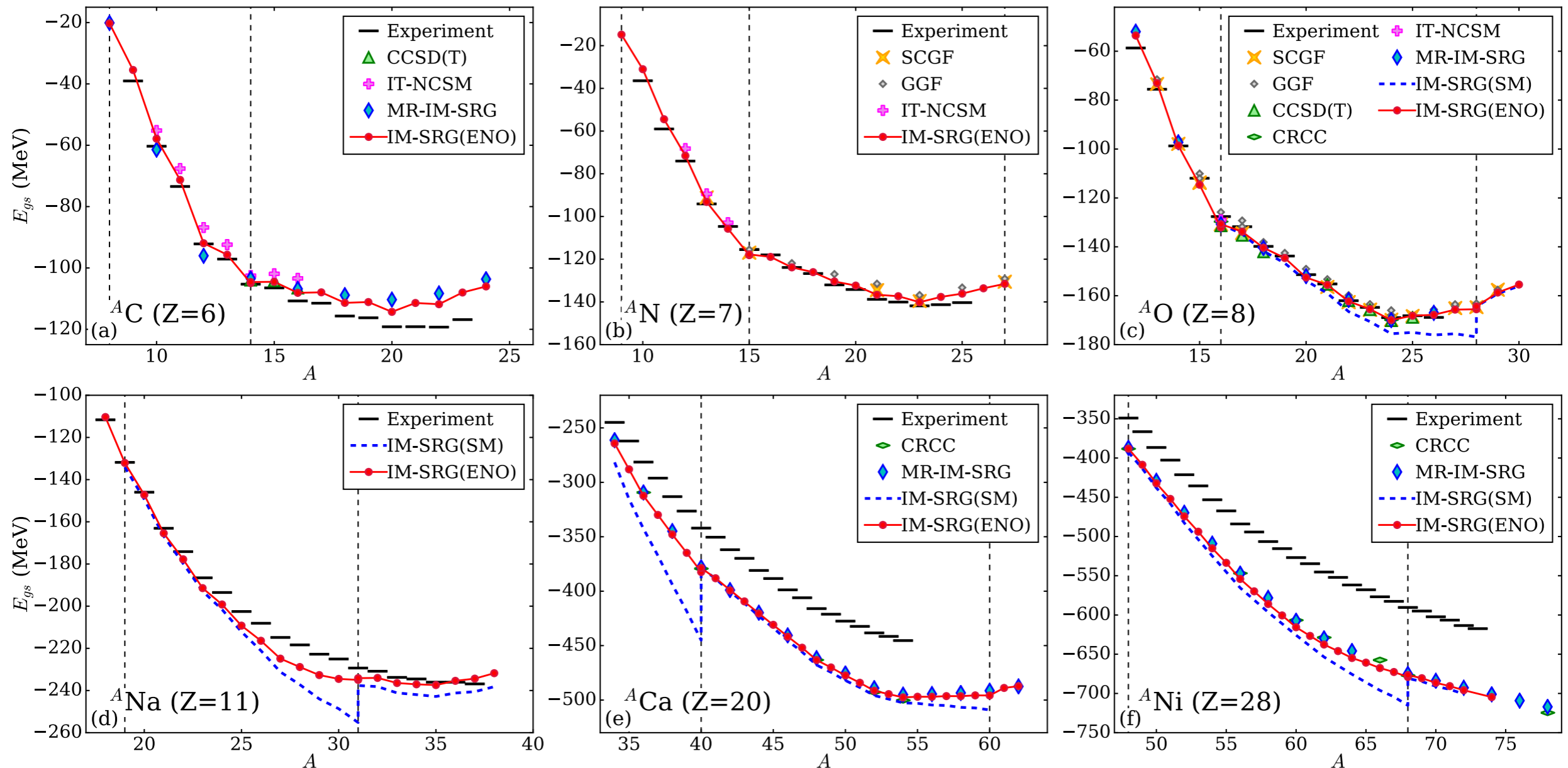


philosopher's stone

H. Hergert et al., Phys. Rep. 621, 165 (2016)

Ab initio world

► Still some problems in reproducing some basic observables in heavier systems.



S. R. Stroberg et al., Phys. Rev. Lett. 118, 032502 (2017)

Phenomenological world



Introduction

SCCM with cranking

Collective and single-particle states in ^{44}S

Summary and Outlook

- ▶ Use of phenomenological interactions (adjusted to data in finite nuclei) is necessary to obtain precise predictions/descriptions of ground state, spectroscopic and reaction data.

► Use of phenomenological interactions (adjusted to data in finite nuclei) is necessary to obtain precise predictions/descriptions of ground state, spectroscopic and reaction data.

LARGE SCALE SHELL MODEL

- Exact diagonalizations within a valence space.
- Effective interactions adapted to the valence space and adjusted to reproduce the evolution of single particle energies (monopoles).
- Very precise description of spectroscopy and transitions of nuclei.
- Limited by the combinatorial increase of the number of configurations.
- Defined in the laboratory frame → Intrinsic shapes can only be inferred from the spectra and electromagnetic moments

► Use of phenomenological interactions (adjusted to data in finite nuclei) is necessary to obtain precise predictions/descriptions of ground state, spectroscopic and reaction data.

SELF-CONSISTENT MEAN FIELD

- Variational approach with simple trial wave functions (HFB) using 'universal' functionals (applicable to the whole nuclear chart).
- Parameters of the functional fitted to bulk properties and masses and radii of finite nuclei.
- Very precise description of ground state properties and collective phenomena.
- Defined in the intrinsic frame
- **Spectroscopy with beyond mean-field techniques (GCM, QRPA, ...)**

LARGE SCALE SHELL MODEL

- Exact diagonalizations within a valence space.
- Effective interactions adapted to the valence space and adjusted to reproduce the evolution of single particle energies (monopoles).
- Very precise description of spectroscopy and transitions of nuclei.
- Limited by the combinatorial increase of the number of configurations.
- Defined in the laboratory frame → Intrinsic shapes can only be inferred from the spectra and electromagnetic moments

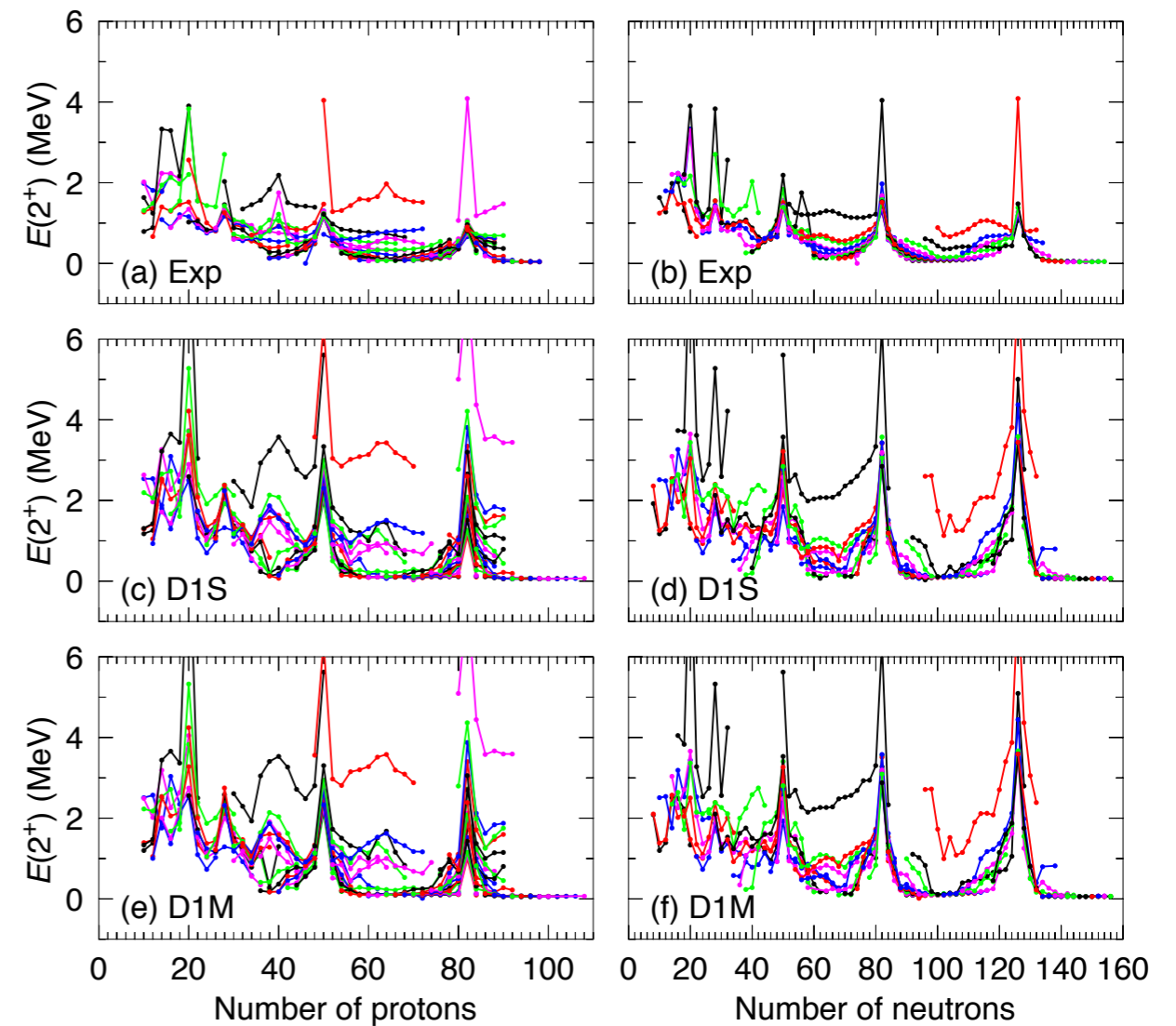
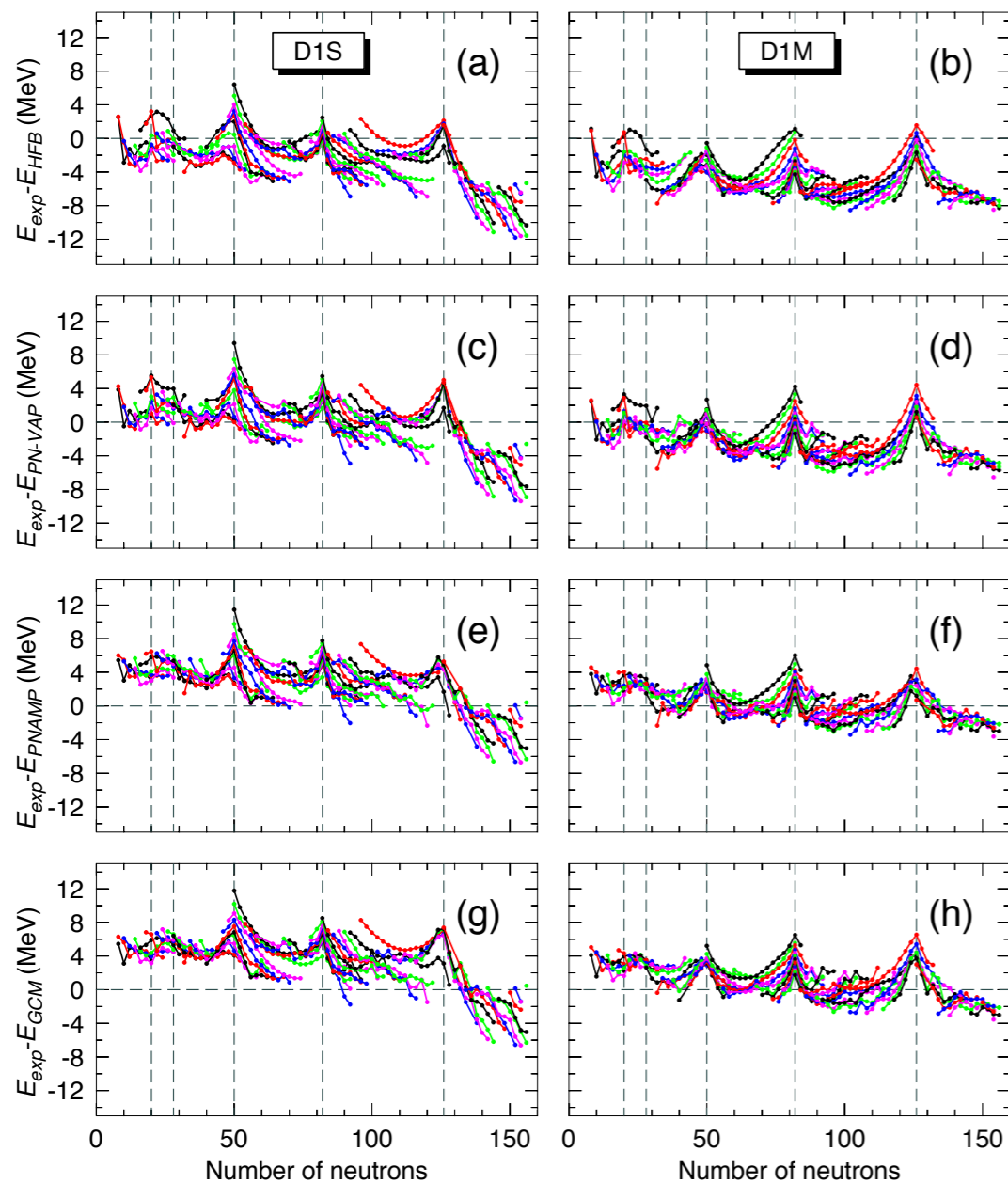
► Use of phenomenological interactions (adjusted to data in finite nuclei) is necessary to obtain precise predictions/descriptions of ground state, spectroscopic and reaction data.

SELF-CONSISTENT MEAN FIELD

- Variational approach with simple trial wave functions (HFB) using 'universal' functionals (applicable to the whole nuclear chart).
- Parameters of the functional fitted to bulk properties and masses and radii of finite nuclei.
- Very precise description of ground state properties and collective phenomena.
- Defined in the intrinsic frame
- **Spectroscopy with beyond mean-field techniques (GCM, QRPA, ...)**

Example of the performance of Gogny D1S/D1M interactions

Global study of nuclear masses and 2^+ excitation energies



Physical Review C 91, 044315 (2015)

Gogny interaction

Effective nucleon-nucleon interaction: Gogny force (DIS/D1M)

$$V(1,2) = \sum_{i=1}^2 e^{-(\vec{r}_1 - \vec{r}_2)^2 / \mu_i^2} (W_i + B_i P^\sigma - H_i P^\tau - M_i P^\sigma P^\tau) \\ + iW_0(\sigma_1 + \sigma_2) \vec{k} \times \delta(\vec{r}_1 - \vec{r}_2) \vec{k} + V_{\text{Coulomb}}(\vec{r}_1, \vec{r}_2) \\ + t_3(1 + x_0 P^\sigma) \delta(\vec{r}_1 - \vec{r}_2) \rho^\alpha ((\vec{r}_1 + \vec{r}_2)/2)$$



Gogny interaction

Effective nucleon-nucleon interaction: Gogny force (DIS/D1M)

$$V(1,2) = \sum_{i=1}^2 e^{-(\vec{r}_1 - \vec{r}_2)^2 / \mu_i^2} (W_i + B_i P^\sigma - H_i P^\tau - M_i P^\sigma P^\tau) \\ + iW_0 (\sigma_1 + \sigma_2) \vec{k} \times \delta(\vec{r}_1 - \vec{r}_2) \vec{k} + V_{\text{Coulomb}}(\vec{r}_1, \vec{r}_2)$$

2-body potential

$$+ t_3 (1 + x_0 P^\sigma) \delta(\vec{r}_1 - \vec{r}_2) \rho^\alpha ((\vec{r}_1 + \vec{r}_2)/2)$$



Gogny interaction

Effective nucleon-nucleon interaction: Gogny force (DIS/D1M)

$$V(1,2) = \sum_{i=1}^2 e^{-(\vec{r}_1 - \vec{r}_2)^2 / \mu_i^2} (W_i + B_i P^\sigma - H_i P^\tau - M_i P^\sigma P^\tau) \\ + iW_0 (\sigma_1 + \sigma_2) \vec{k} \times \delta(\vec{r}_1 - \vec{r}_2) \vec{k} + V_{\text{Coulomb}}(\vec{r}_1, \vec{r}_2)$$

2-body potential

$$+ t_3 (1 + x_0 P^\sigma) \delta(\vec{r}_1 - \vec{r}_2) \rho^\alpha ((\vec{r}_1 + \vec{r}_2)/2)$$

Density dependent term



Gogny interaction

Effective nucleon-nucleon interaction: Gogny force (DIS/D1M)

$$V(1,2) = \sum_{i=1}^2 e^{-(\vec{r}_1 - \vec{r}_2)^2 / \mu_i^2} (W_i + B_i P^\sigma - H_i P^\tau - M_i P^\sigma P^\tau) \\ + iW_0 (\sigma_1 + \sigma_2) \vec{k} \times \delta(\vec{r}_1 - \vec{r}_2) \vec{k} + V_{\text{Coulomb}}(\vec{r}_1, \vec{r}_2)$$

2-body potential

$$+ t_3 (1 + x_0 P^\sigma) \delta(\vec{r}_1 - \vec{r}_2) \rho^\alpha ((\vec{r}_1 + \vec{r}_2)/2)$$

Density dependent term



Other alternatives: Skyrme, relativistic Lagrangians, BCPM, ...

Gogny EDF with cranked states

- *Initial intrinsic states: PN-VAP*

$$E^{N,Z}[\Phi] = \frac{\langle \Phi | \hat{H}_{2b} \hat{P}^N \hat{P}^Z | \Phi \rangle}{\langle \Phi | \hat{P}^N \hat{P}^Z | \Phi \rangle} + \varepsilon_{DD}^{N,Z}(\Phi) - \lambda_{q_{20}} \langle \Phi | \hat{Q}_{20} | \Phi \rangle - \lambda_{q_{22}} \langle \Phi | \hat{Q}_{22} | \Phi \rangle - \omega \langle \Phi | \hat{J}_x | \Phi \rangle$$

Gogny EDF with cranked states

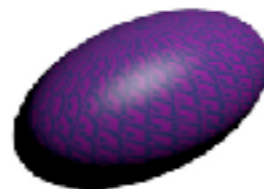
- Initial intrinsic states: PN-VAP

$$E^{N,Z}[\Phi] = \frac{\langle \Phi | \hat{H}_{2b} \hat{P}^N \hat{P}^Z | \Phi \rangle}{\langle \Phi | \hat{P}^N \hat{P}^Z | \Phi \rangle} + \varepsilon_{DD}^{N,Z}(\Phi) - \lambda_{q_{20}} \langle \Phi | \hat{Q}_{20} | \Phi \rangle - \lambda_{q_{22}} \langle \Phi | \hat{Q}_{22} | \Phi \rangle - \omega \langle \Phi | \hat{J}_x | \Phi \rangle$$

$I_c=0$



$I_c=2$



$I_c=4$



Gogny EDF with cranked states

- Initial intrinsic states: PN-VAP

$$E^{N,Z}[\Phi] = \frac{\langle \Phi | \hat{H}_{2b} \hat{P}^N \hat{P}^Z | \Phi \rangle}{\langle \Phi | \hat{P}^N \hat{P}^Z | \Phi \rangle} + \varepsilon_{DD}^{N,Z}(\Phi) - \lambda_{q20} \langle \Phi | \hat{Q}_{20} | \Phi \rangle - \lambda_{q22} \langle \Phi | \hat{Q}_{22} | \Phi \rangle - \omega \langle \Phi | \hat{J}_x | \Phi \rangle$$

- Intermediate Particle Number and Angular Momentum Projected states

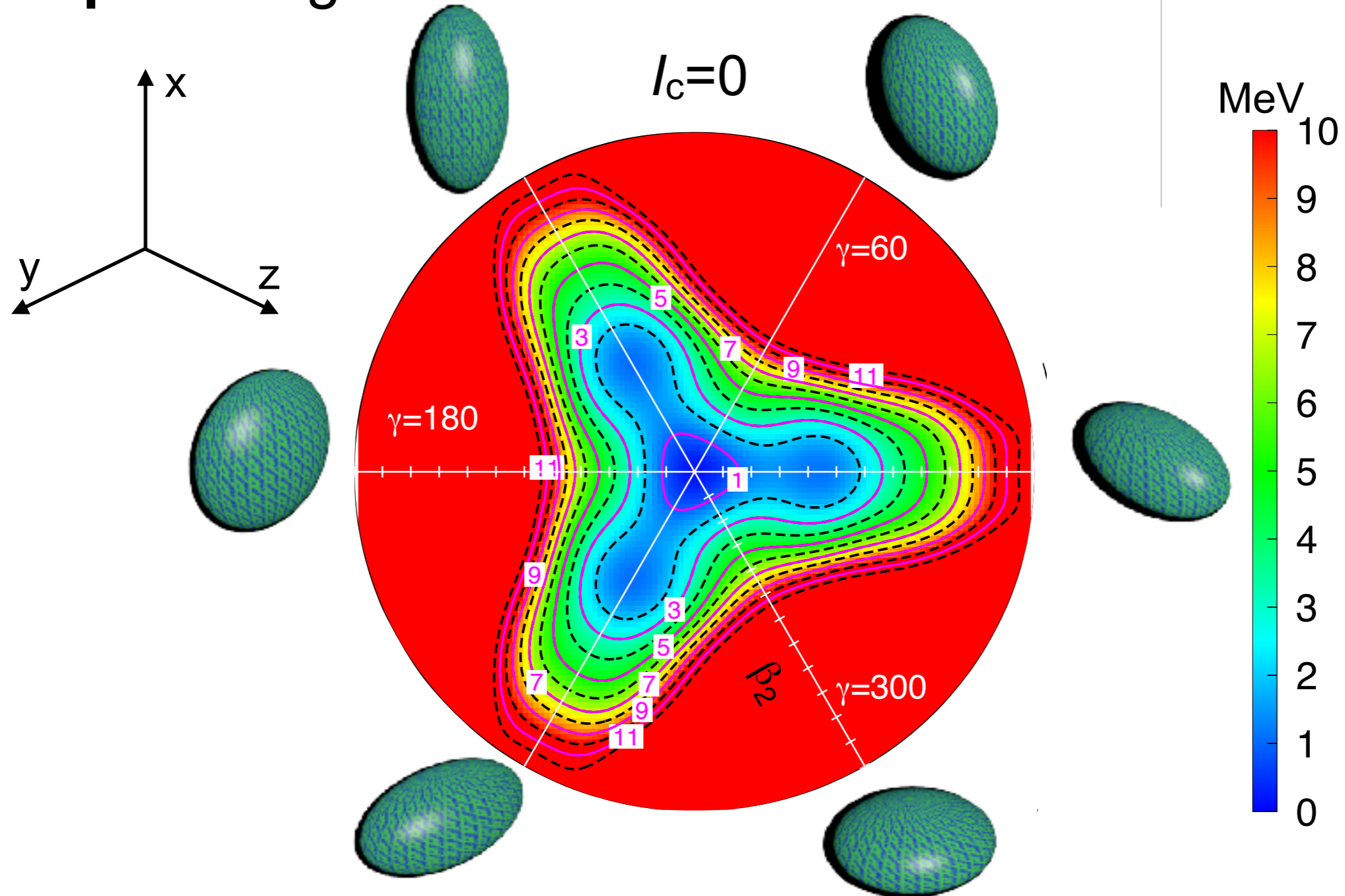
$$|IMK; NZ; \beta\gamma; \omega\rangle = \frac{2I+1}{8\pi^2} \int \mathcal{D}_{KK'}^{I*}(\Omega) \hat{R}(\Omega) \hat{P}^N \hat{P}^Z | \Phi(\beta, \gamma, \omega) \rangle d\Omega$$

- Final GCM states $|IM; NZ; \sigma\rangle = \sum_{K\beta\gamma\omega} f_{K\beta\gamma\omega}^{I;NZ;\sigma} |IMK; NZ; \beta\gamma; \omega\rangle$

$$\sum_{K'\beta'\gamma'\omega'} \left(\mathcal{H}_{K\beta\gamma\omega; K'\beta'\gamma'\omega'}^{I;NZ} - E^{I;NZ;\sigma} \mathcal{N}_{K\beta\gamma\omega; K'\beta'\gamma'\omega'}^{I;NZ} \right) f_{K'\beta'\gamma'\omega'}^{I;NZ;\sigma} = 0$$

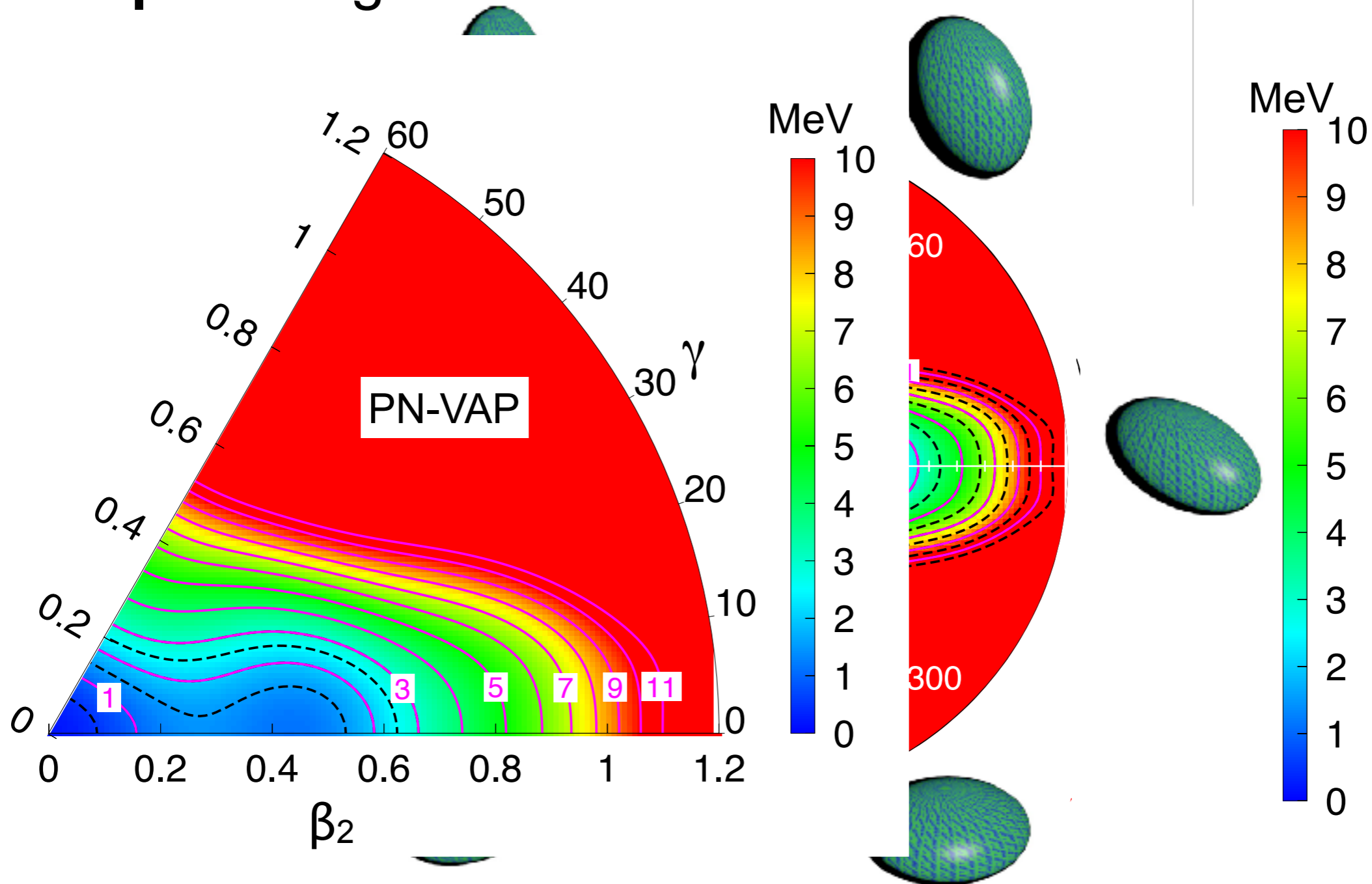
Gogny EDF with cranked states

Example: ^{32}Mg triaxial+TRSB



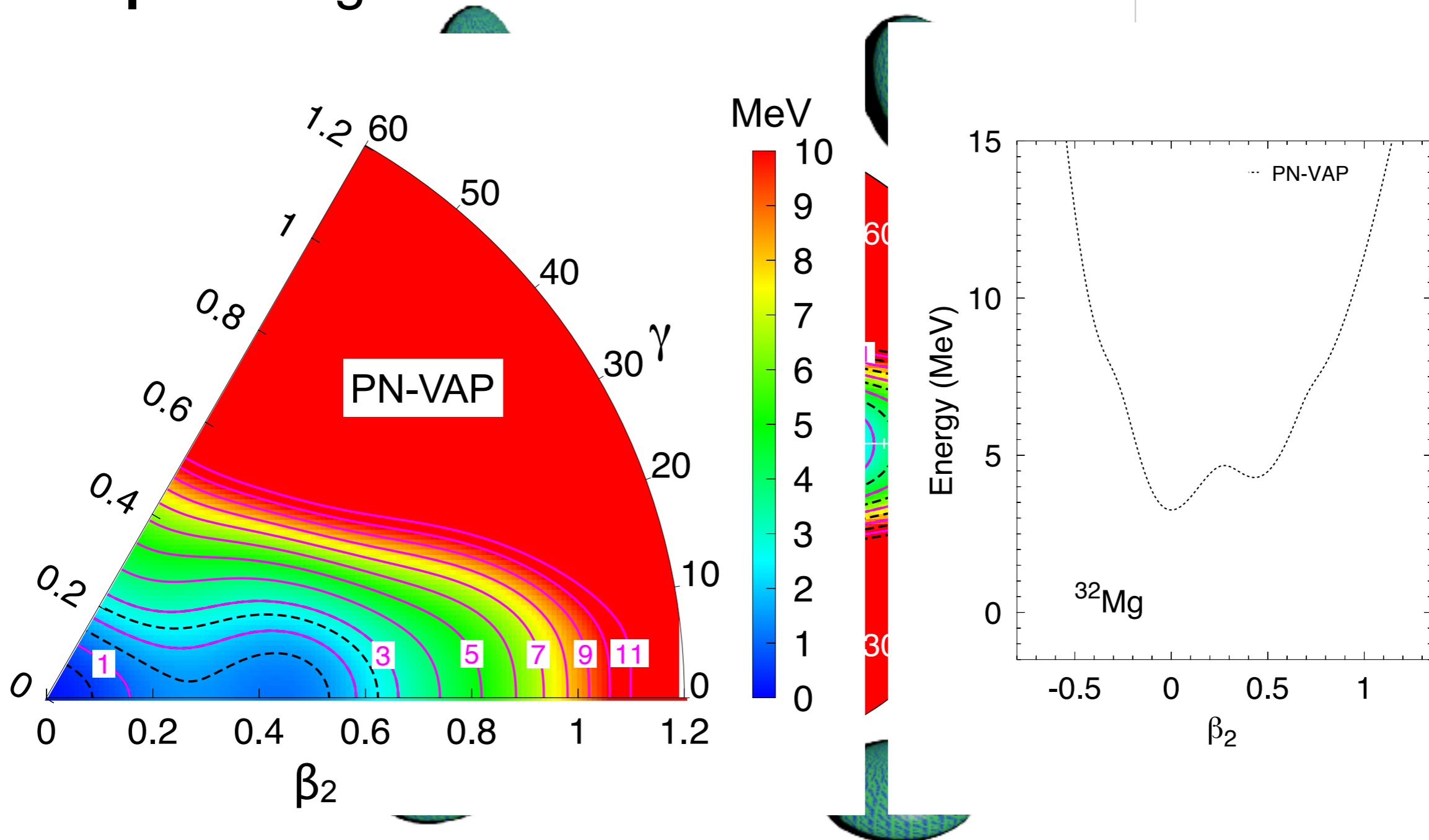
Gogny EDF with cranked states

Example: ^{32}Mg triaxial+TRSB



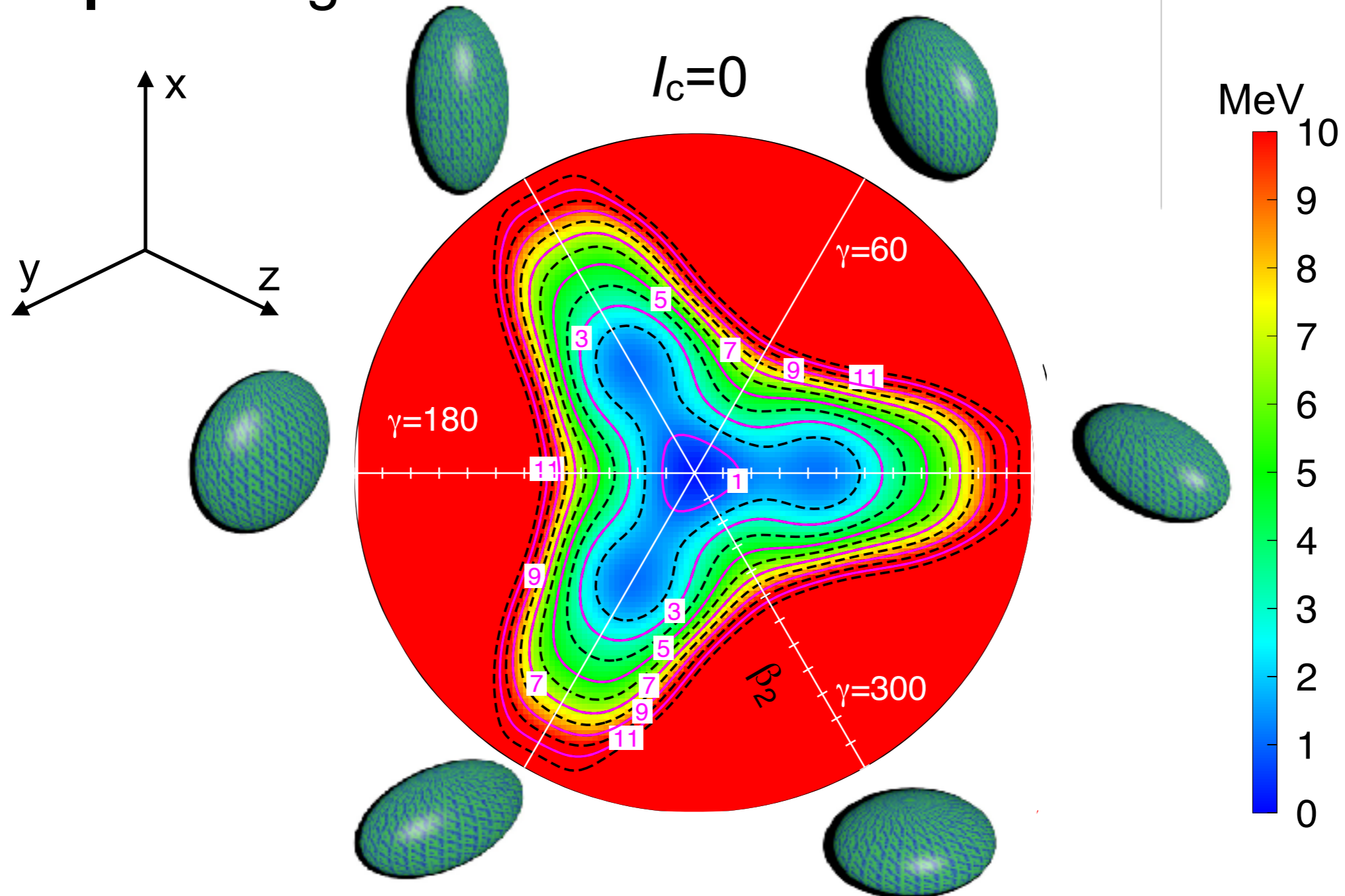
Gogny EDF with cranked states

Example: ^{32}Mg triaxial+TRSB



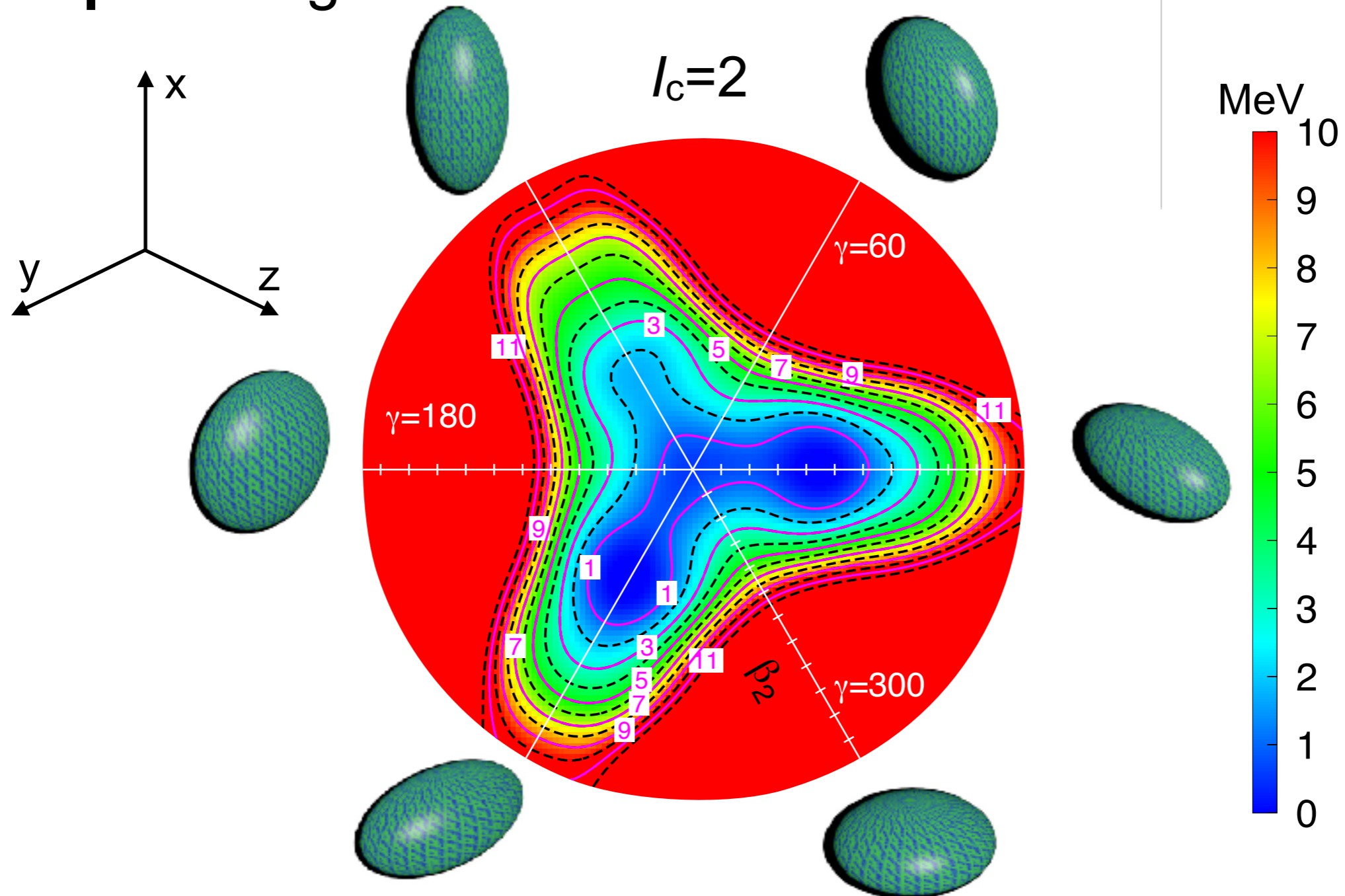
Gogny EDF with cranked states

Example: ^{32}Mg triaxial+TRSB



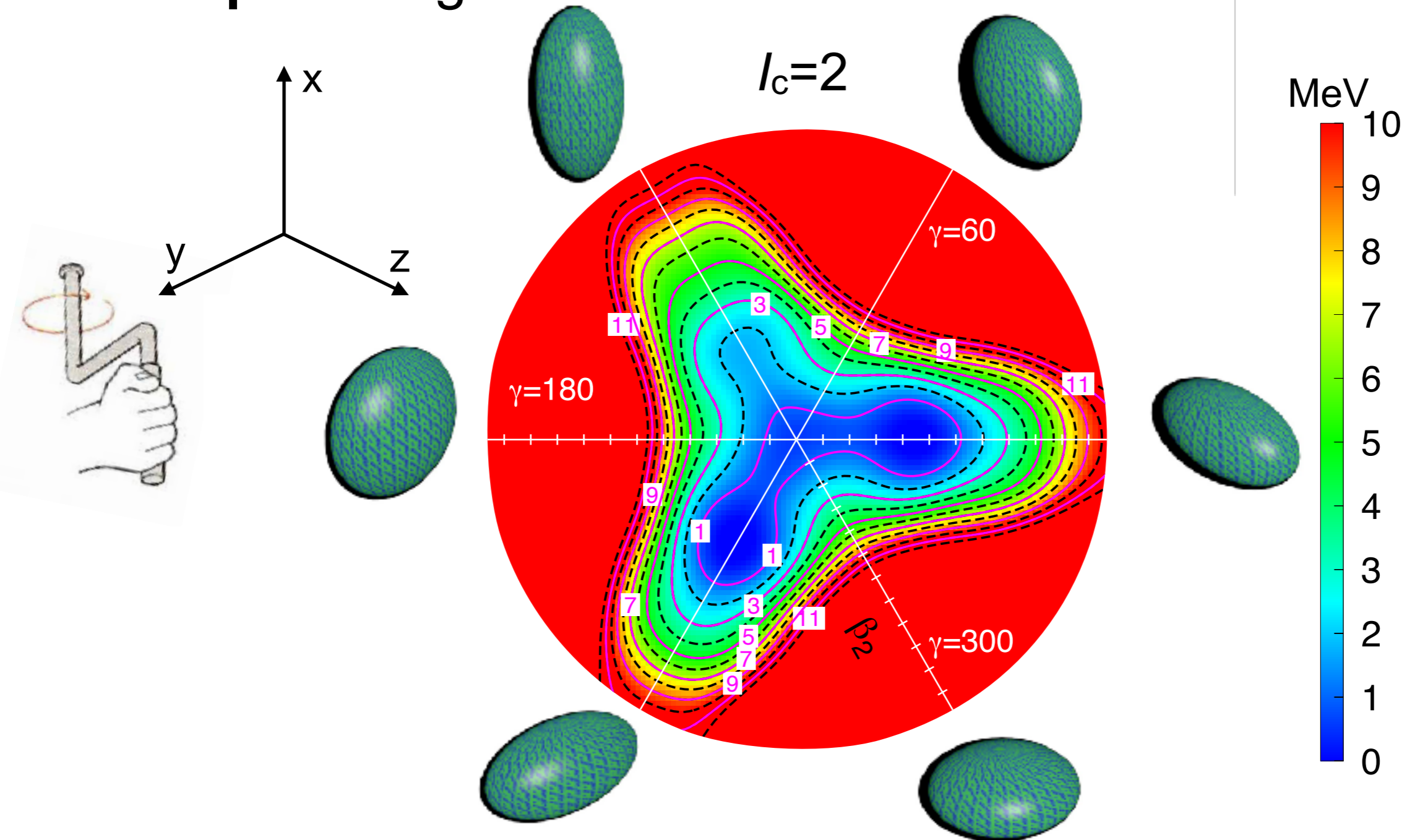
Gogny EDF with cranked states

Example: ^{32}Mg triaxial+TRSB



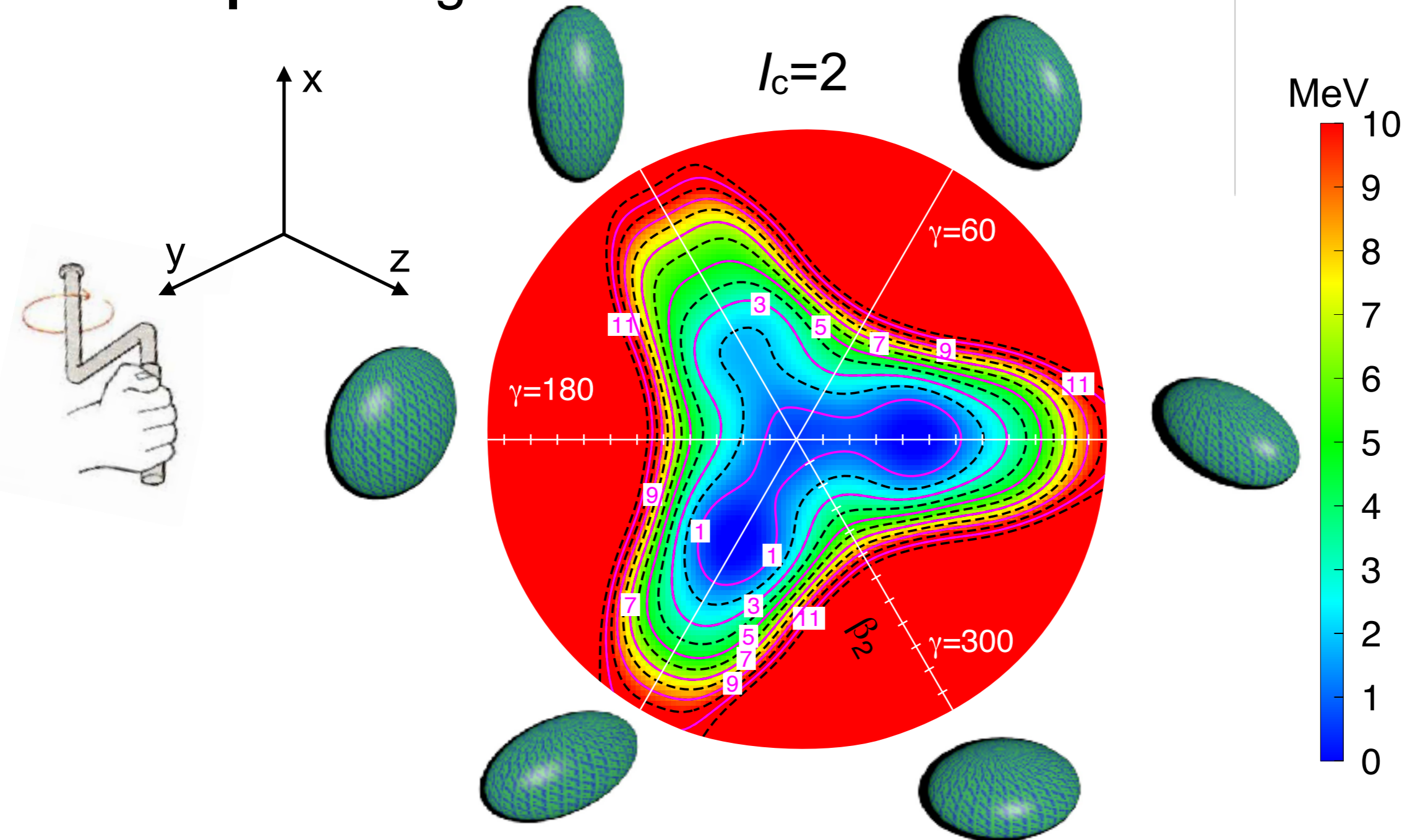
Gogny EDF with cranked states

Example: ^{32}Mg triaxial+TRSB



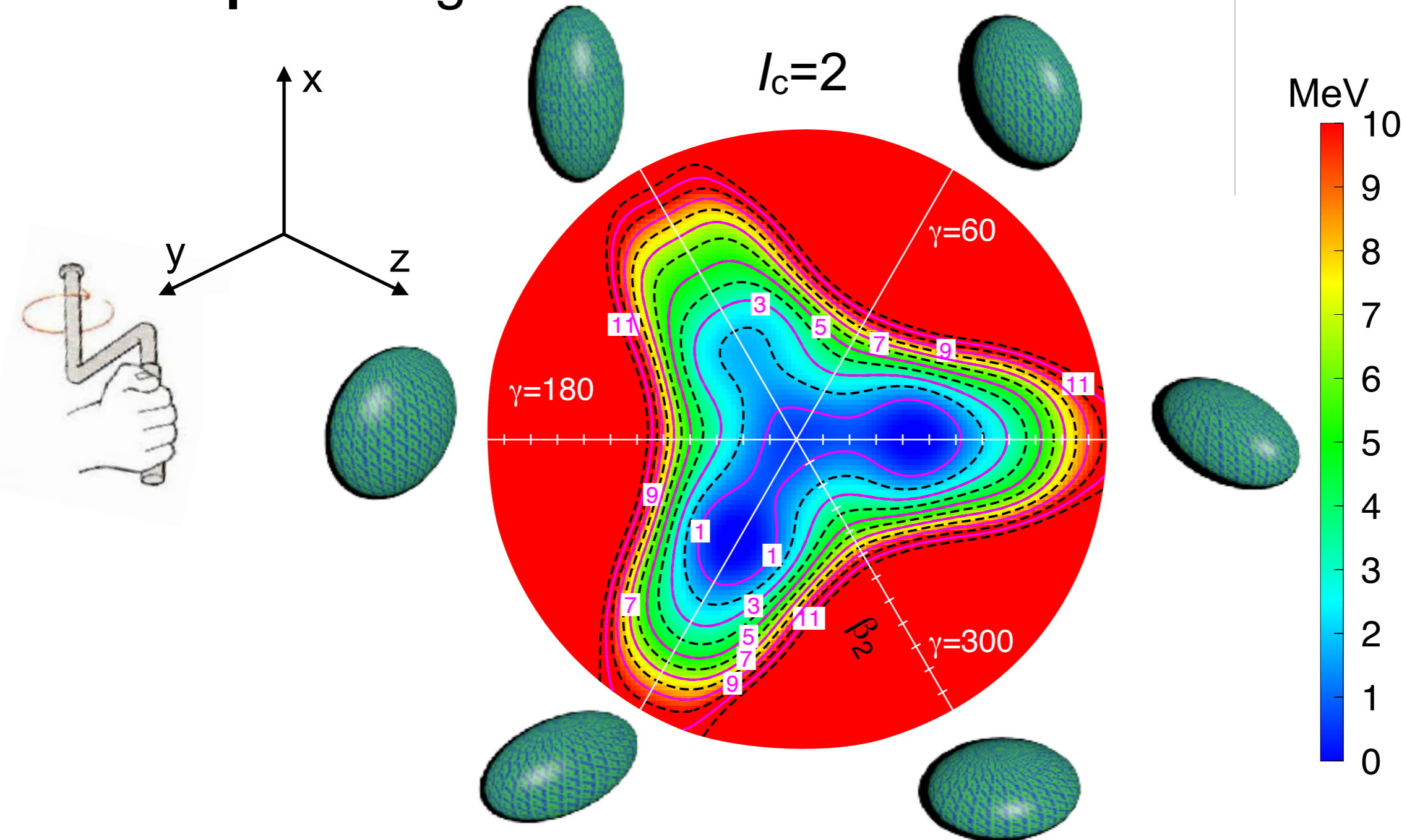
Gogny EDF with cranked states

Example: ^{32}Mg triaxial+TRSB



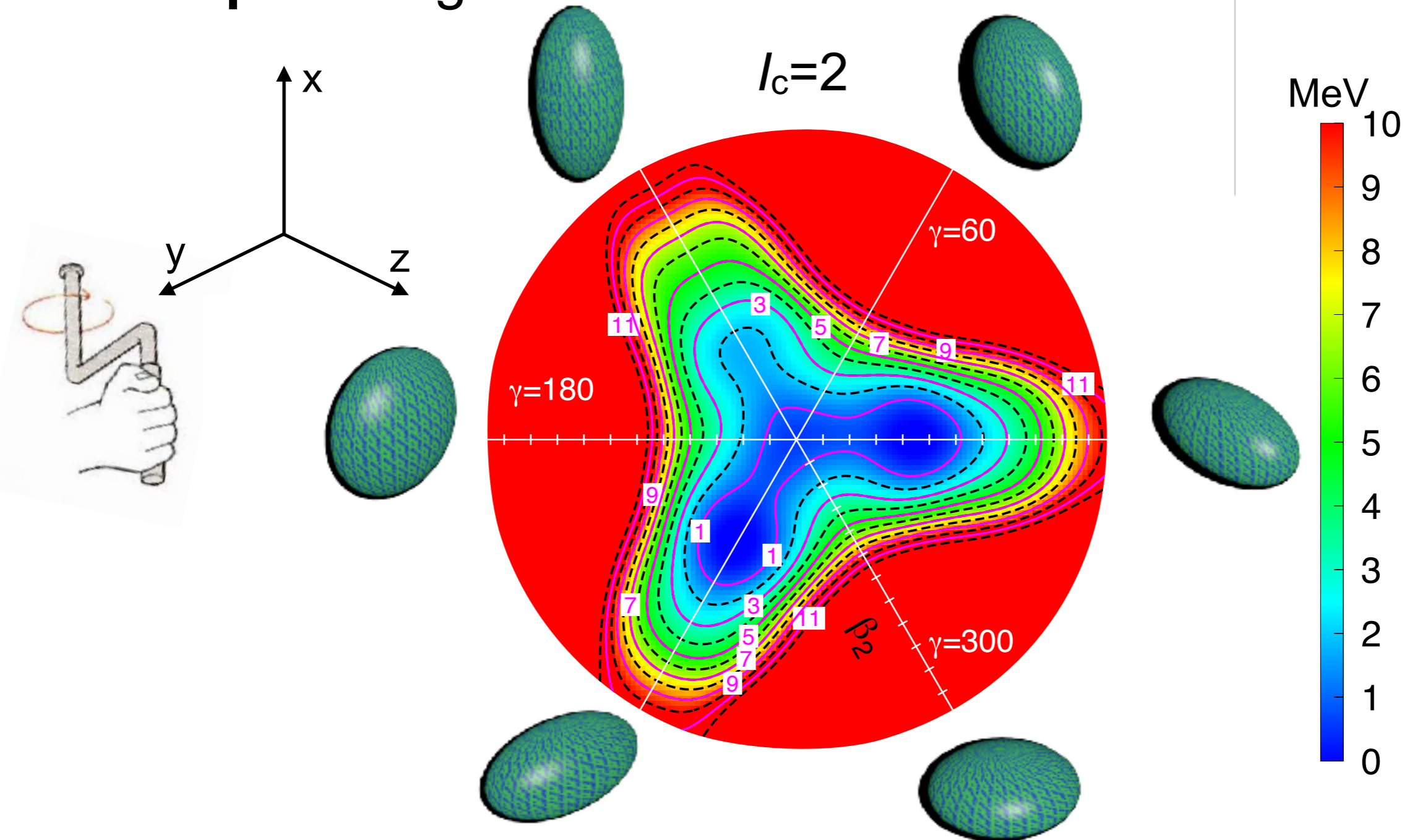
Gogny EDF with cranked states

Example: ^{32}Mg triaxial+TRSB



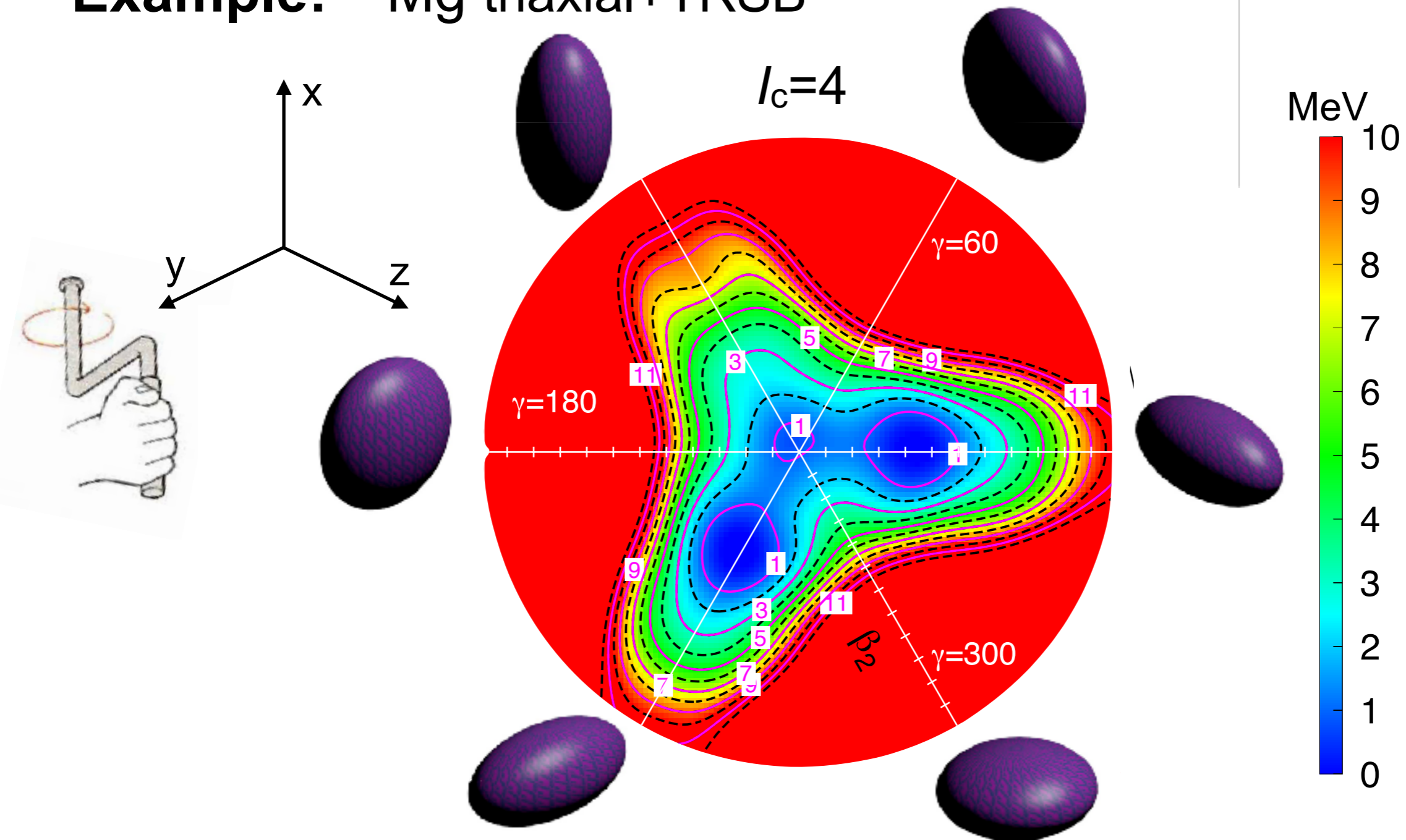
Gogny EDF with cranked states

Example: ^{32}Mg triaxial+TRSB



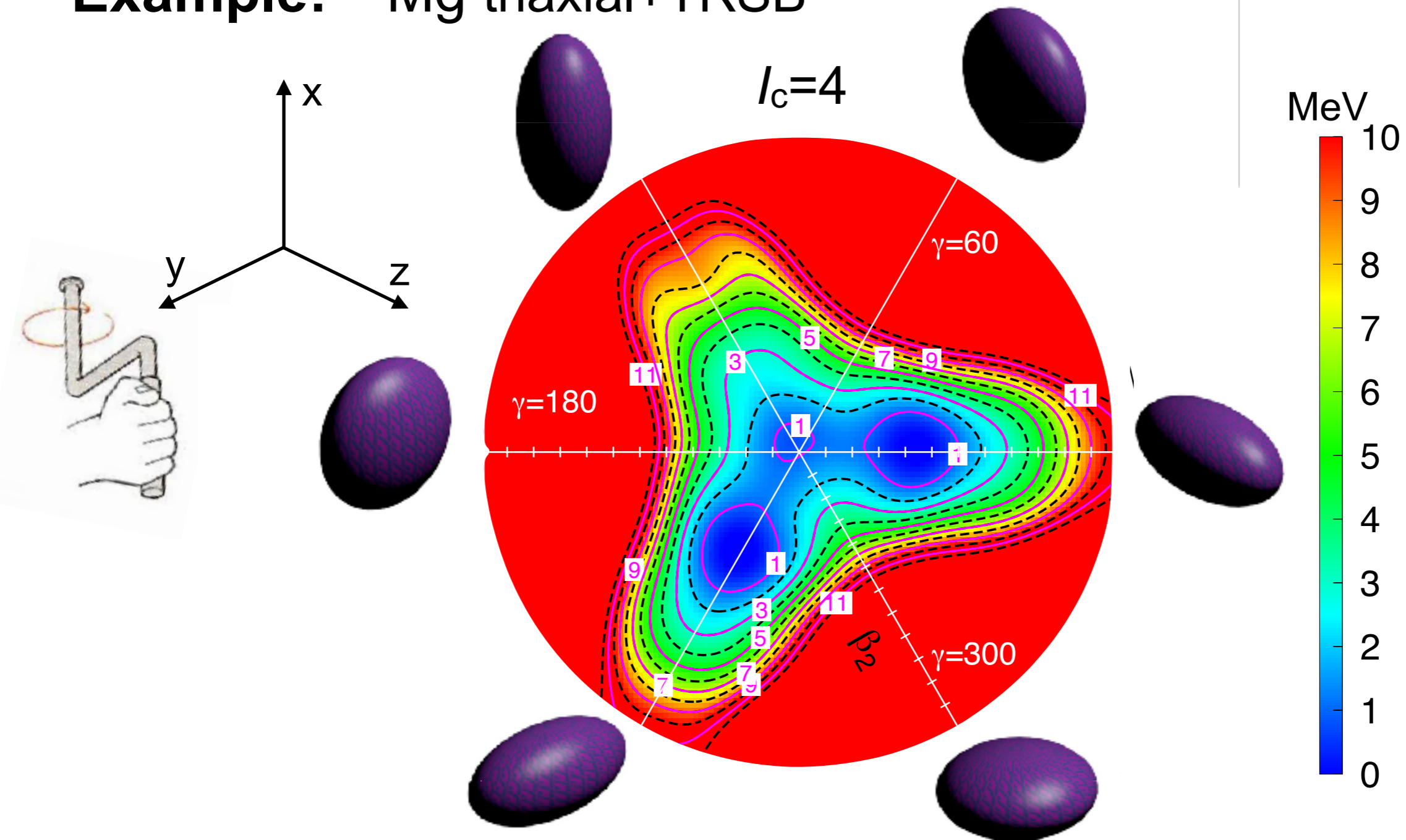
Gogny EDF with cranked states

Example: ^{32}Mg triaxial+TRSB



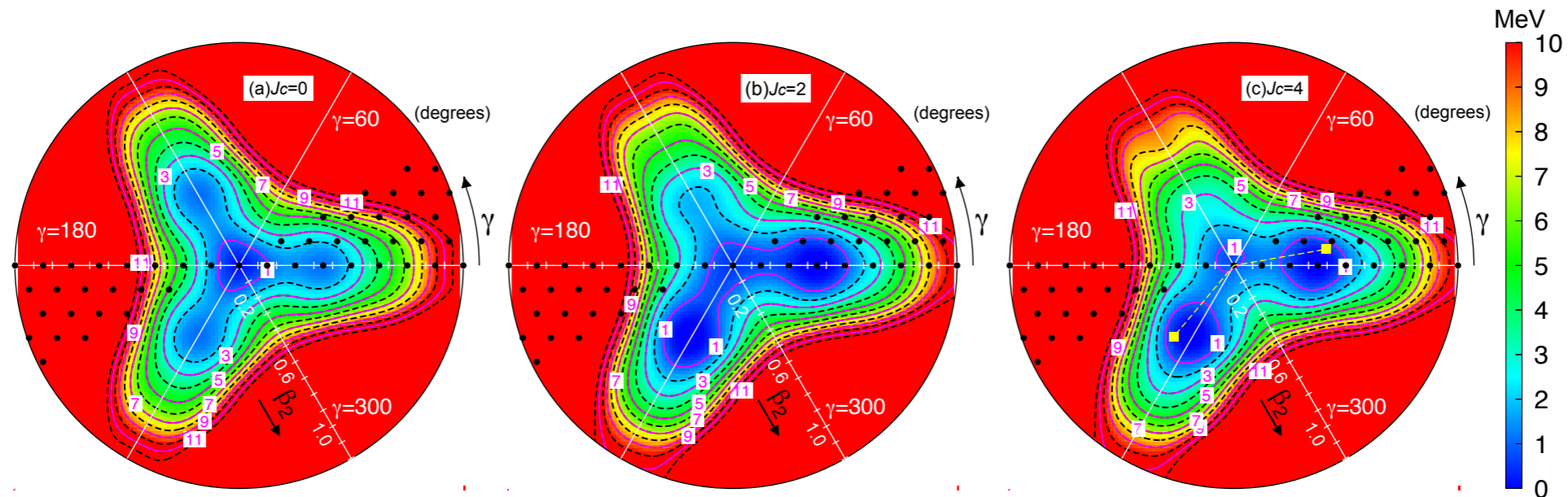
Gogny EDF with cranked states

Example: ^{32}Mg triaxial+TRSB



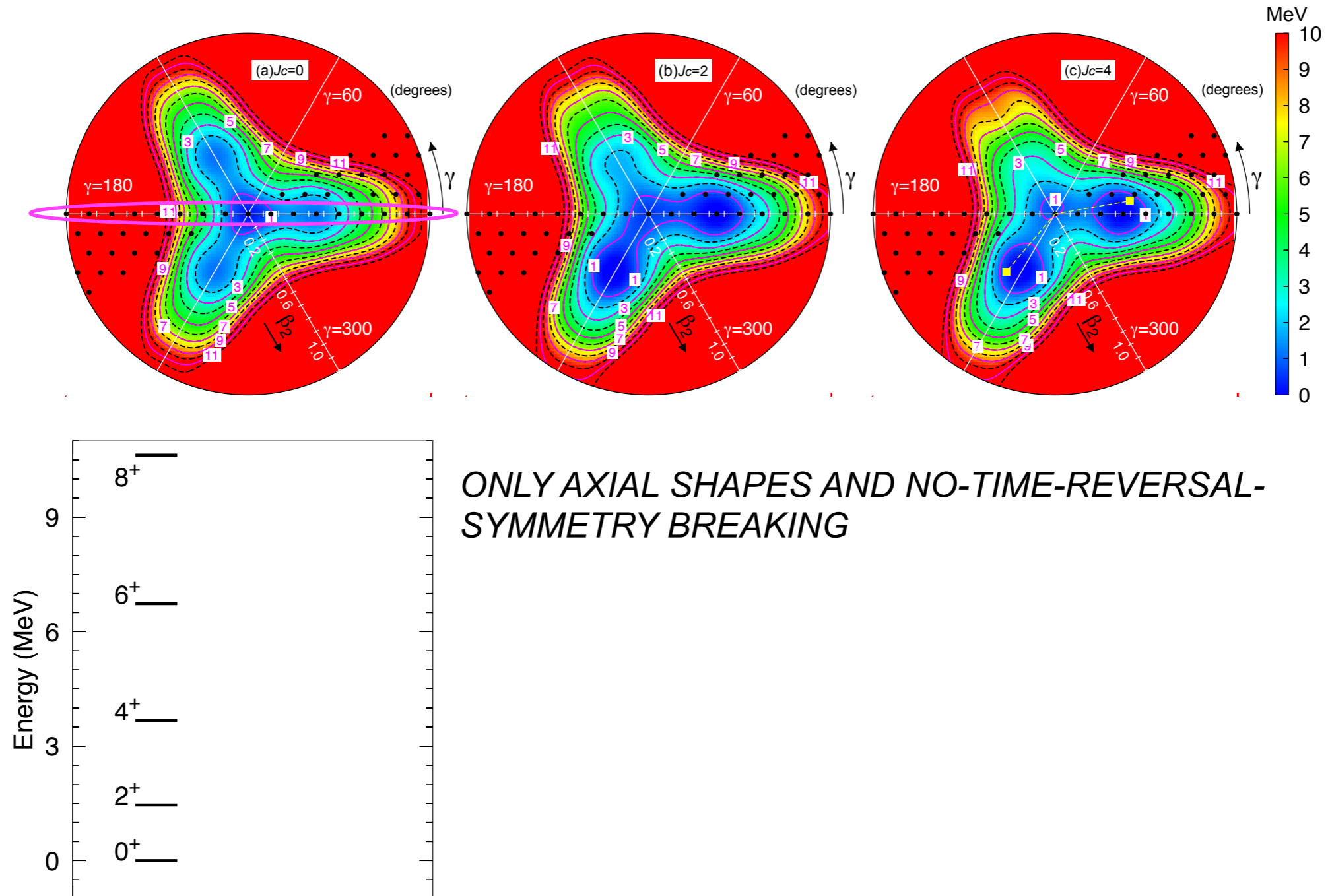
EDF with cranked states

Example: ^{32}Mg ; Effect on the energies



EDF with cranked states

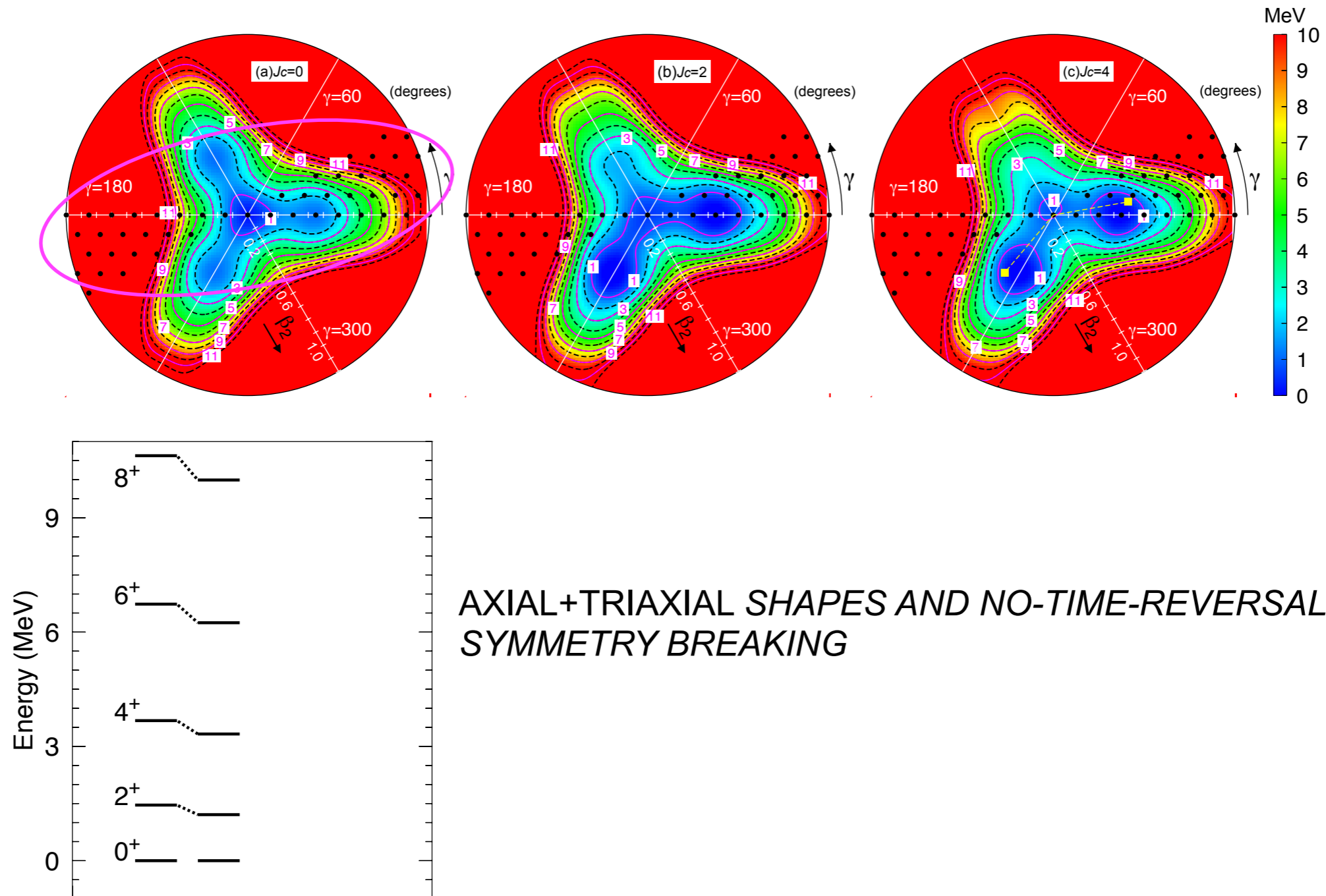
Example: ^{32}Mg ; Effect on the energies



M. Borrajo, T.R.R, J.L. Egido, PLB 746, 341 (2015)

EDF with cranked states

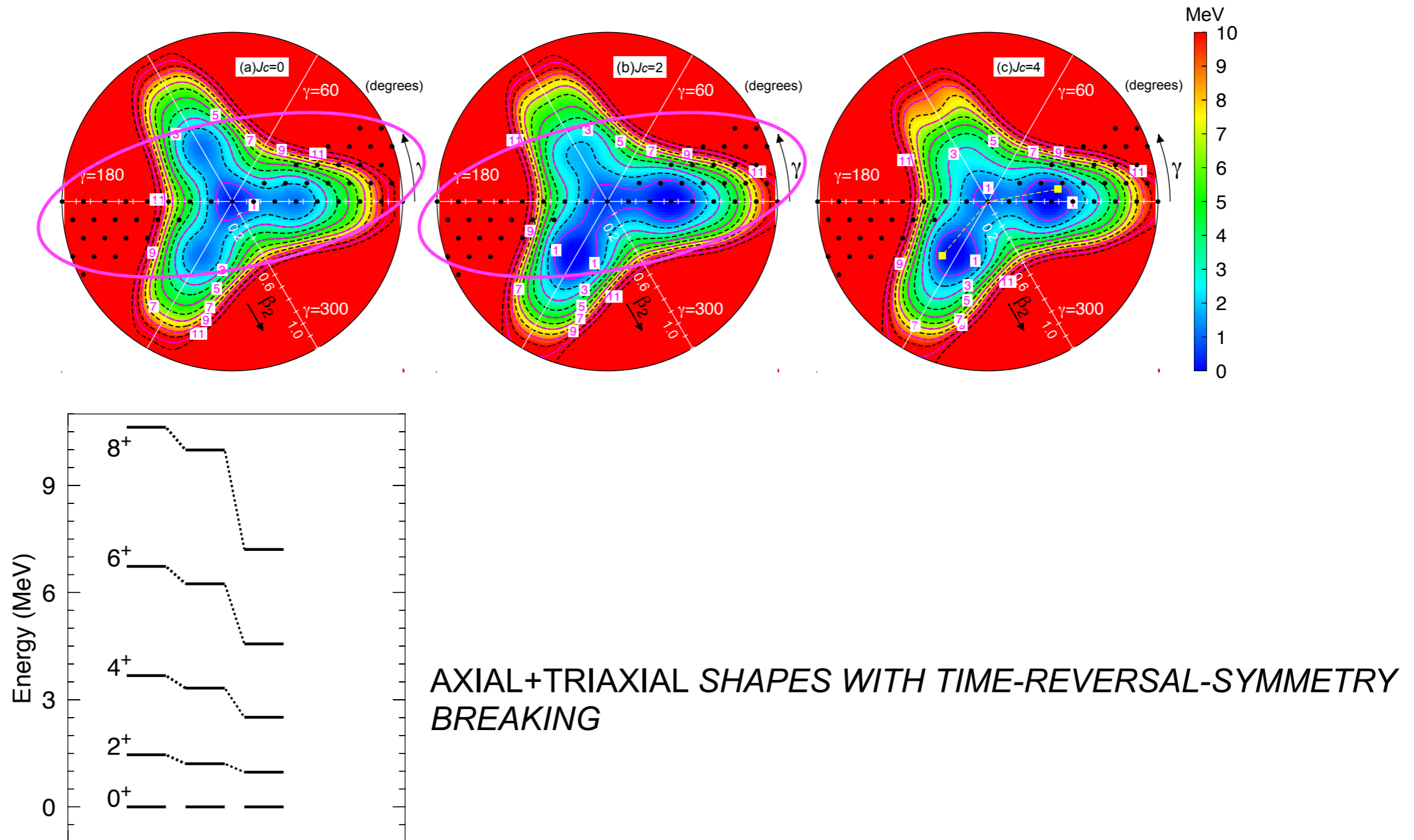
Example: ^{32}Mg ; Effect on the energies



M. Borrajo, T.R.R, J.L. Egido, PLB 746, 341 (2015)

EDF with cranked states

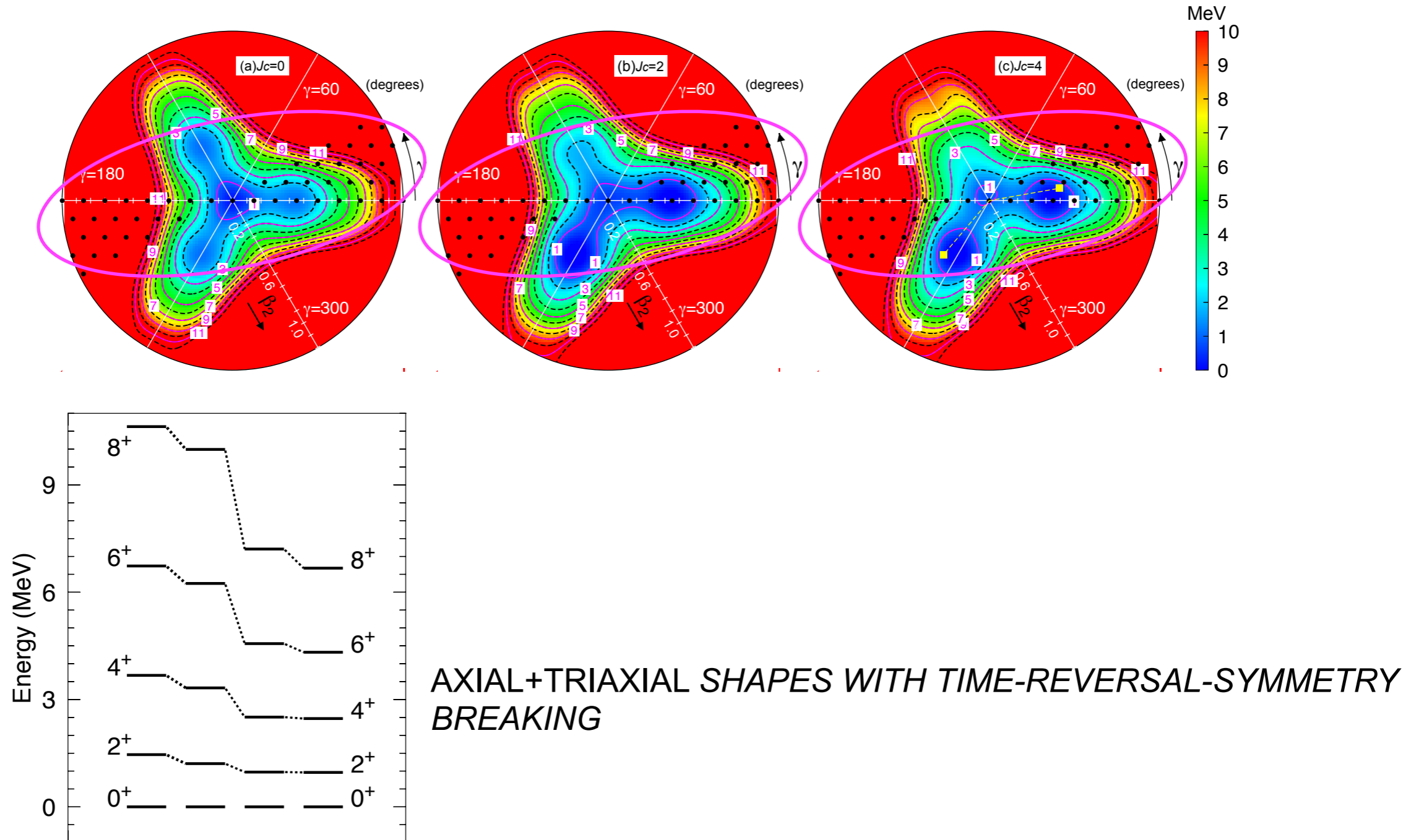
Example: ^{32}Mg ; Effect on the energies



M. Borrajo, T.R.R, J.L. Egido, PLB 746, 341 (2015)

EDF with cranked states

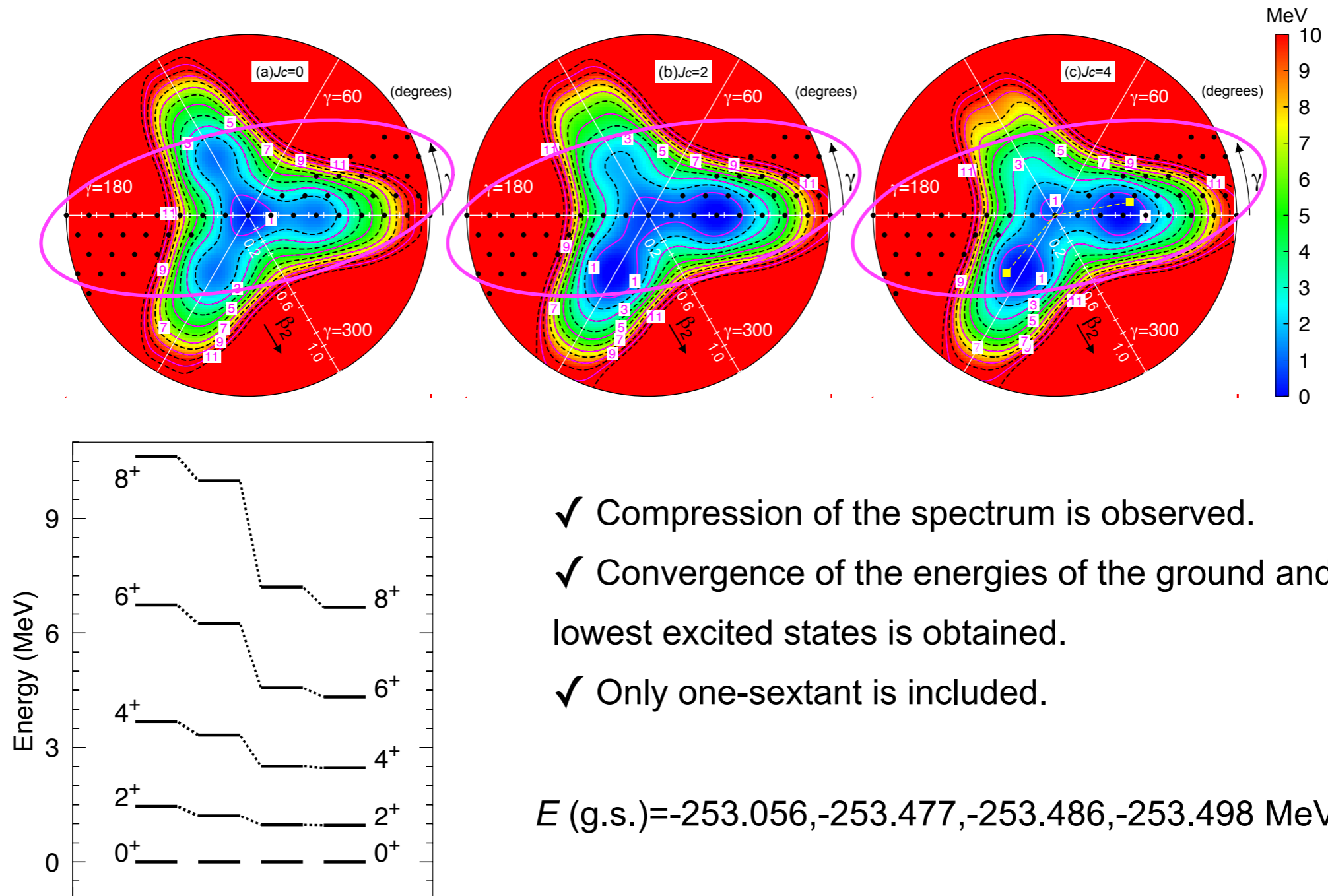
Example: ^{32}Mg ; Effect on the energies



M. Borrajo, T.R.R, J.L. Egido, PLB 746, 341 (2015)

EDF with cranked states

Example: ^{32}Mg ; Effect on the energies



M. Borrajo, T.R.R, J.L. Egido, PLB 746, 341 (2015)

EDF with cranked states



Introduction

SCCM with cranking

Collective and single-particle states in ^{44}S

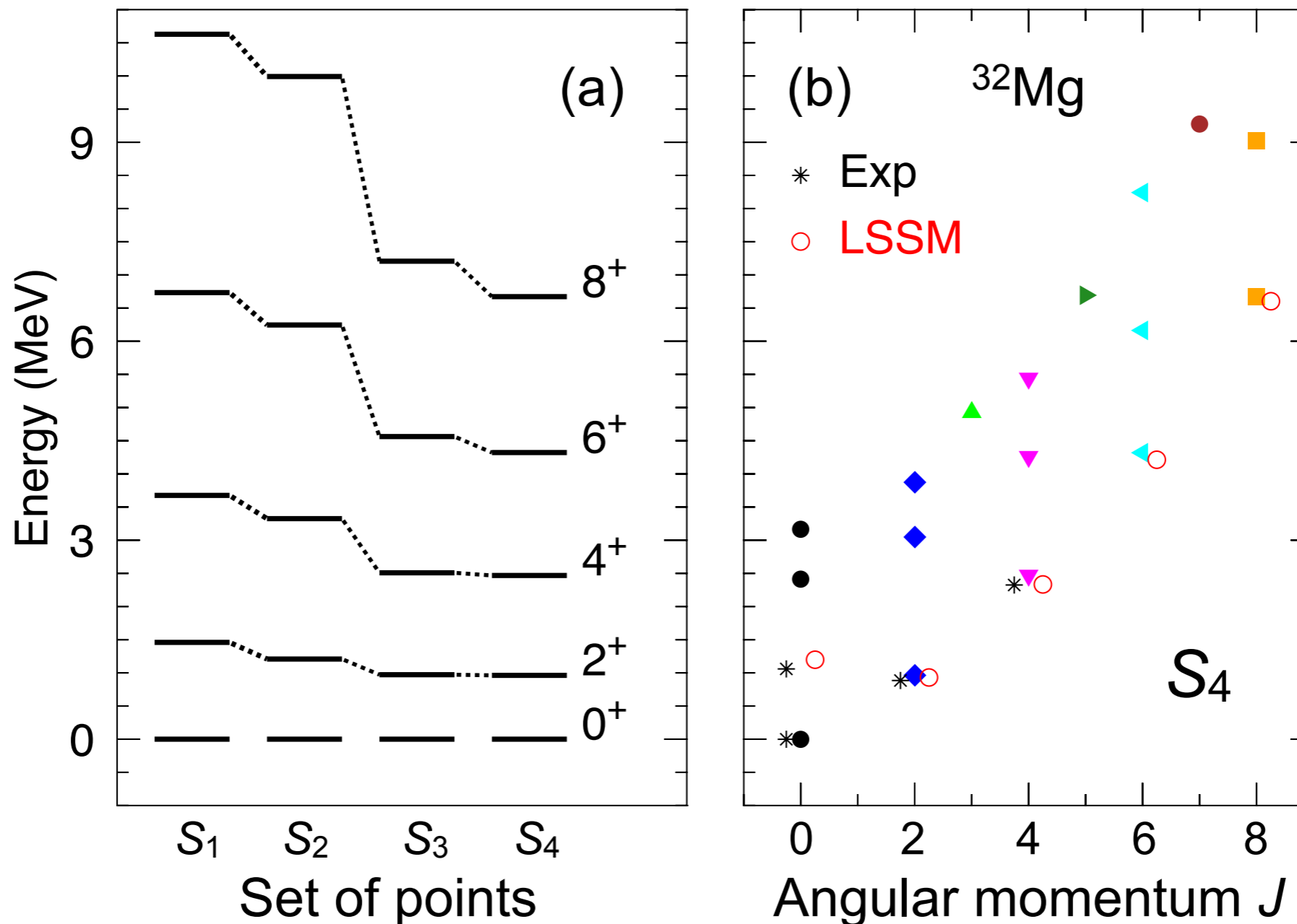
Summary and Outlook

Example: ^{32}Mg ; Comparison to SM and experiment

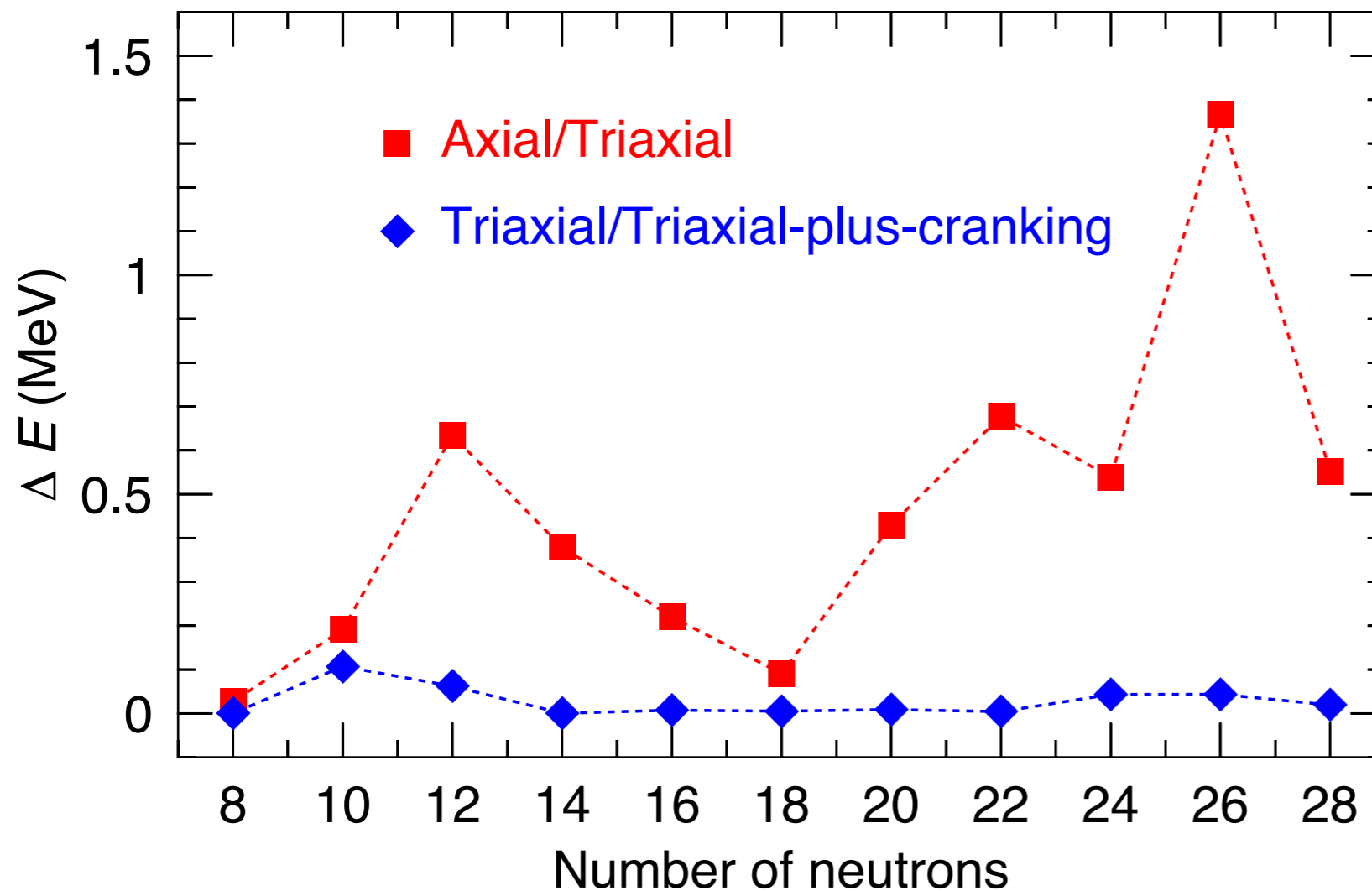
M. Borrajo, T.R.R, J.L. Egidio, PLB 746, 341 (2015)

EDF with cranked states

Example: ^{32}Mg ; Comparison to SM and experiment

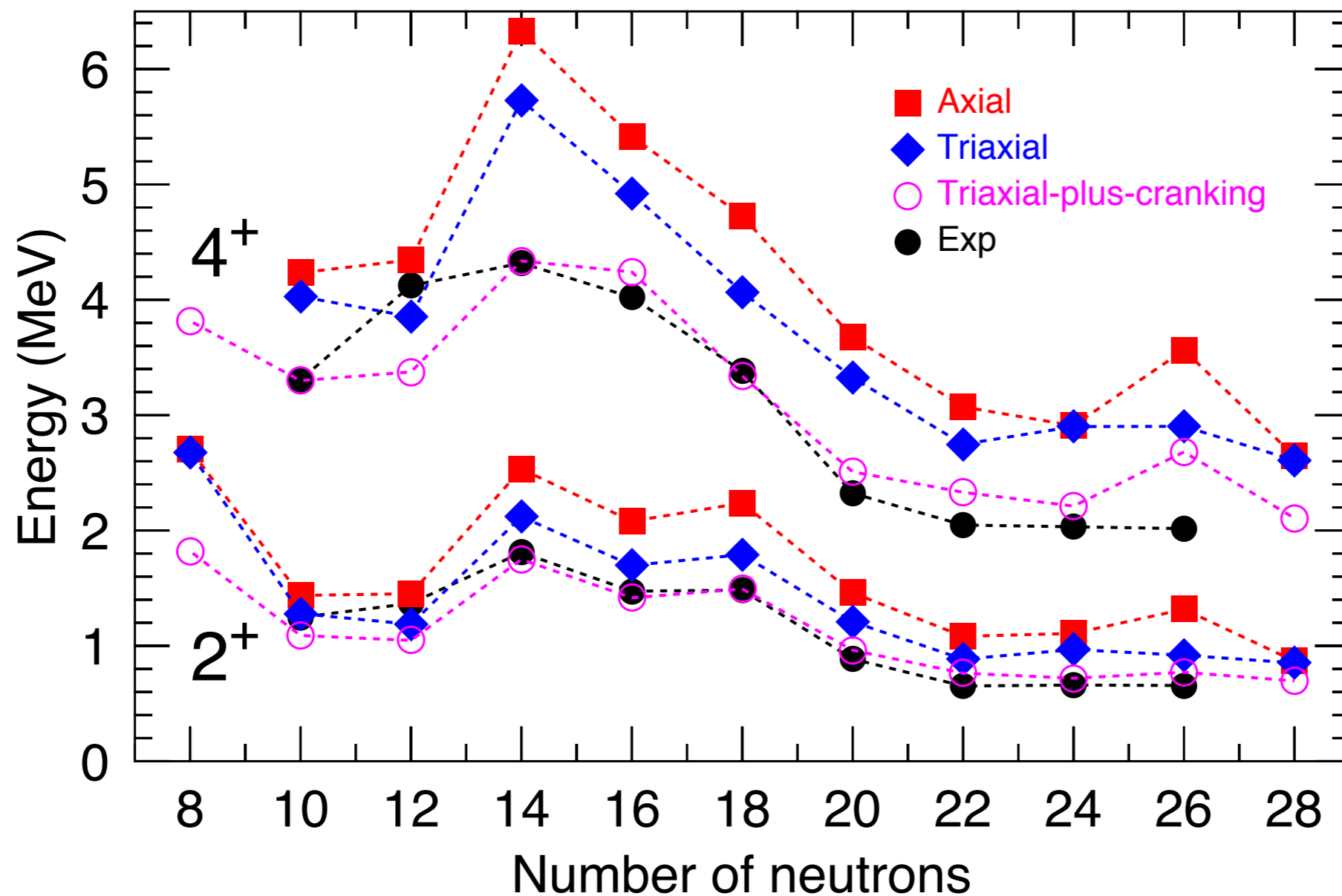


Systematics in the Magnesium isotopic chain



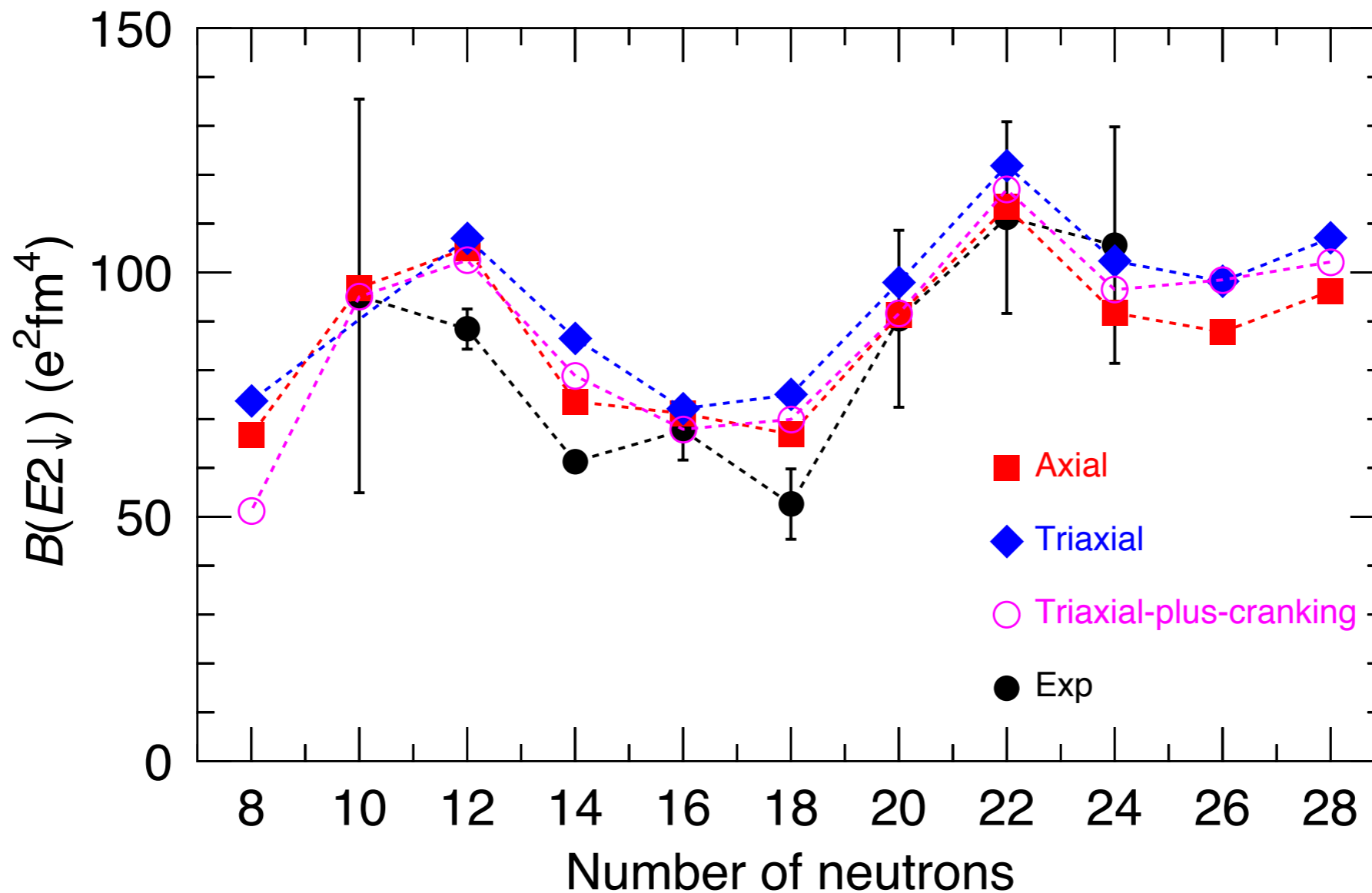
✓ Convergence of the energies of the ground states

Systematics in the Magnesium isotopic chain



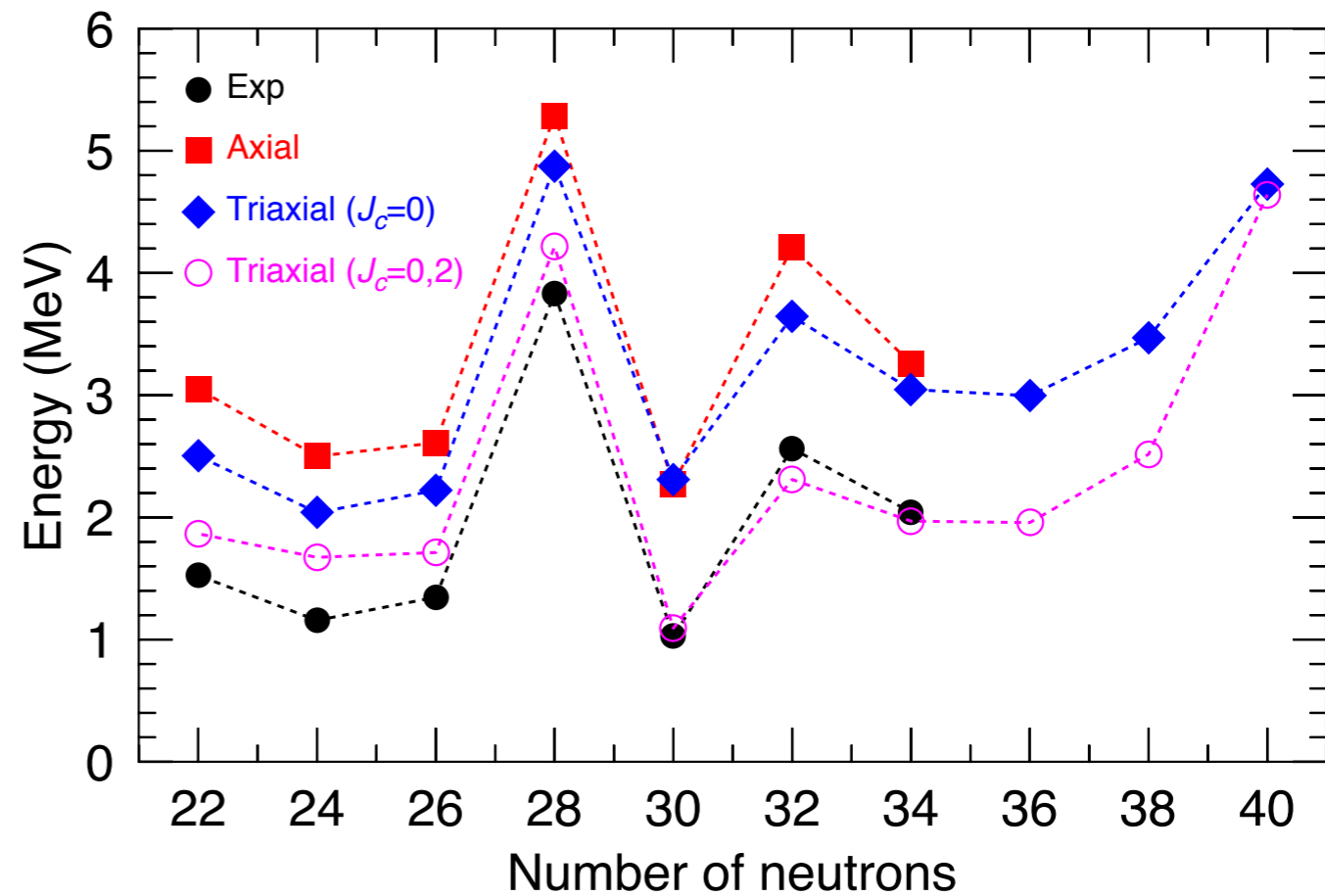
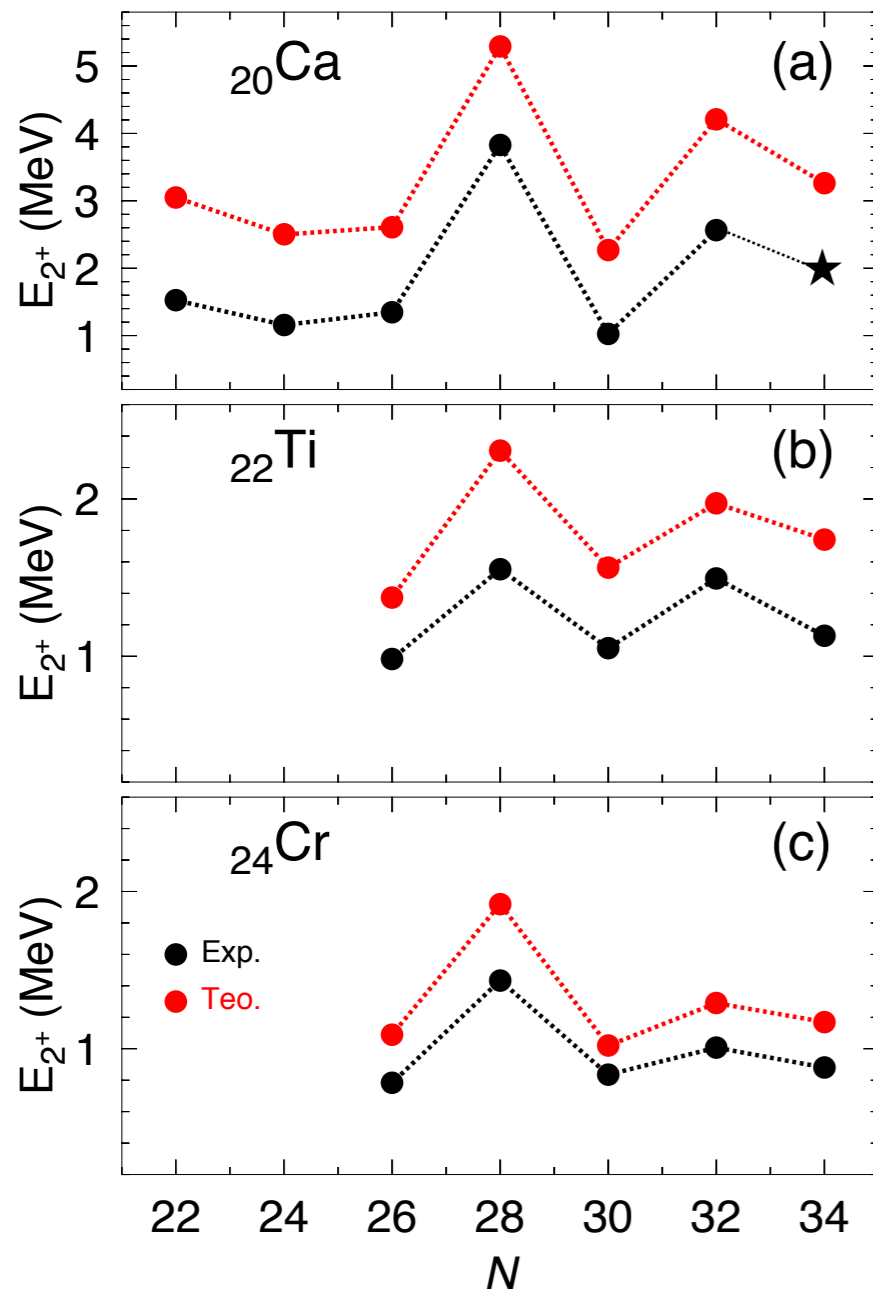
M. Borrajo, T.R.R, J.L. Egido, PLB 746, 341 (2015)

Systematics in the Magnesium isotopic chain



EDF with cranked states

$N=32$ and/or $N=34$ in Ca, Ti and Cr isotopes.

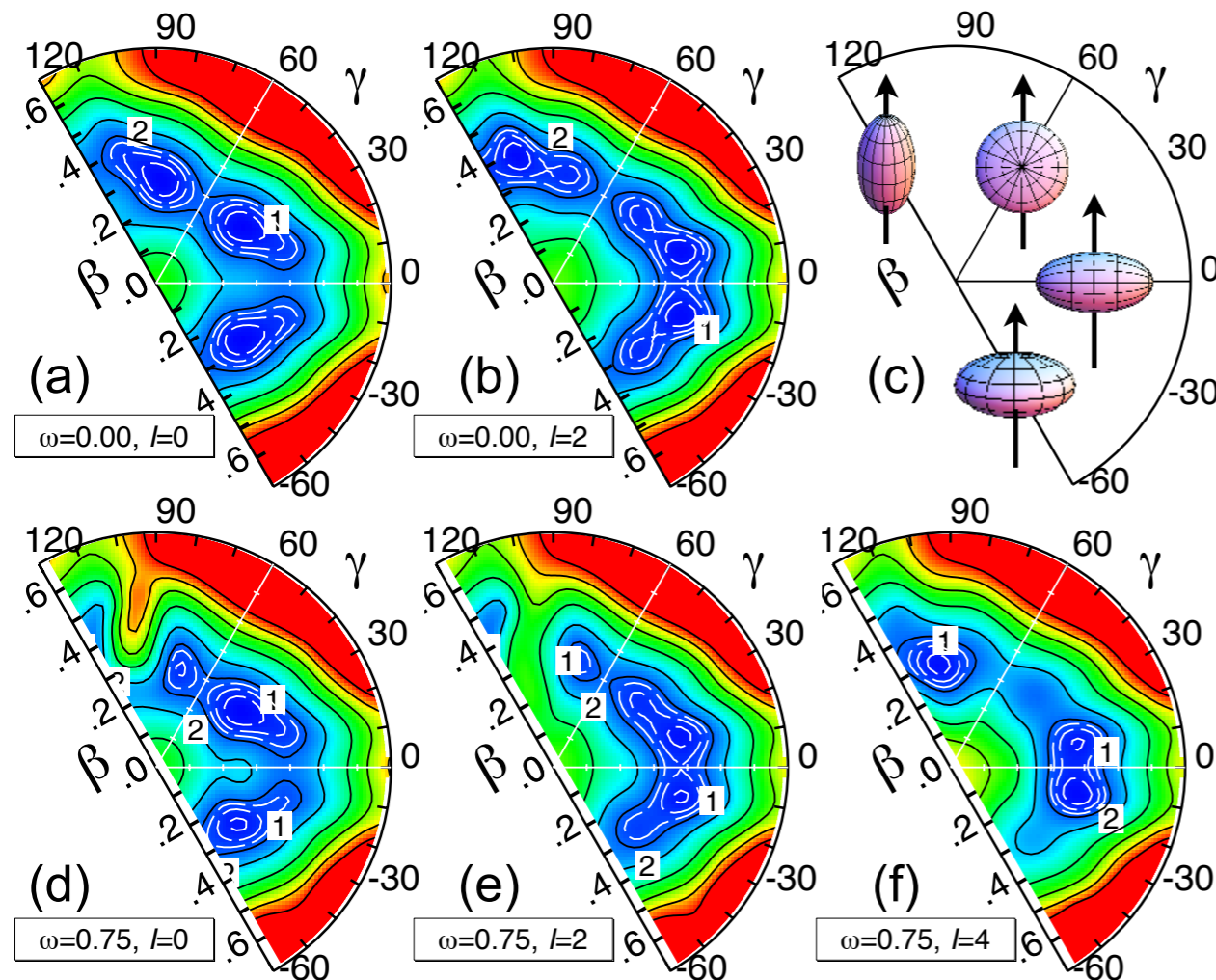


T.R.R, M. Borrajo, J.L. Egido, unpublished

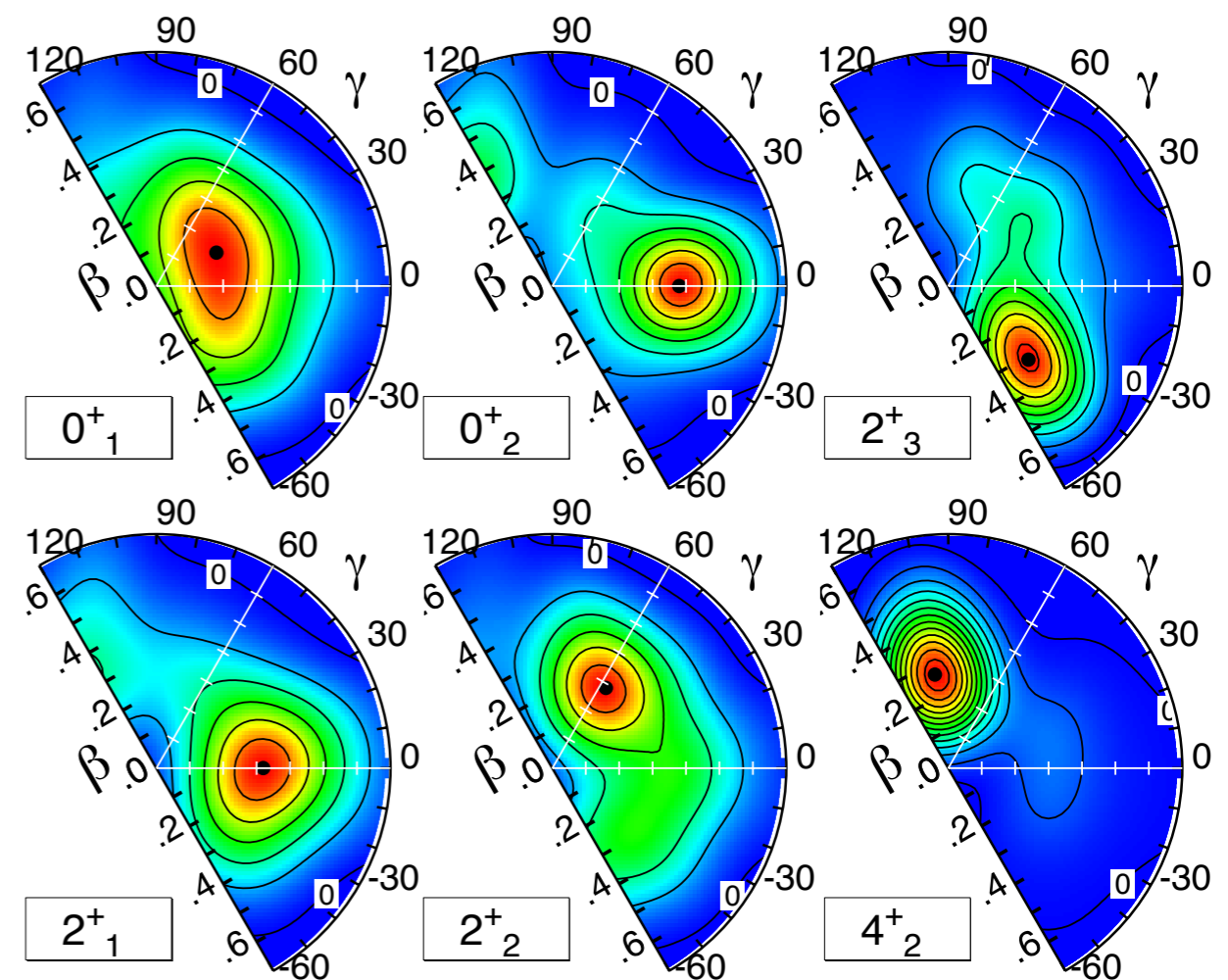
EDF with cranked states

^{44}S

Potential energy surfaces

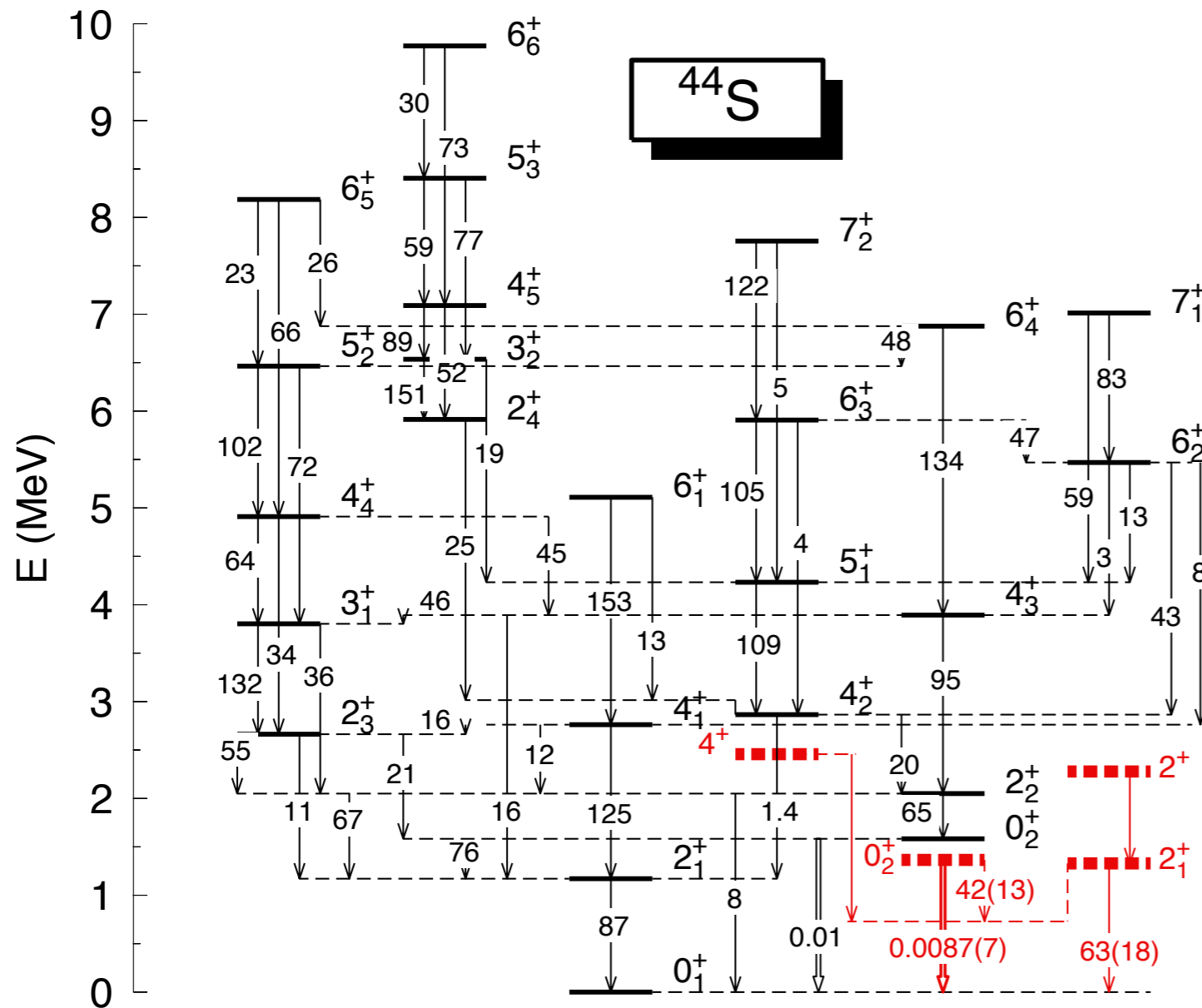


Collective wave functions

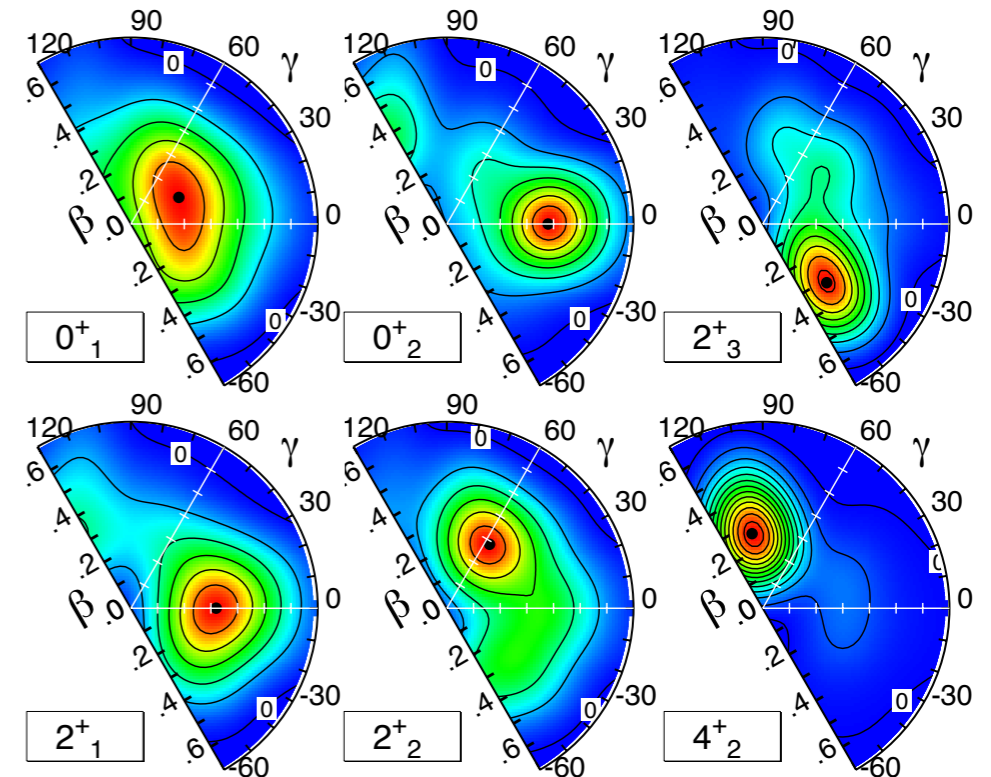


J.L. Egido, M. Borrajo, TRR, PRL 116, 054319 (2016)

EDF with cranked states



Collective wave functions



EDF with cranked states

^{44}S

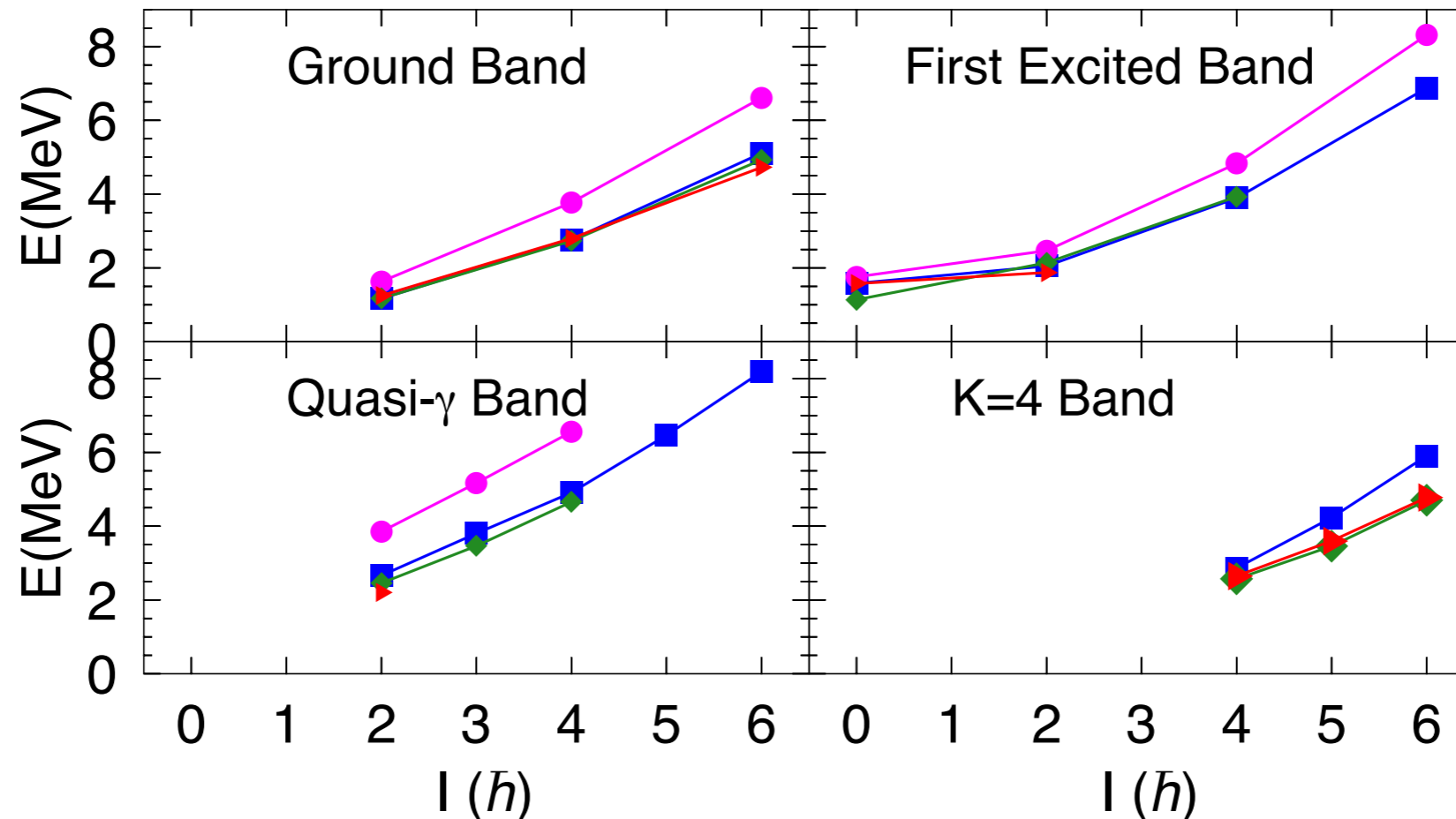


FIG. 4: (Color online) Comparison of several theories: Triangles, red lines, Tokyo group [22]; diamonds, green lines, Madrid-Strasbourg collaboration [31]; boxes, blue lines, this work; circles, magenta lines, our former work without angular frequency dependence [20].

Summary and outlook

Time-reversal symmetry (cranking states) allows for a quantitative agreement with the experimental energy spectra.

Outlook

- Quasiparticle states:
 - Odd-nuclei. (Bally-Bender-Heenen, Borrajo-Egido, ...)
 - Single-particle excitations. (Qi-Egido, ...)
- $T_z=0$ pairing, pn pairing.
- Isospin projection. (Satula et al., ...)
- Parameters of the next generation of functionals should be fitted taking into account beyond-mean-field effects and/or extracted from *ab initio* methods (...).



RECENT DEVELOPMENTS IN MICROSCOPIC THEORIES FOR NUCLEAR STRUCTURE

Gogny 2017

November 29- December 1, 2017
Universidad Autónoma de Madrid (Spain)

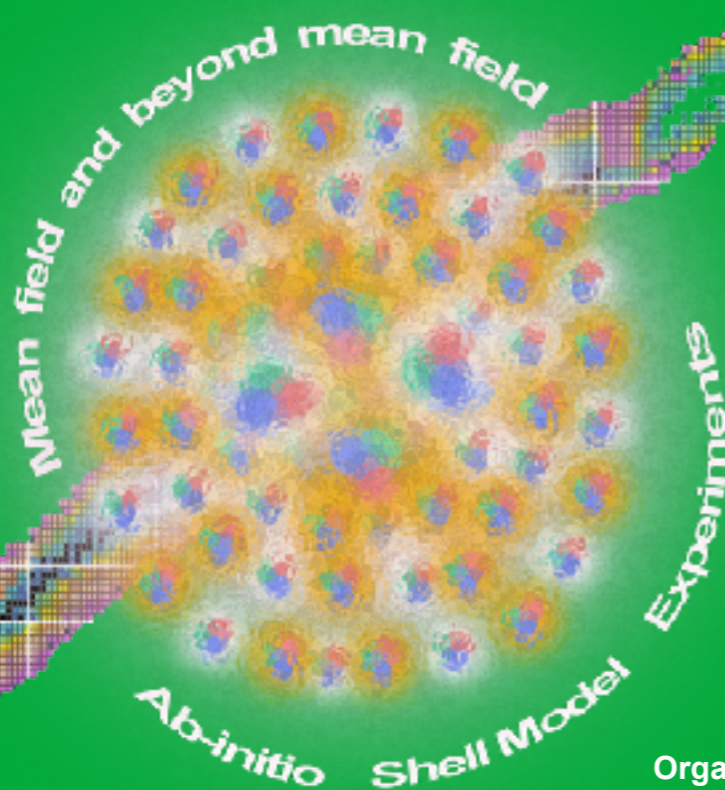


We look forward to welcoming you in Madrid!

<http://eventos.uam.es/go/gogny2017>

Advisory Committee

M. Baldo
K. Bennaceur
J. F. Berger
G. F. Berstch
G. Co'
J. P. Delaroche
M. Dupuis
M Girod
H. Goutte
S. Hilaire
D. Lacroix
H. Nakada
J. Navarro
N. Pillet
A. Polls
P. Ring
A. Ríos
R. Rodríguez-Guzmán
P. Schuck
X. Viñas
M Warda
W. Younes



Organizing Committee

L. M. Robledo
A. Poves
T. R. Rodríguez
M. Anguiano
E. Moya de Guerra
P. Sarriguren

Information

Milvia Soumbounou
<http://eventos.uam.es/go/gogny2017>
gogny2017@uam.es

