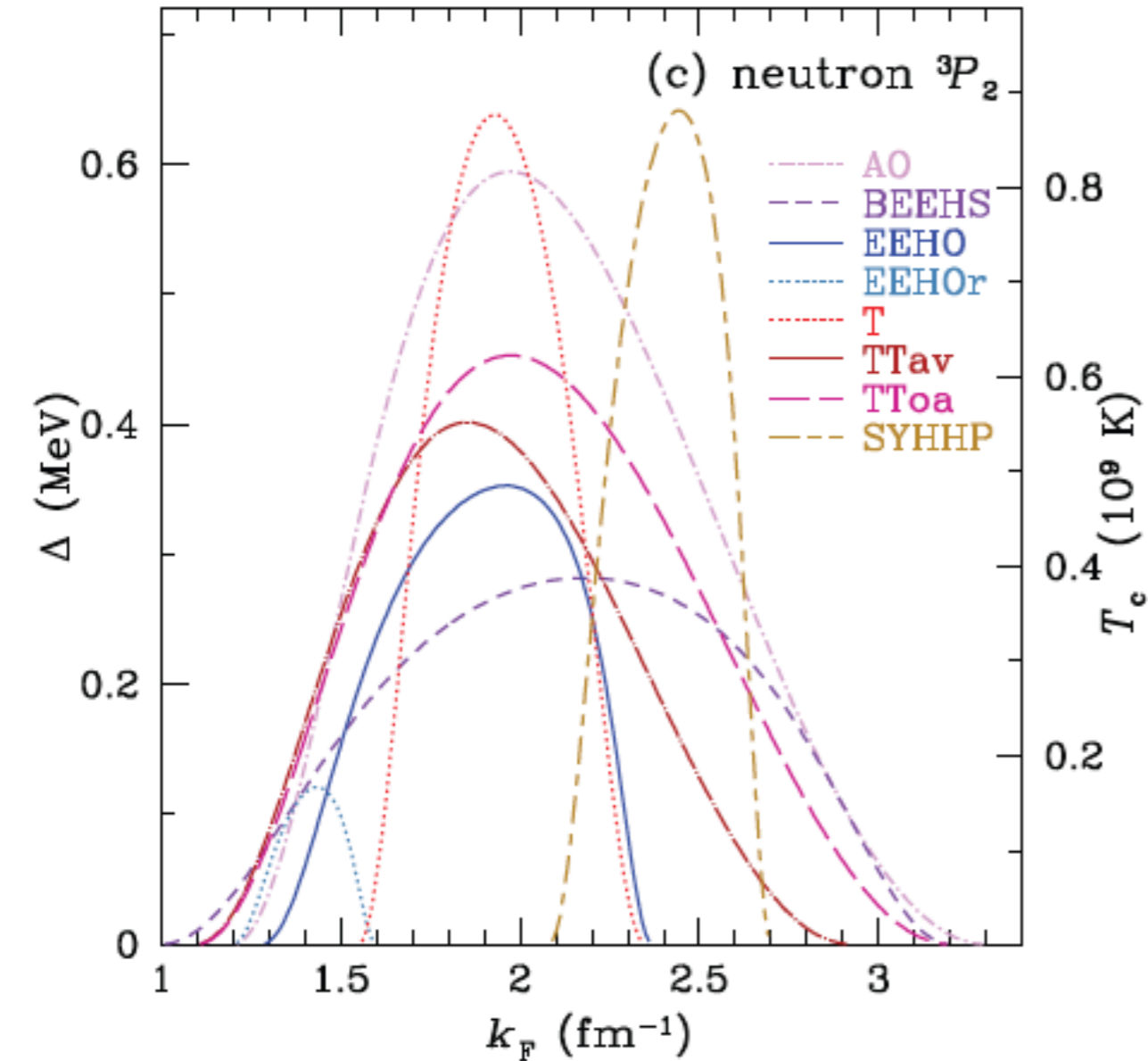


Pairing in high-density neutron matter

Arnau Rios Huguet
Lecturer in Nuclear Theory
Department of Physics
University of Surrey

- 1. Neutron star motivation**
2. Infinite matter BCS
3. Beyond-BCS with SCGF methods

Cooling of CasA

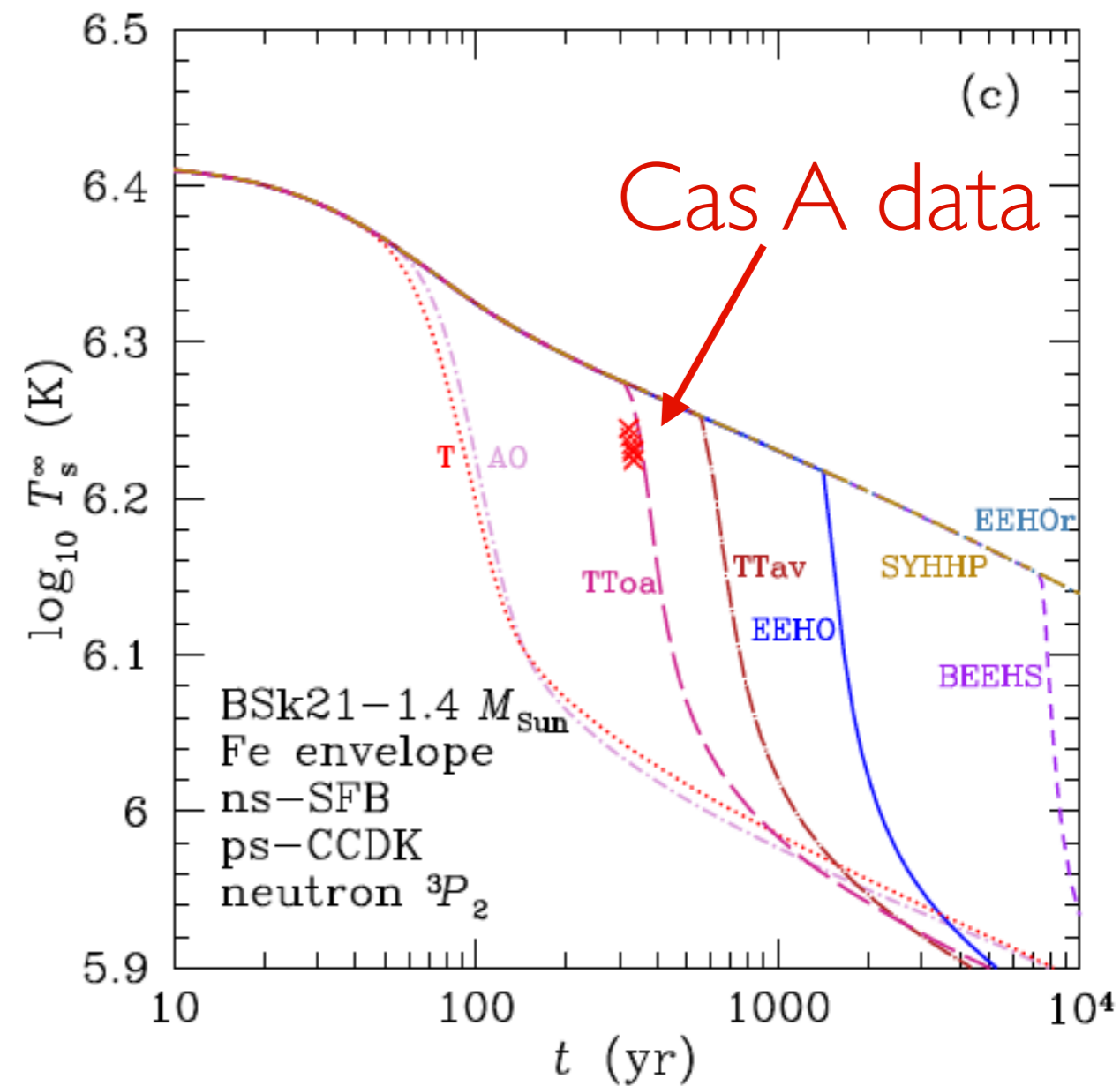


Ho, et al., PRC **91** 015806 (2015)

Page, et al., PRL **106** 081101 (2011)

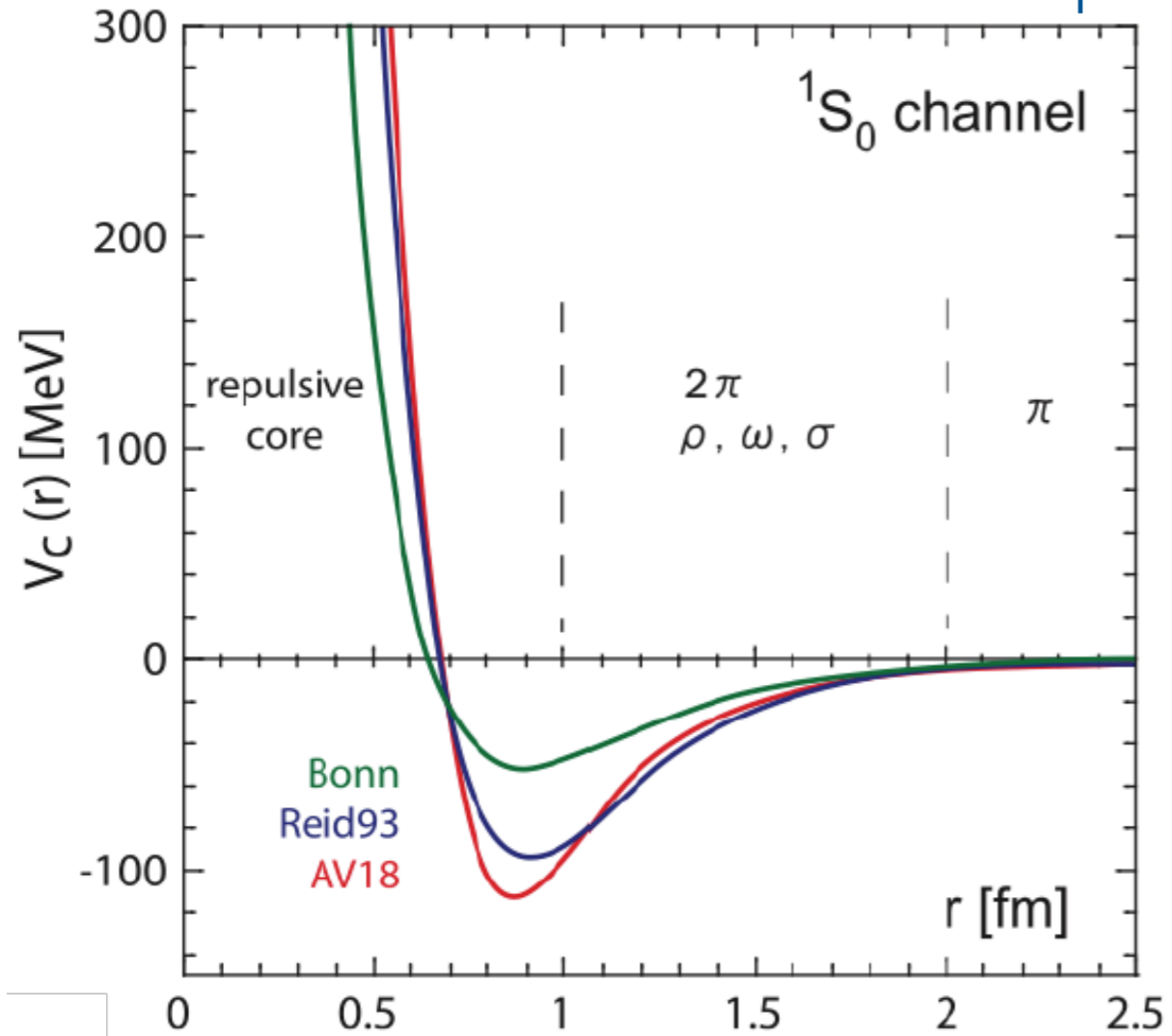
Ingredients

- (a) Mass of pulsar
- (b) EoS (determines radius)
- (c) Internal composition
- (d) **Pairing gaps** (1S_0 & 3P_2 channels)
- (e) Atmosphere composition



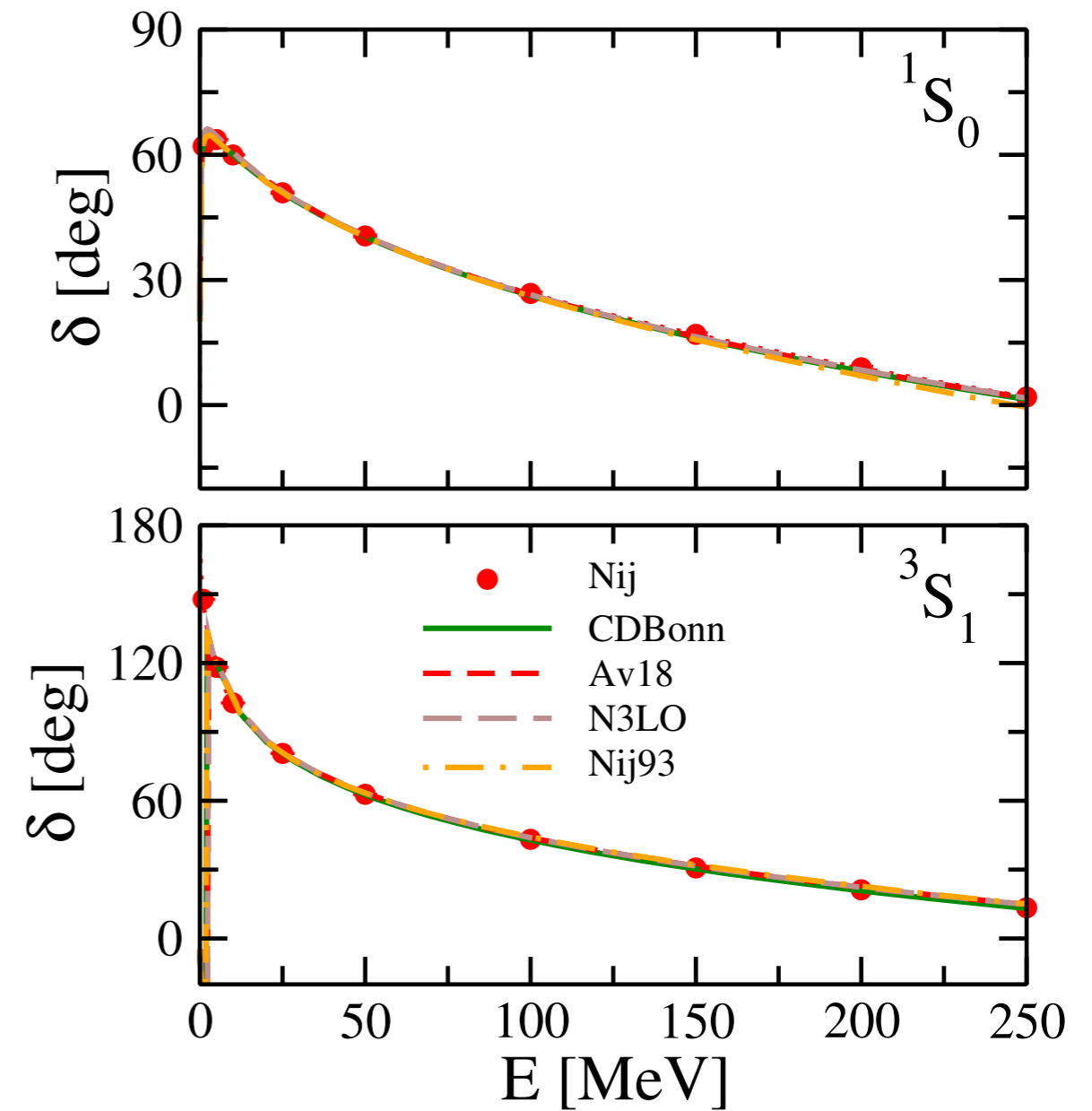
Name	Process	Emissivity ($\text{erg cm}^{-3} \text{s}^{-1}$)
Modified Urca (neutron branch)	$n+n \rightarrow n+p+e^-+\bar{\nu}_e$	$\sim 2 \times 10^{21} RT_9^8$
	$n+p+e^- \rightarrow n+n+\nu_e$	
Modified Urca (proton branch)	$p+n \rightarrow p+p+e^-+\bar{\nu}_e$	$\sim 10^{21} RT_9^8$
	$p+p+e^- \rightarrow p+n+\nu_e$	
Bremsstrahlungs	$n+n \rightarrow n+n+\nu+\bar{\nu}$	$\sim 10^{19} RT_9^8$
	$n+p \rightarrow n+p+\nu+\bar{\nu}$	
	$p+p \rightarrow p+p+\nu+\bar{\nu}$	
Cooper pair	$n+n \rightarrow [nn]+\nu+\bar{\nu}$	$\sim 5 \times 10^{21} RT_9^7$
	$p+p \rightarrow [pp]+\nu+\bar{\nu}$	$\sim 5 \times 10^{19} RT_9^7$
Direct Urca (nucleons)	$n \rightarrow p+e^-+\bar{\nu}_e$	$\sim 10^{27} RT_9^6$
	$p+e^- \rightarrow n+\nu_e$	

NN interaction is not unique

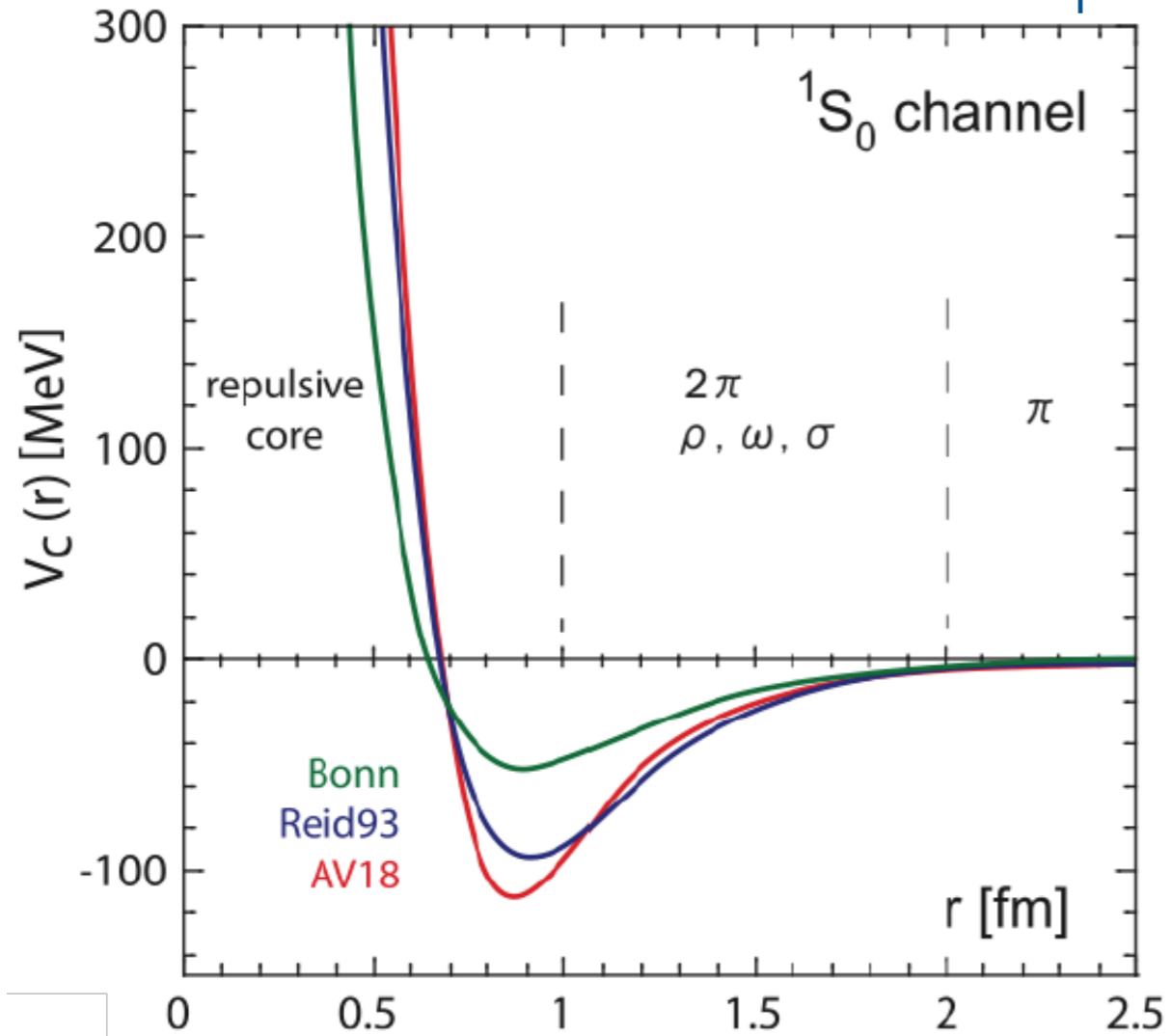


S.Aoki, et al. *Comput. Sci. Dis.* | 015009 (2008)

...but phase-shift equivalent!

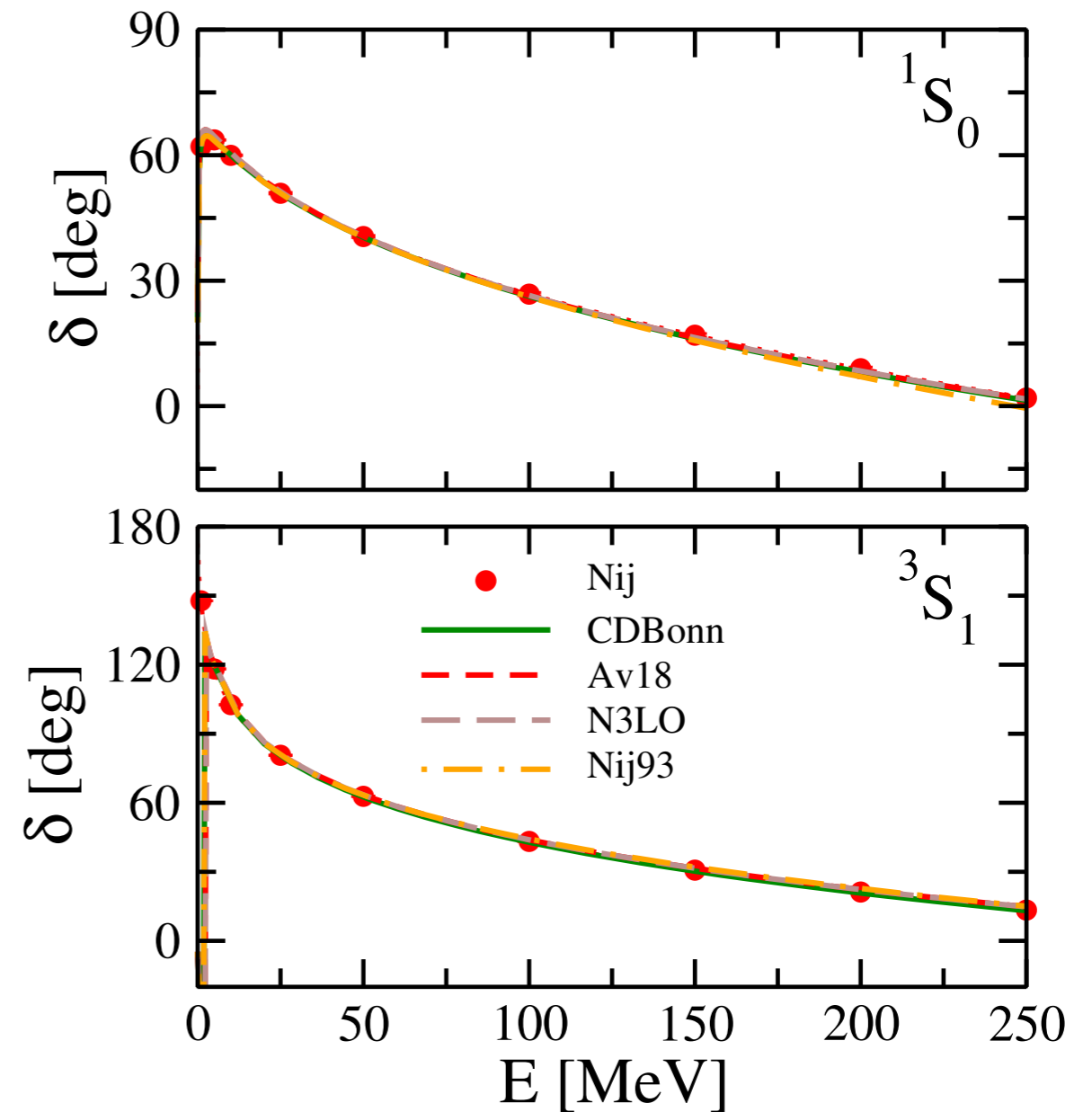


NN interaction is not unique



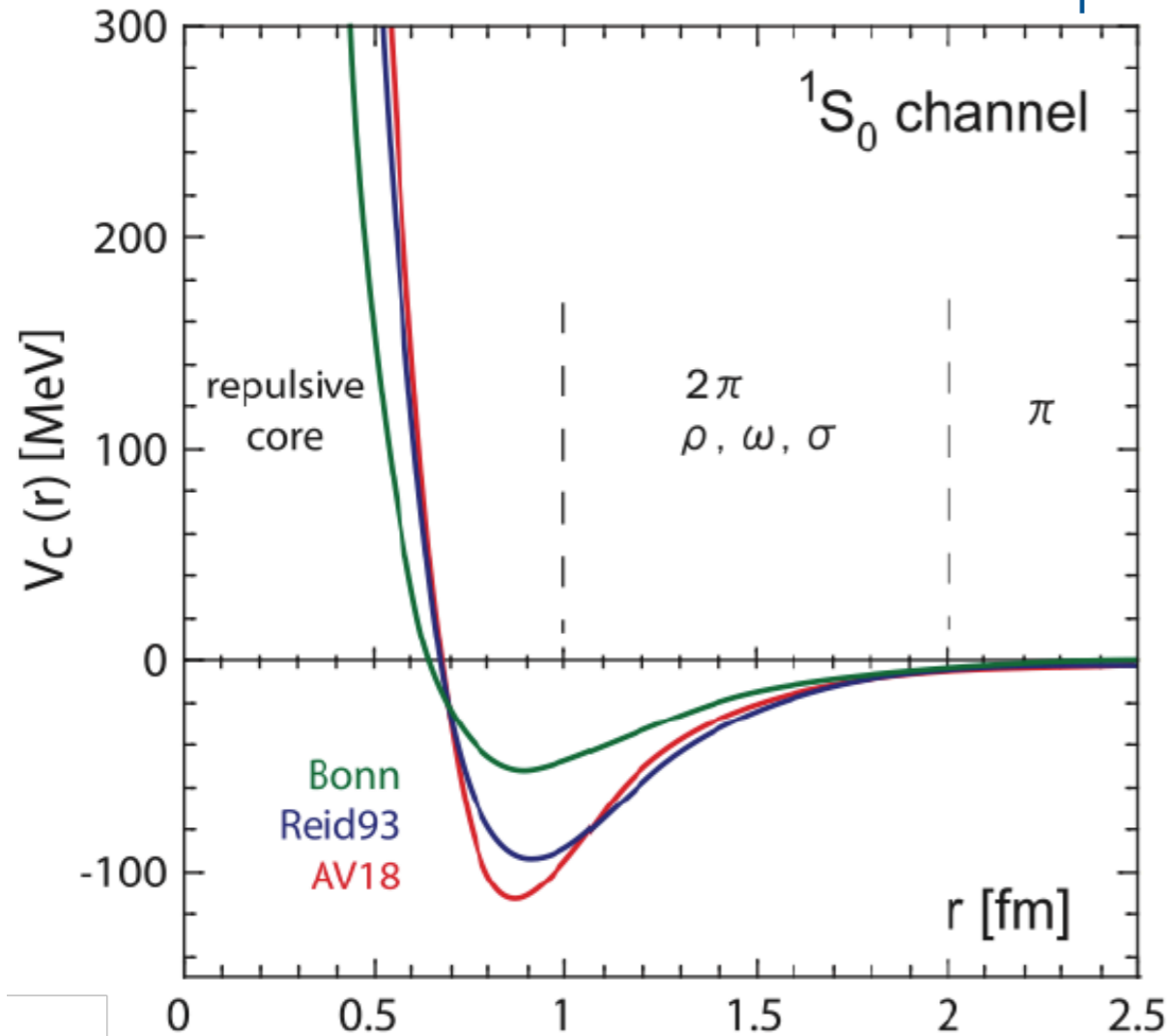
S.Aoki, et al. *Comput. Sci. Dis.* | 015009 (2008)

...but phase-shift equivalent!



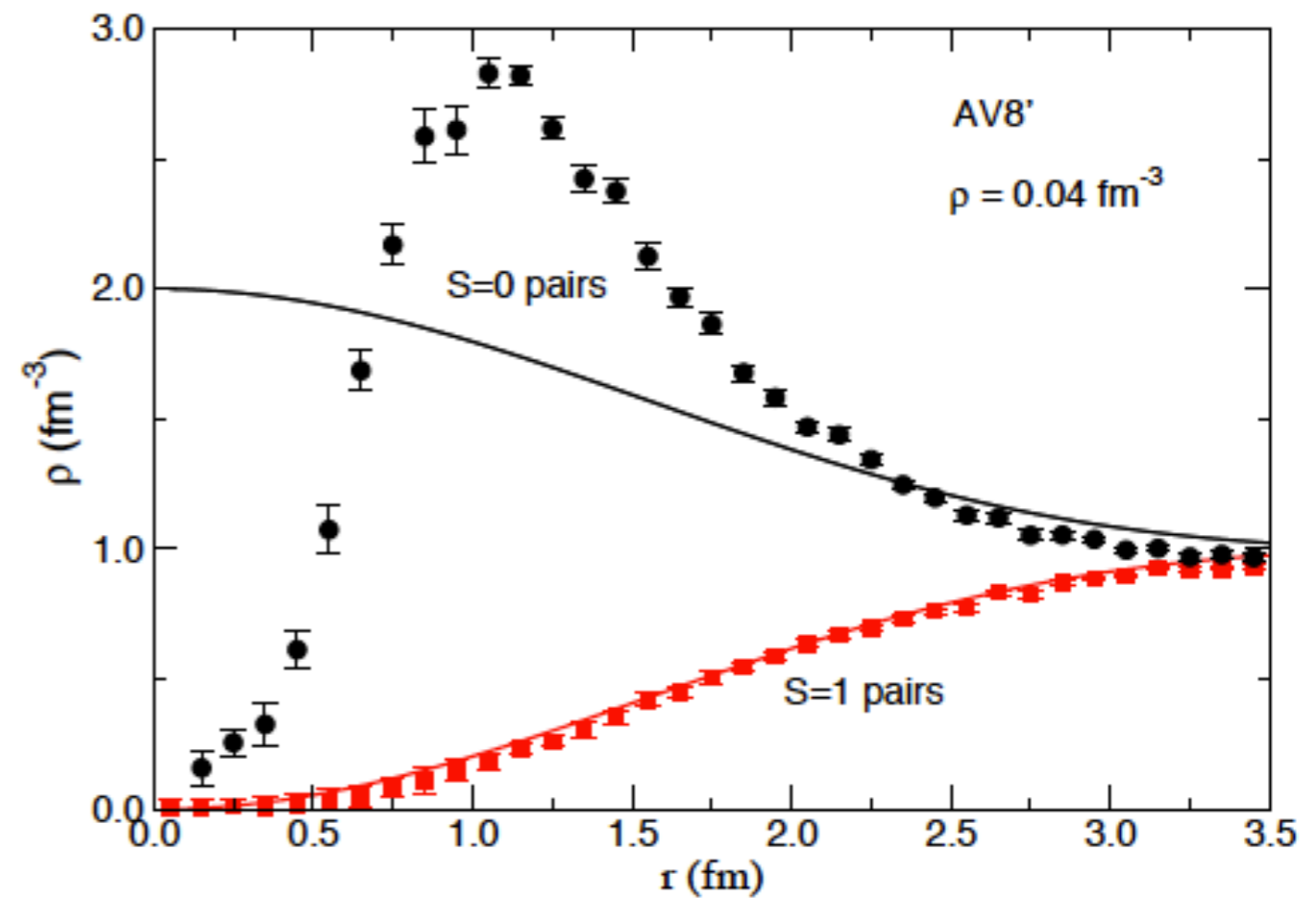
- Non-uniqueness of nucleon forces ✗

NN interaction is not unique



S.Aoki, et al. *Comput. Sci. Dis.* **1** 015009 (2008)

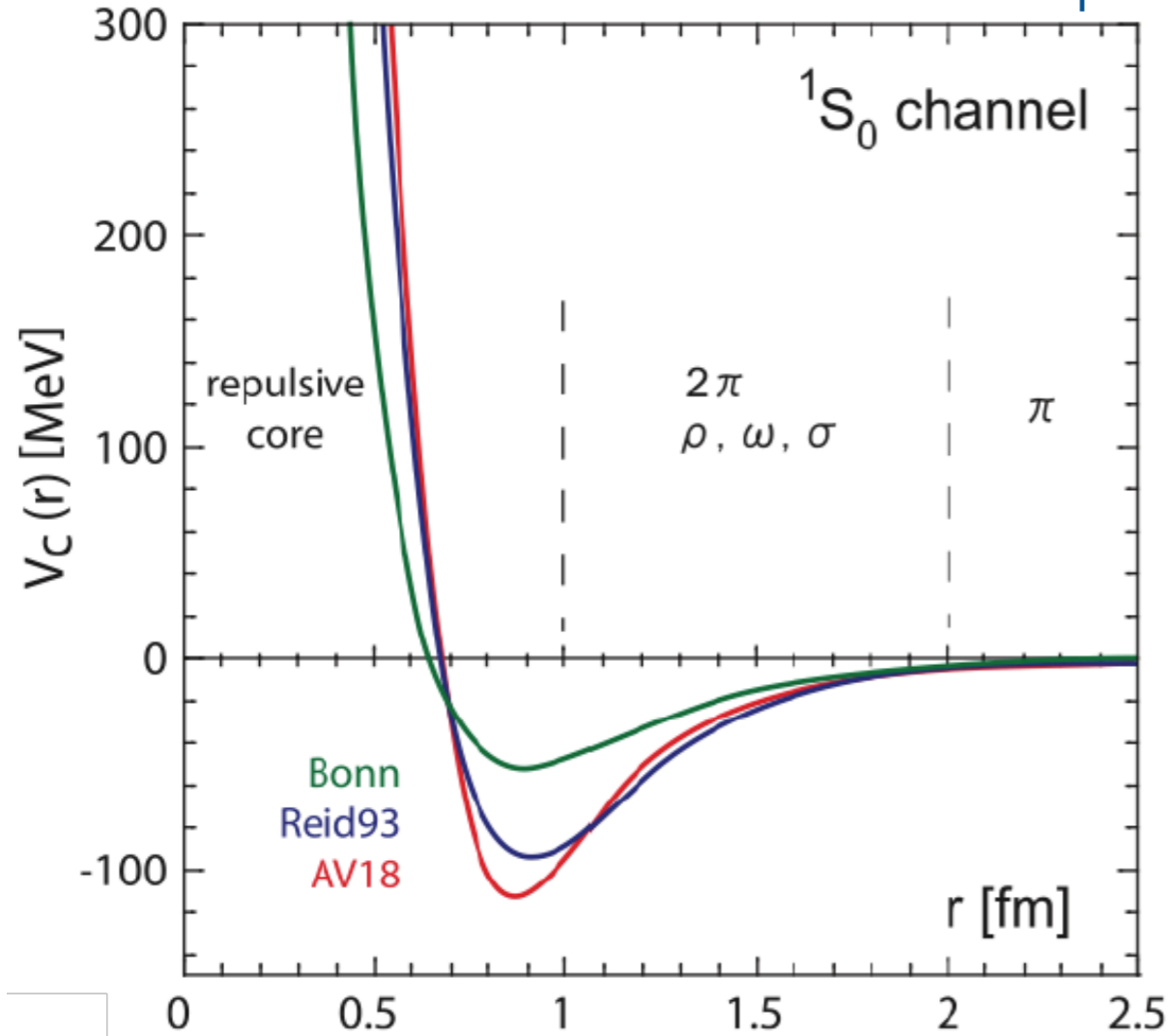
Strong short-range correlations



Carlson et al., *Phys. Rev. C* **68** 025802 (2003)

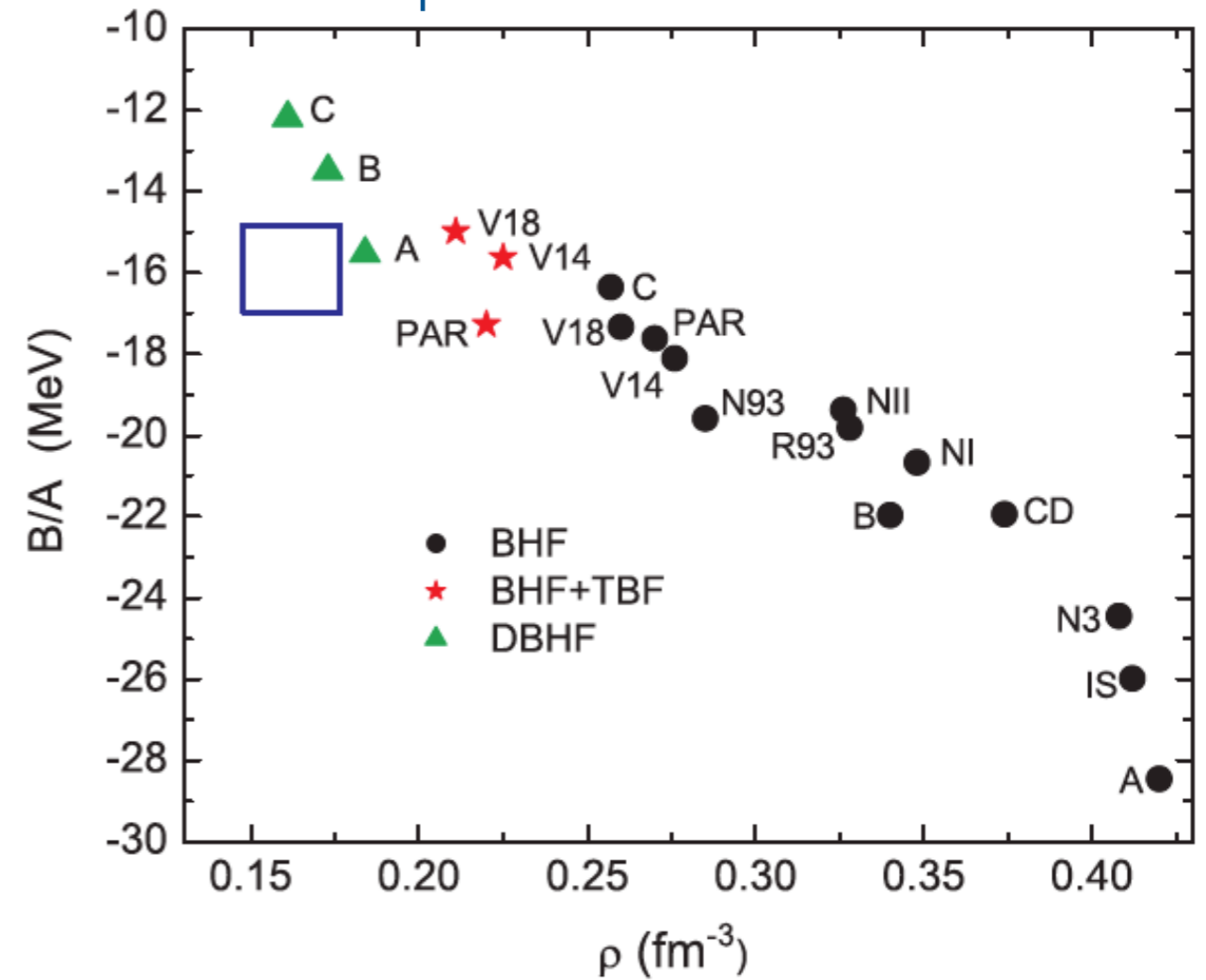
- Non-uniqueness of nucleon forces ✗
- Short-range core needs many-body treatment ✗

NN interaction is not unique



S.Aoki, et al. *Comput. Sci. Dis.* **1** 015009 (2008)

Saturation point of nuclear matter



Li, Lombardo, Schulze et al. *PRC* **74** 047304 (2006)

- Non-uniqueness of nucleon forces ✗
- Short-range core needs many-body treatment ✗
- Three-body forces needed for saturation ✗

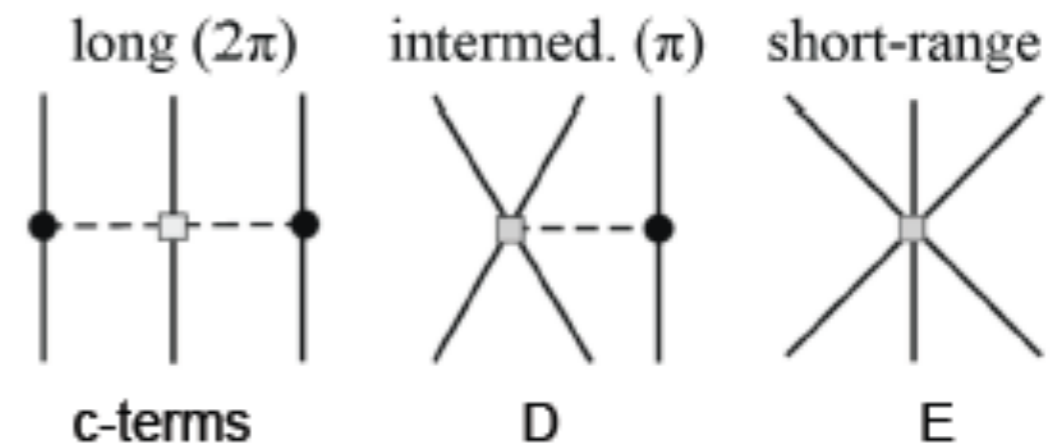
	NN	3N	4N
LO			
NLO			
N ² LO			
N ³ LO			

Chiral perturbation theory

- π and N as dof
- Systematic expansion
- 2N at N³LO - LECs from π N, NN
- 3N at N²LO - 2 more LECs
- (Often further renormalized)

$$\mathcal{O}\left(\frac{Q}{\Lambda}\right)$$

$$\Lambda \sim 1 \text{ GeV}$$



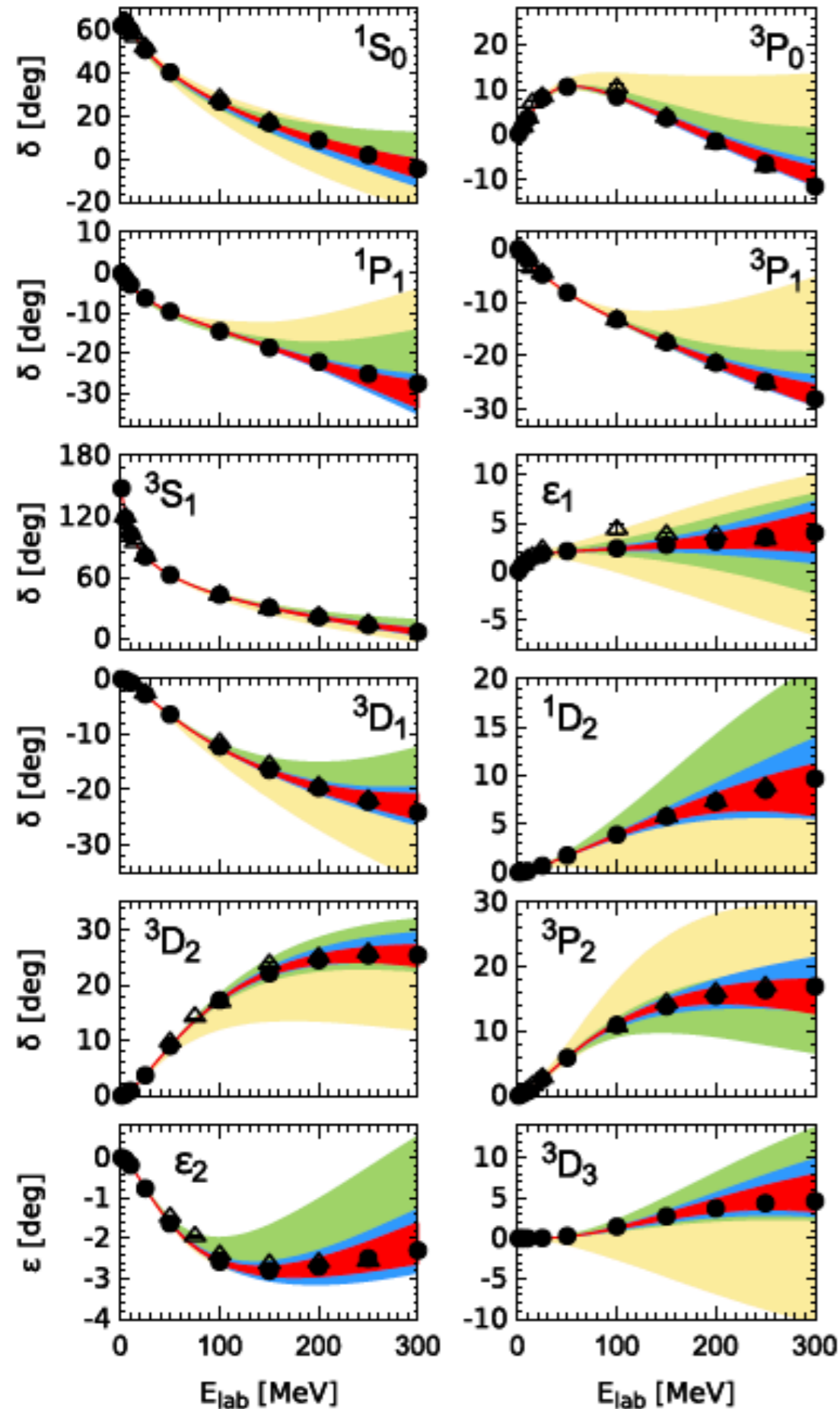
Weinberg, *Phys. Lett. B* **251** 288 (1990), *NPB* **363** 3 (1991)

Entem & Machleidt, *PRC* **68**, 041001(R) (2003)

Tews, Schwenk et al., *PRL* **110**, 032504 (2013)

Epelbaum, Friese & Meissner, *PRL* **115**, 122301 (2015)

NN forces from EFTs of QCD

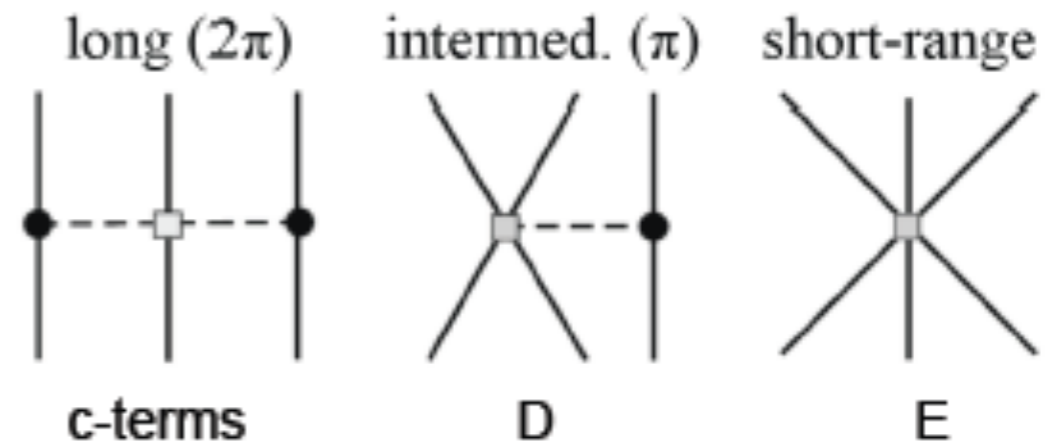


Chiral perturbation theory

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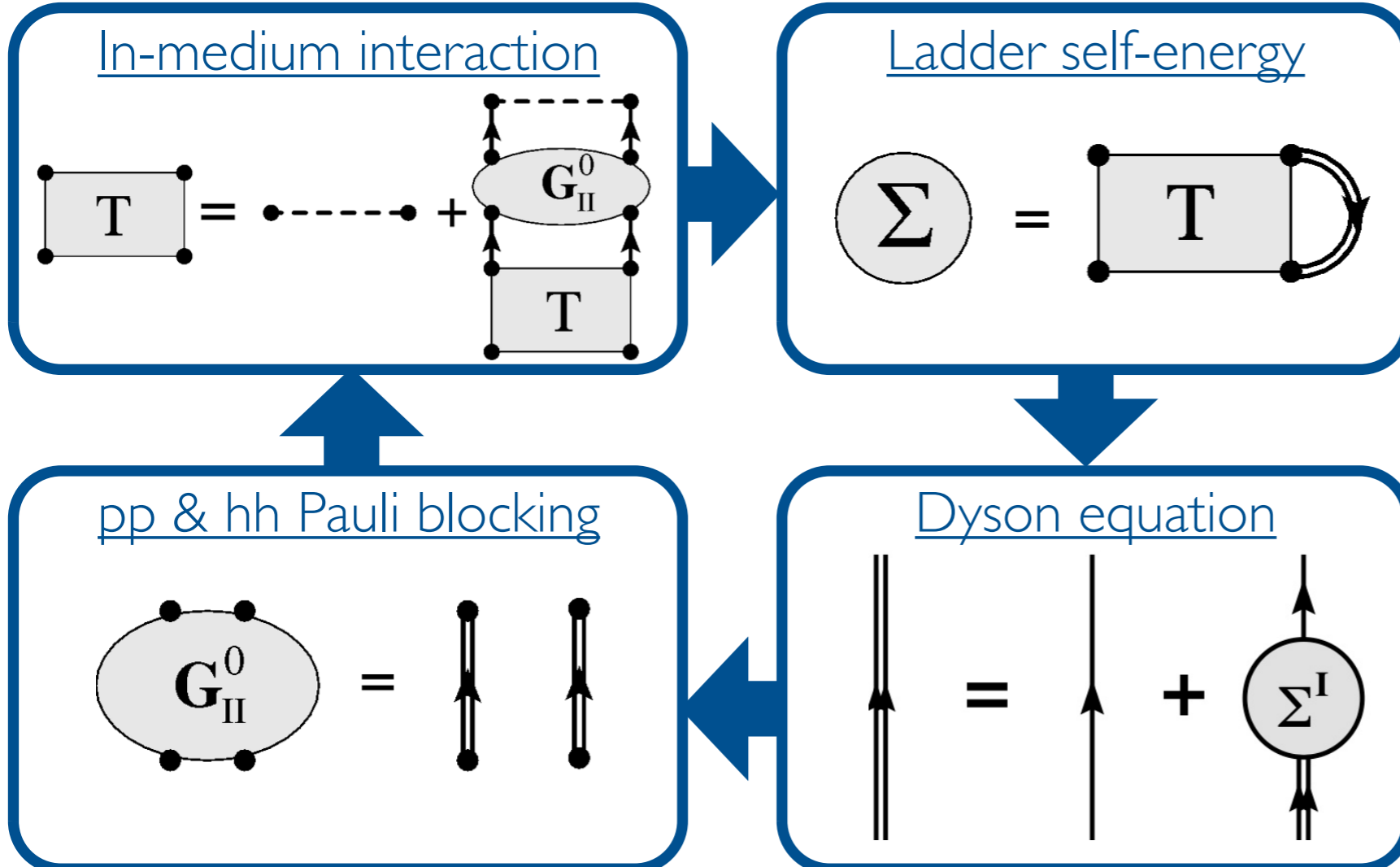
Weinberg, *Phys. Lett. B* **251** 288 (1990), *NPB* **363** 3 (1991)

Entem & Machleidt, *PRC* **68**, 041001(R) (2003)

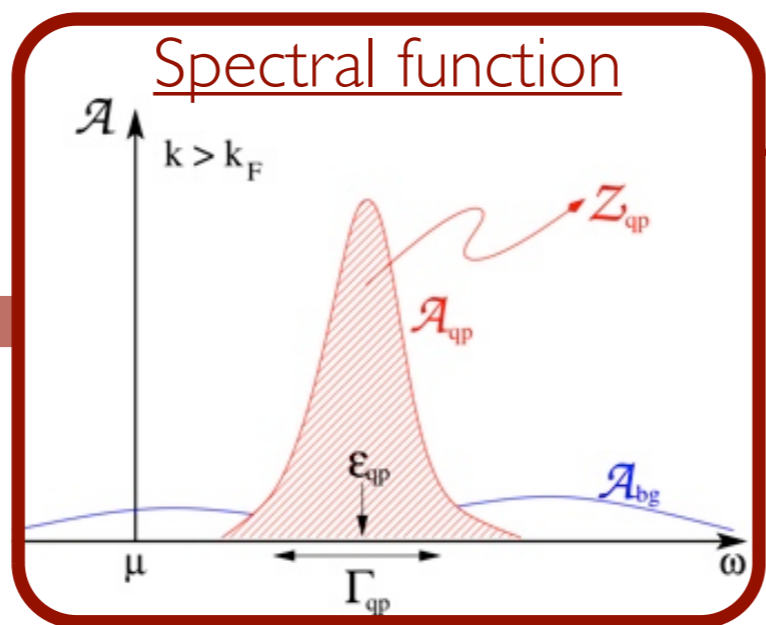
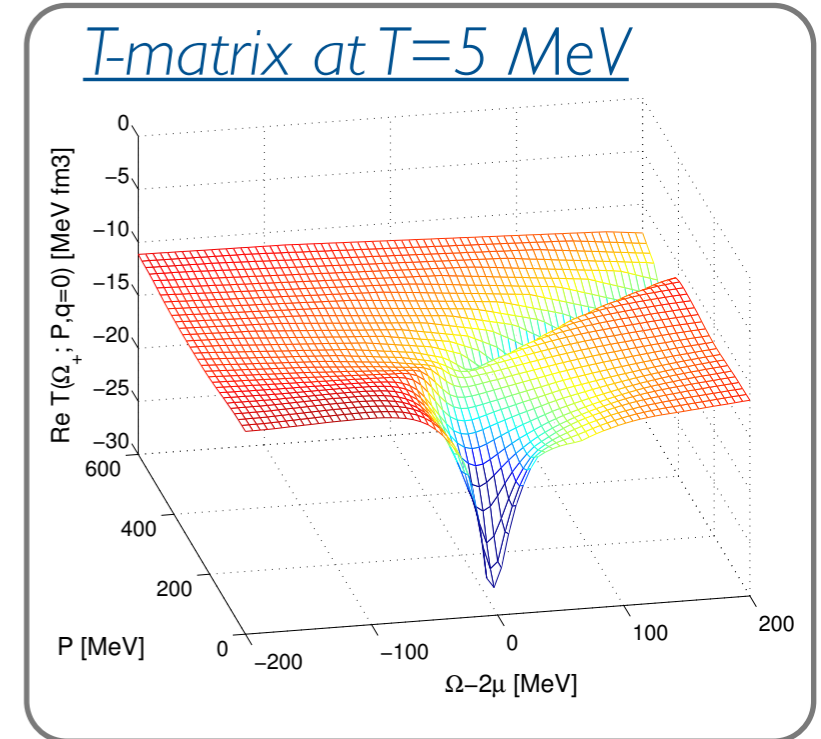
Tews, Schwenk et al., *PRL* **110**, 032504 (2013)

Epelbaum, Friese & Meissner, *PRL* **115**, 122301 (2015) 5

SCGF Ladder approximation



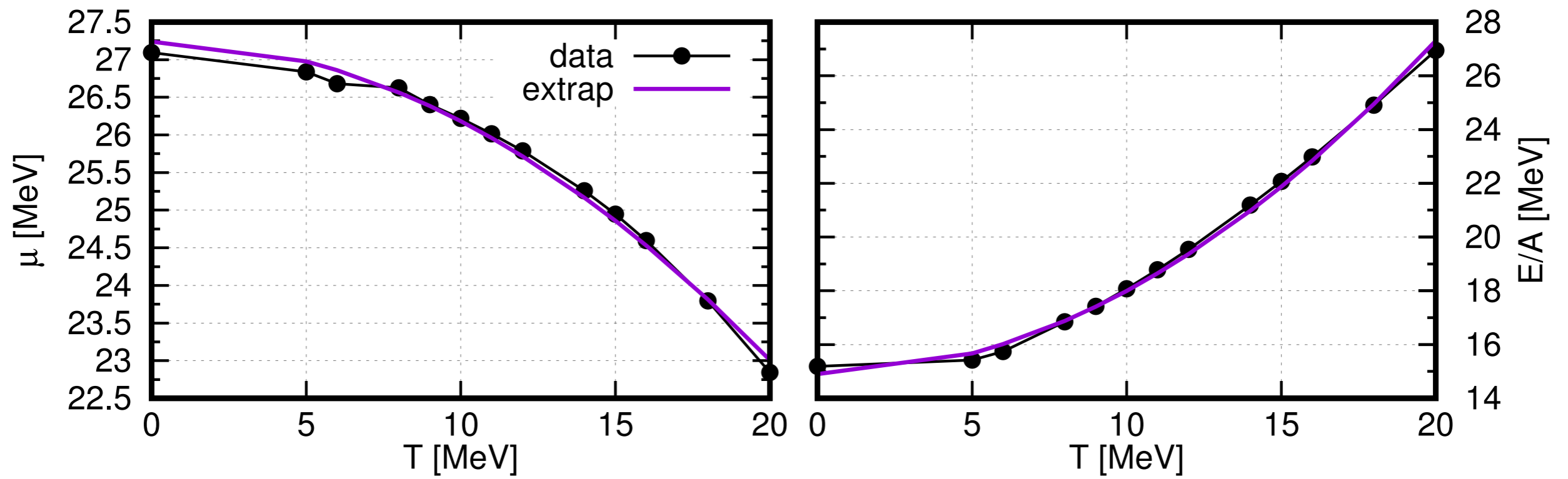
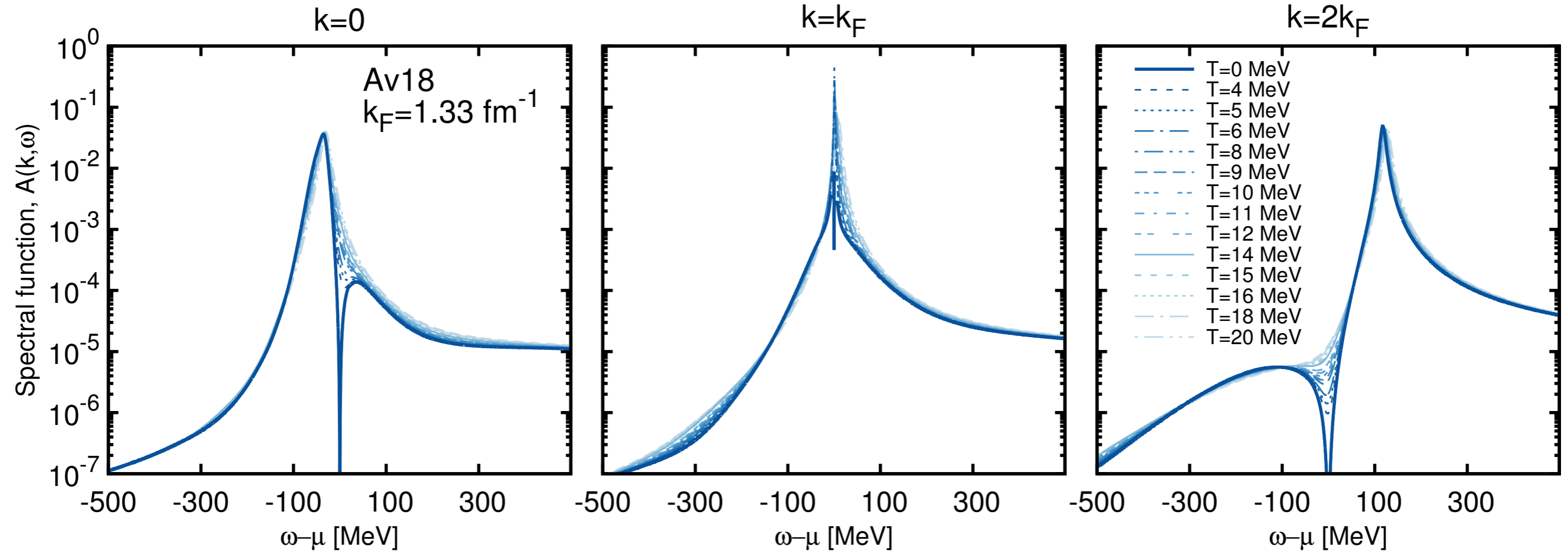
- Self-consistent **resummation**
- **Energy** and **momentum** integral
- @**Finite T** (Matsubara)



One-body properties
Momentum distribution
Thermodynamics & EoS
Transport

Ramos, Polls & Dickhoff, NPA **503** 1 (1989)
 Alm *et al.*, PRC **53** 2181 (1996)
 Dewulf *et al.*, PRL **90** 152501 (2003)
 Frick & Muther, PRC **68** 034310 (2003)
 Rios, PhD Thesis, U. Barcelona (2007)
 Soma & Bozek, PRC **78** 054003 (2008)
 Rios & Soma PRL **108** 012501 (2012)

$T=0$ extrapolations



Momentum distribution

Single-particle occupation

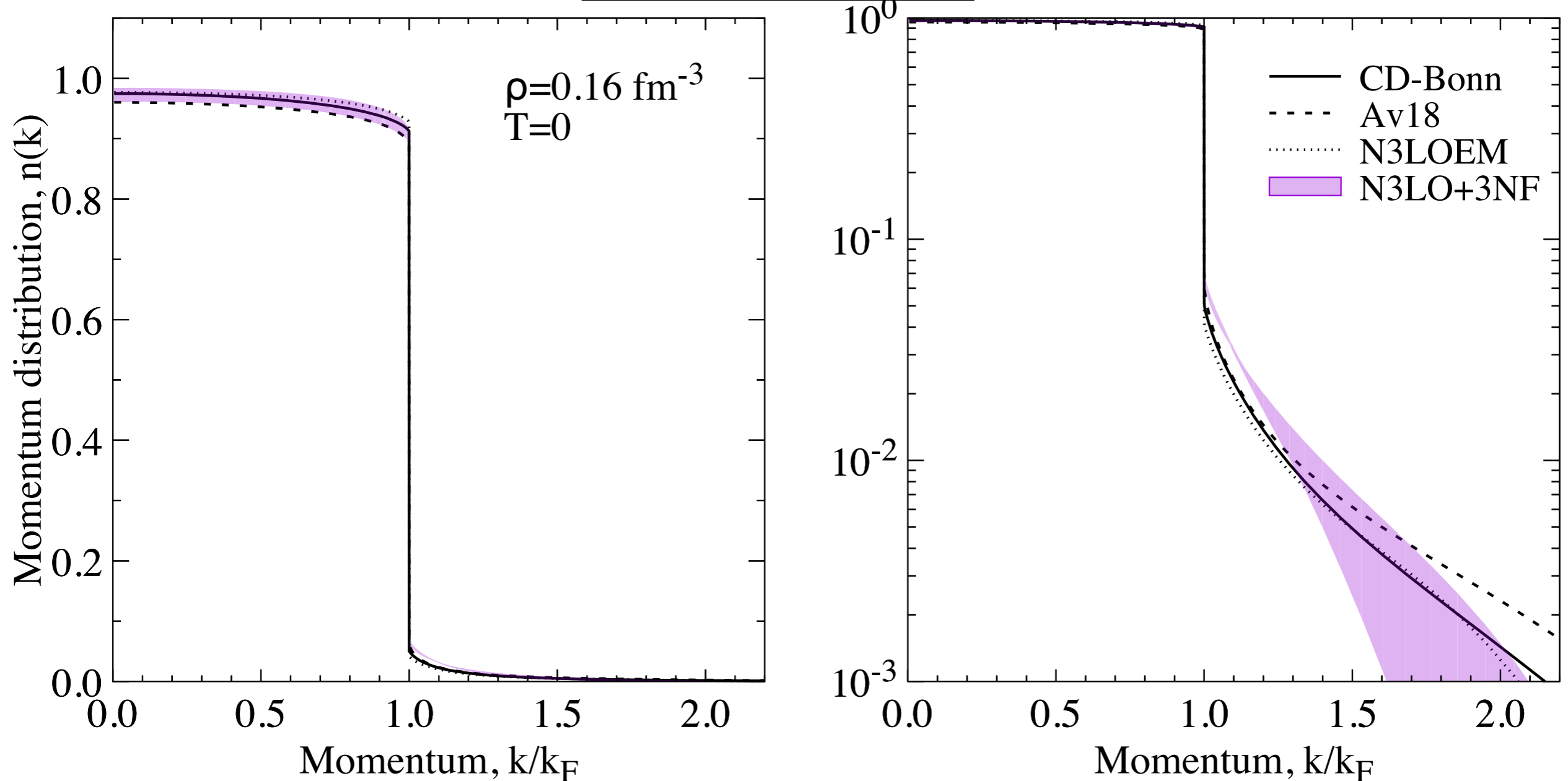
$$n(k) = \langle a_k^\dagger a_k \rangle$$

$$\nu \int \frac{d^3 k}{(2\pi)^3} n(k) = \rho$$



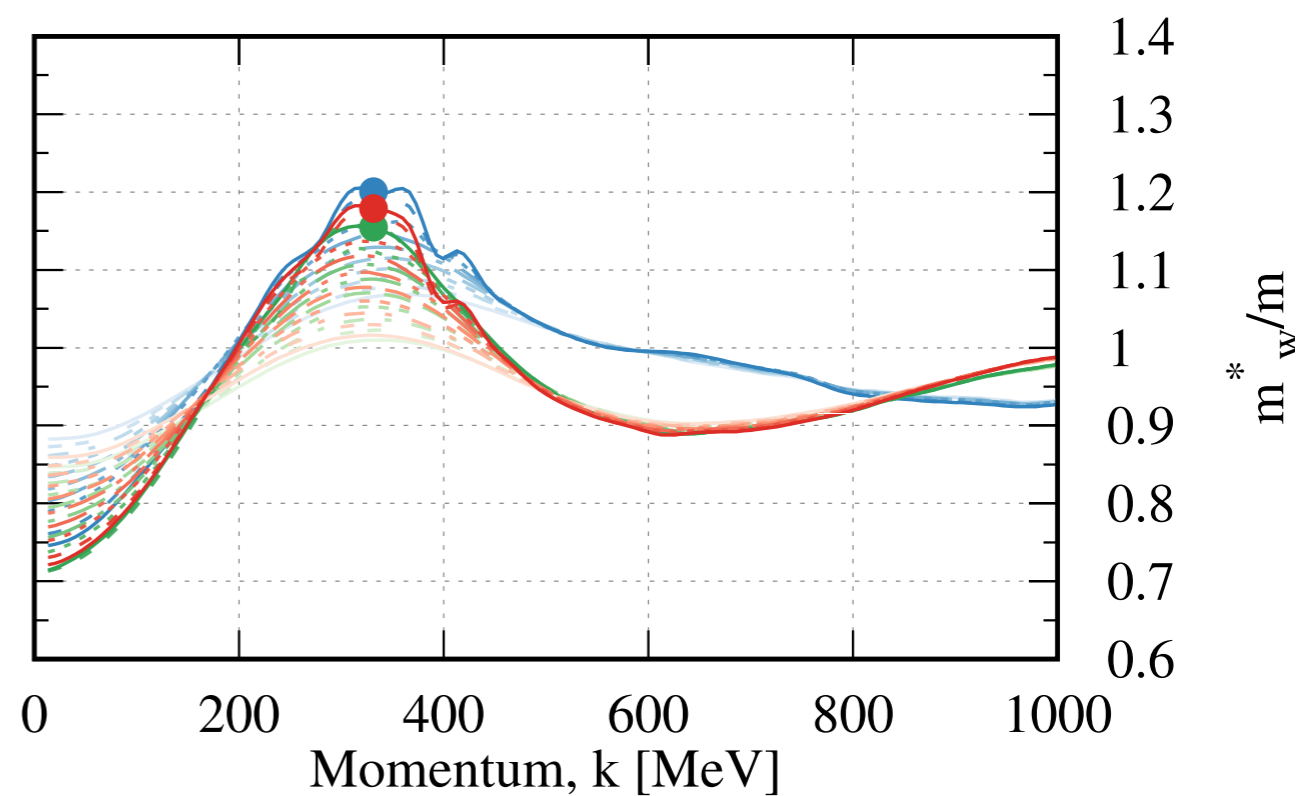
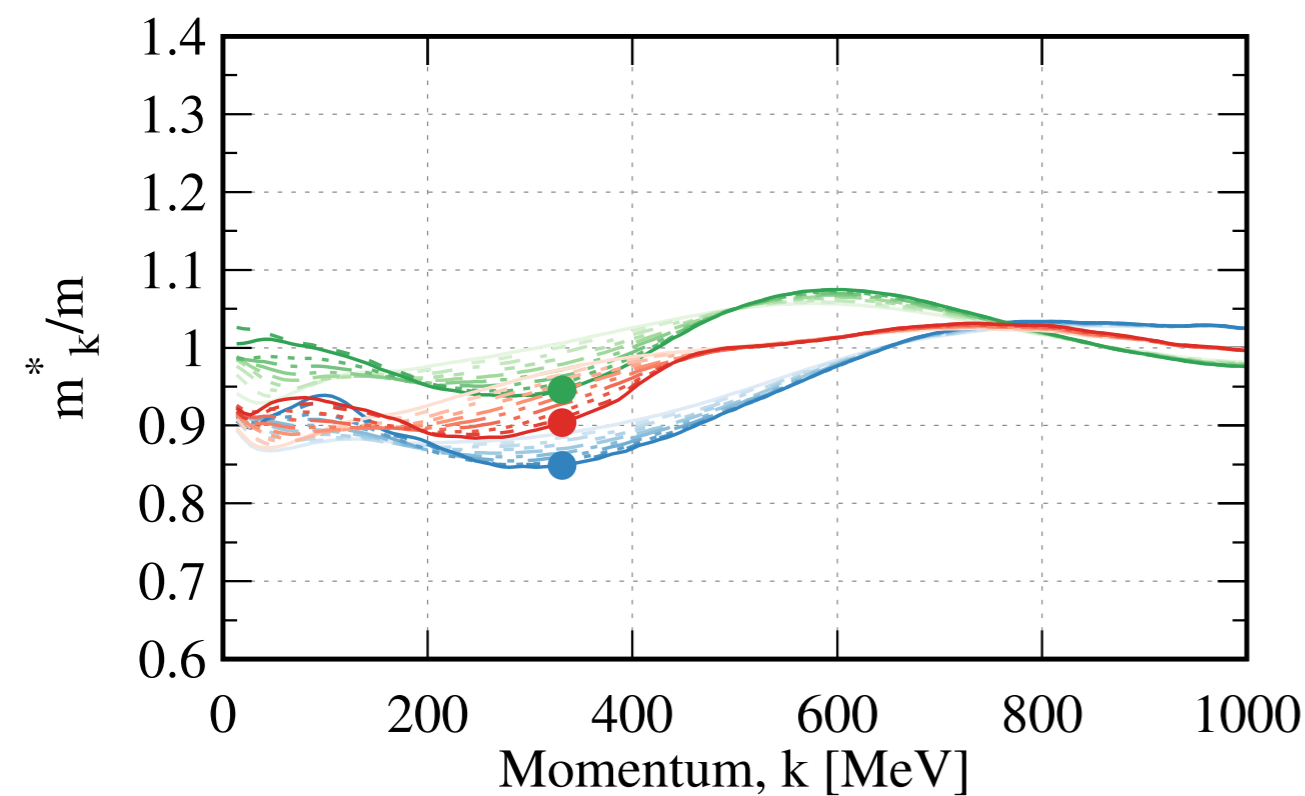
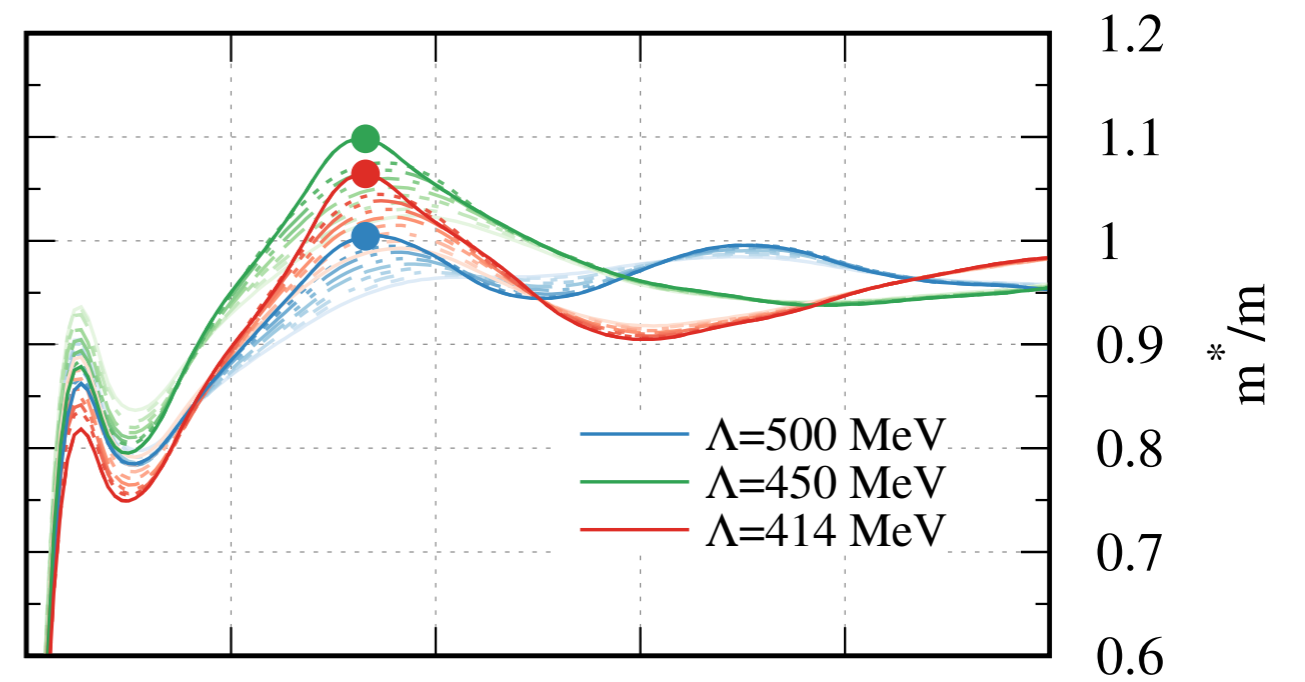
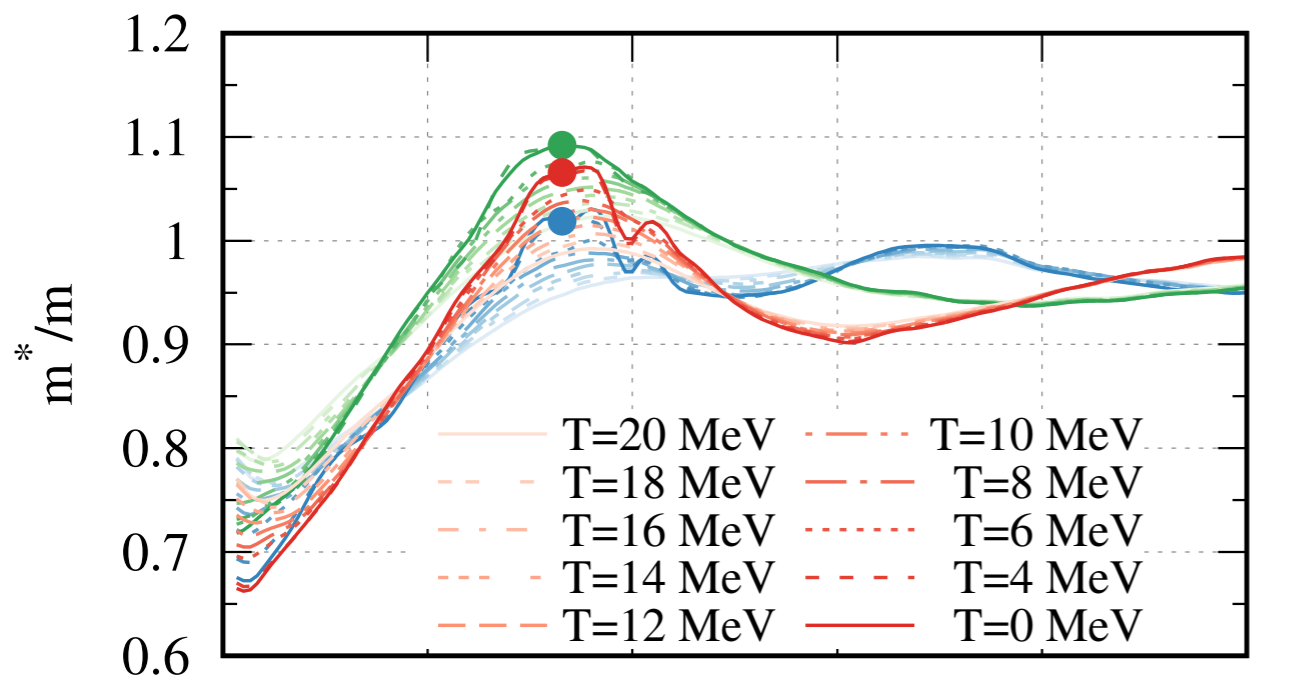
UNIVERSITY OF
SURREY

Neutron matter

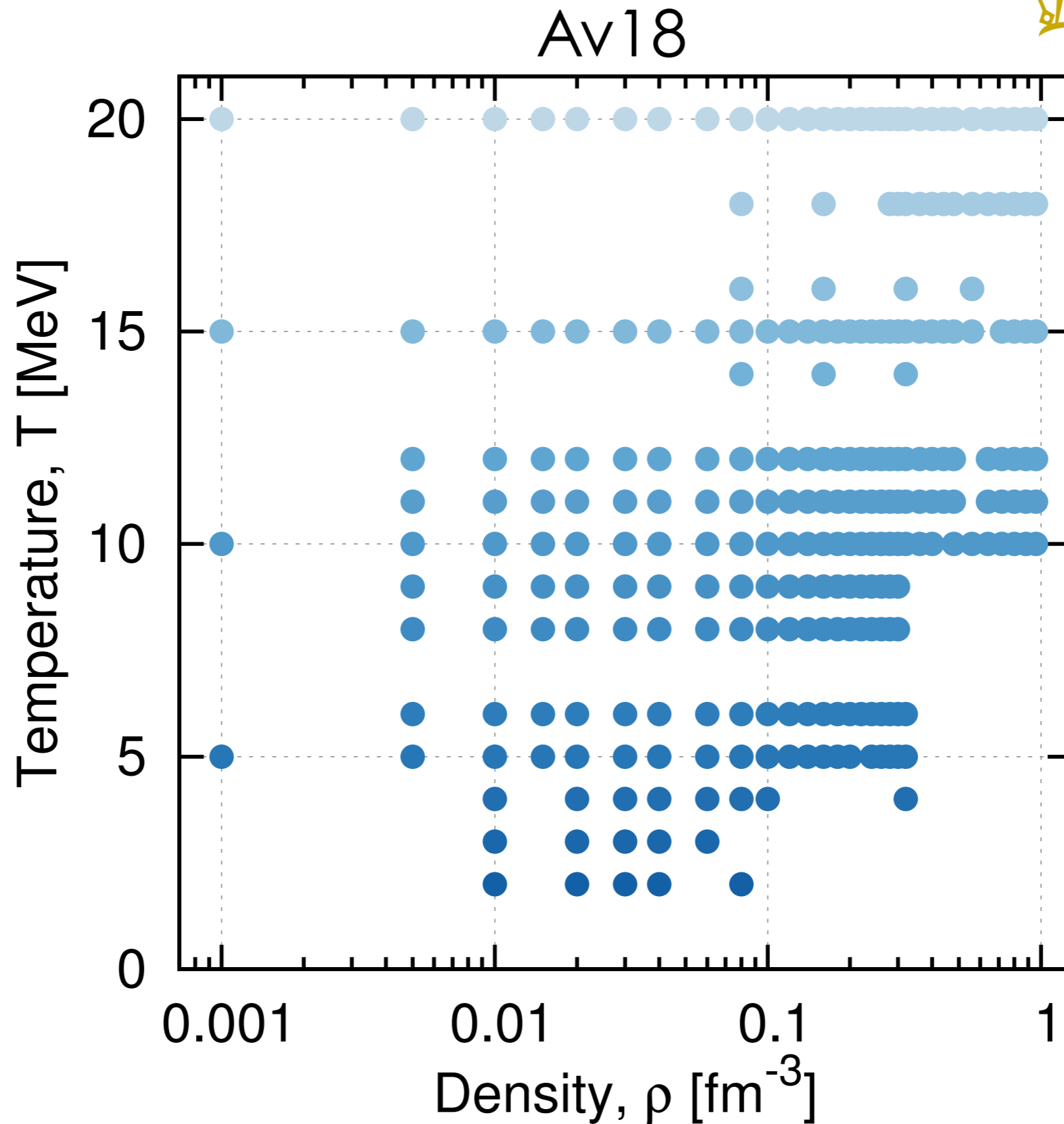


- Dependence on NN interaction under control
- PNM: 4-5% depletion at low k
- N3LO+3NF = N3LO 2NF + N2LO 3NF @ $\Lambda = 414-500 \text{ MeV}$ (cutoff variation only)

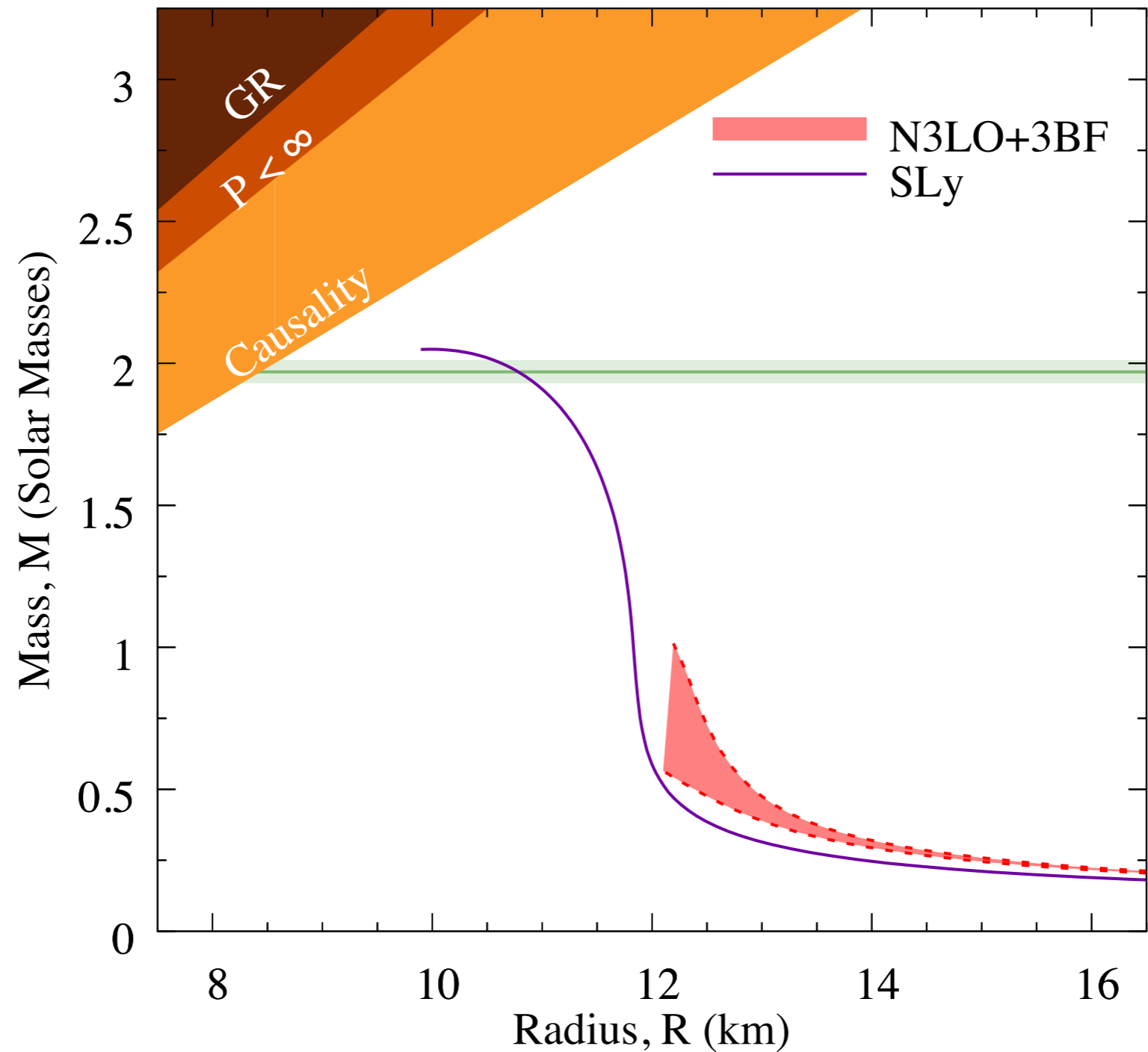
Neutron matter $\rho=0.16 \text{ fm}^{-3}$



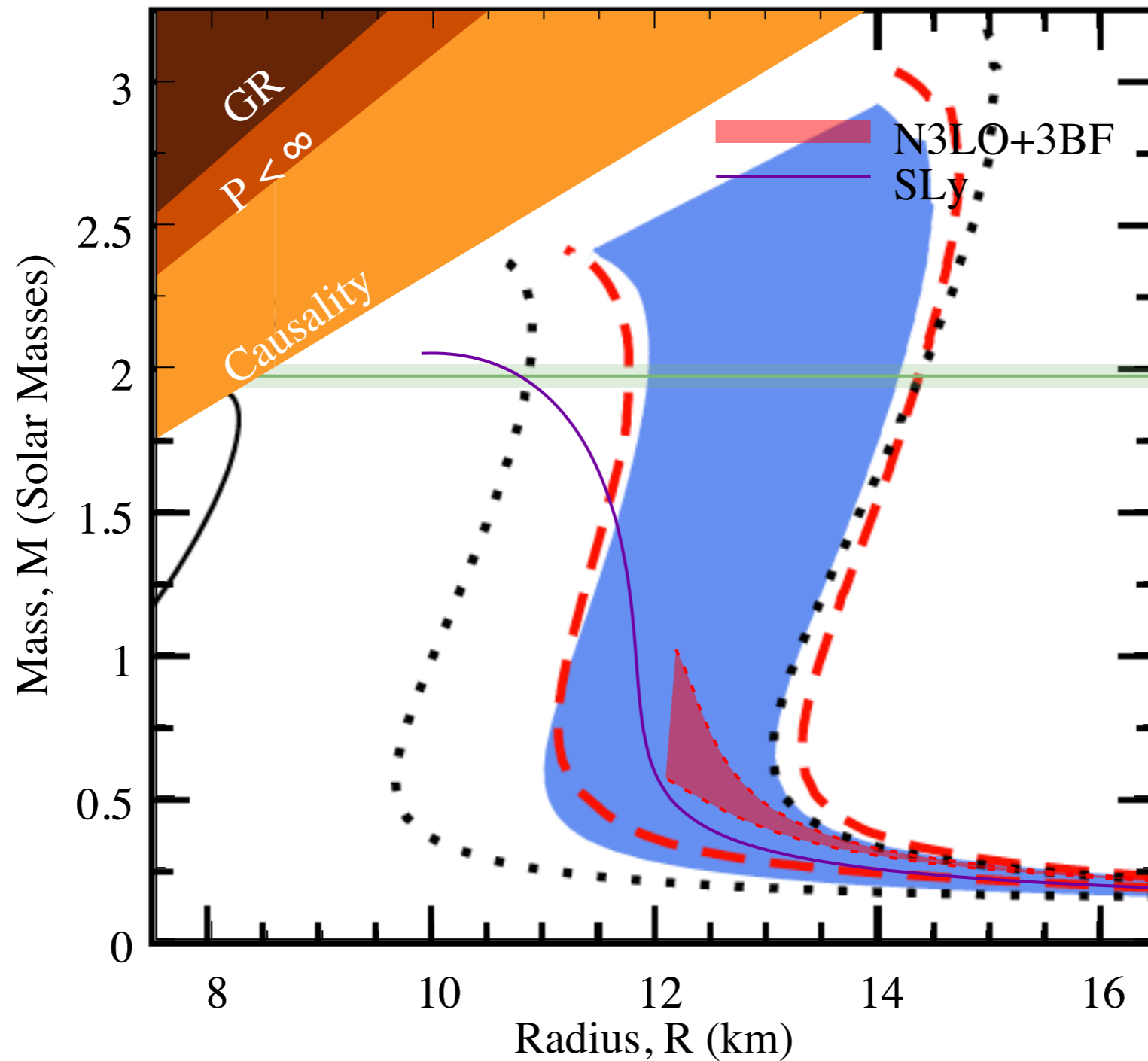
Coraggio, Holt, et al. PRC **89** 044321 (2014)



Self-energy, spectral function & thermodynamics



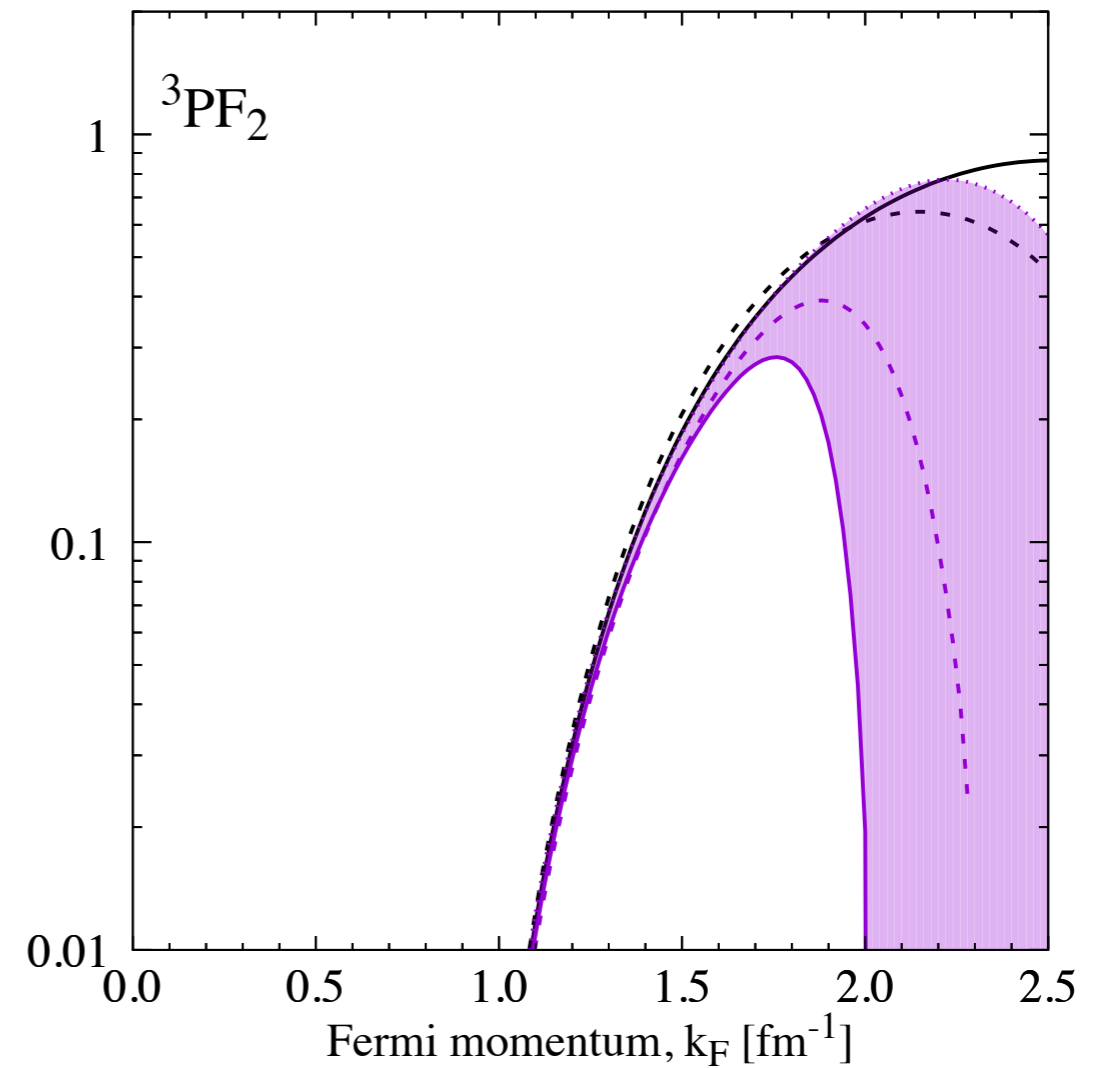
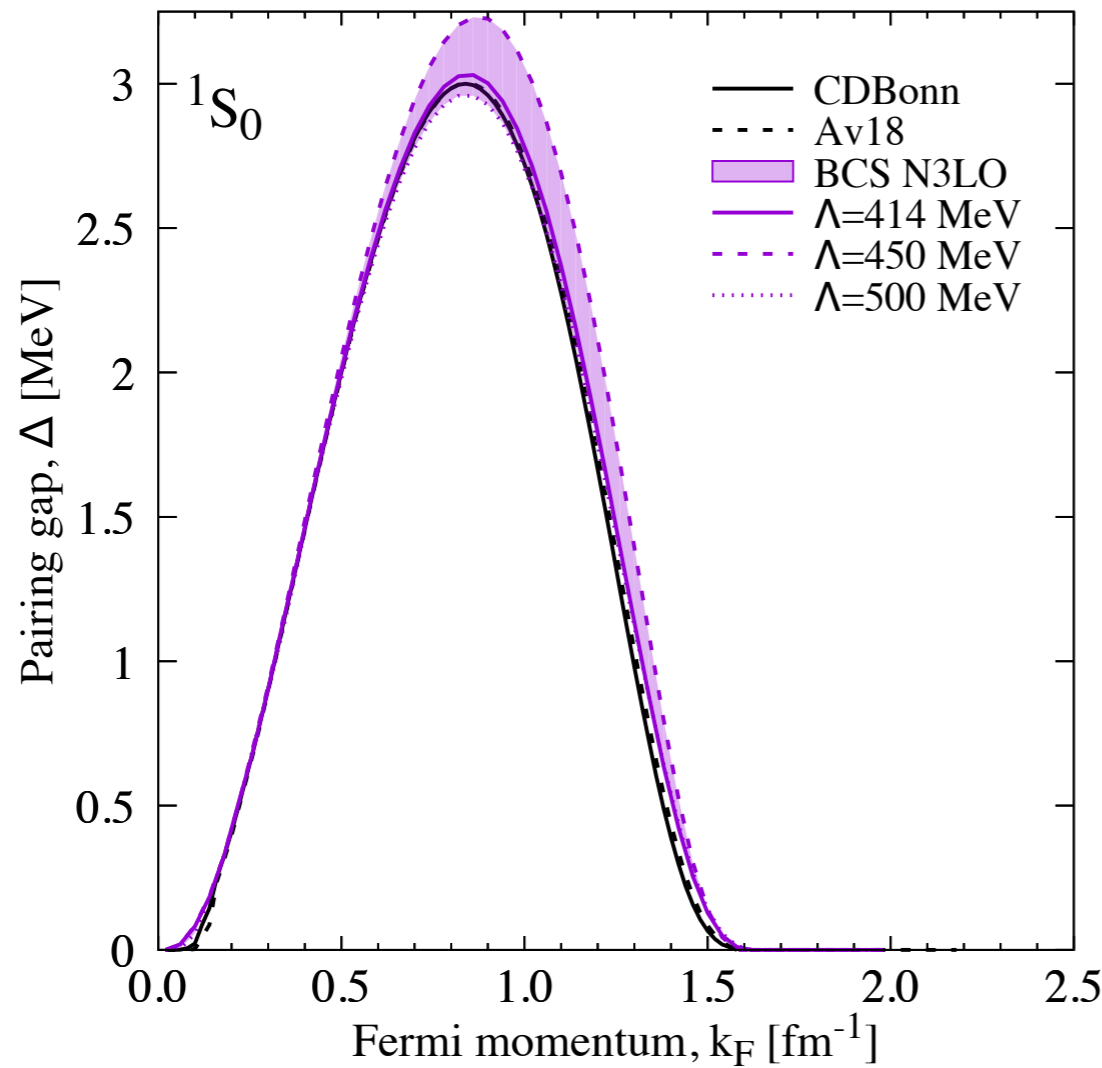
- Mass-Radius relation from SCGF calculations
- Cut-off variation (N3LO) and/or SRG evolution



Hebeler, Lattimer, Pethick, Schwenk *ApJ* **773** 11 (2013)

- Mass-Radius relation from SCGF calculations
- Cut-off variation (N3LO) and/or SRG evolution

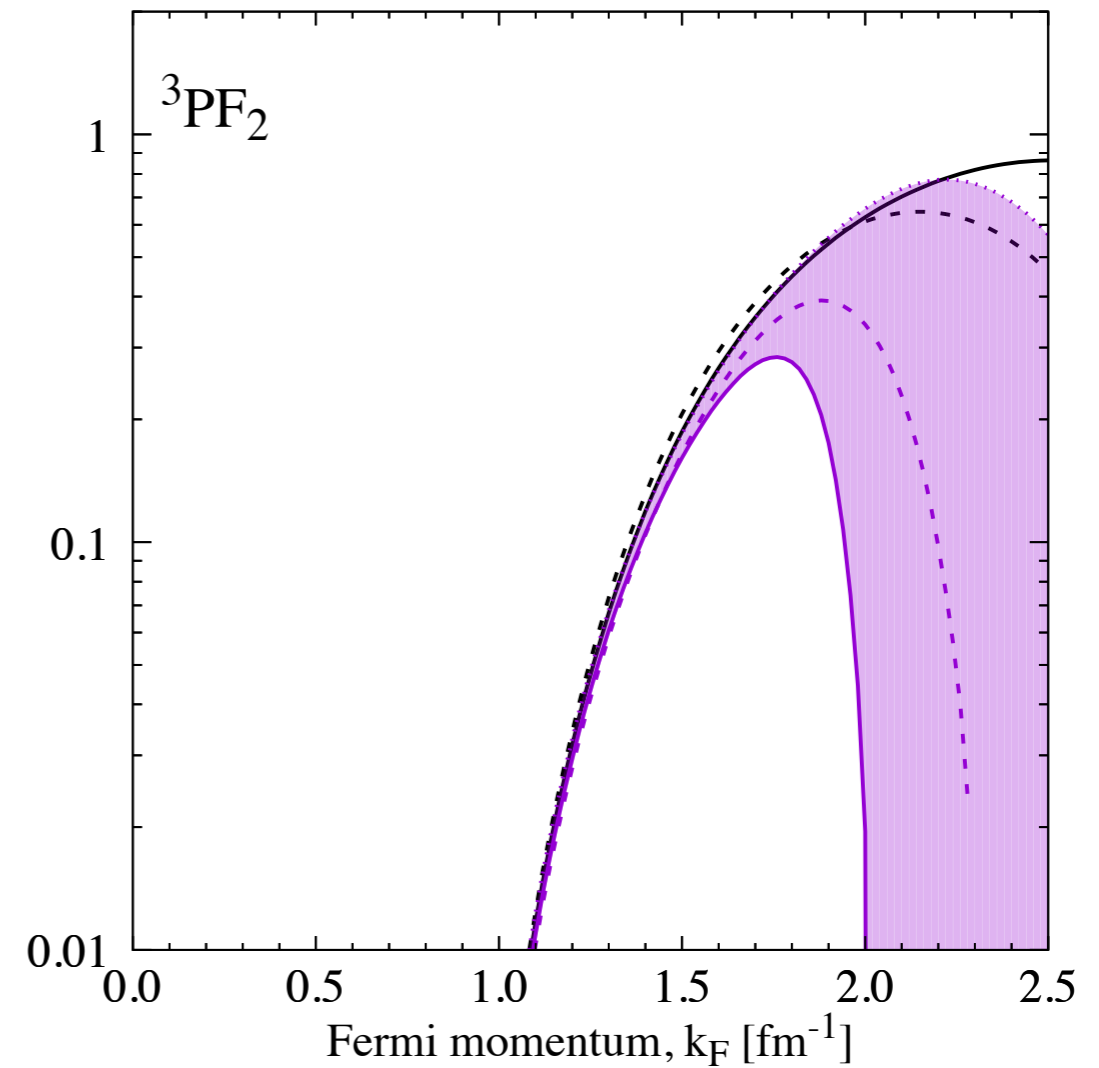
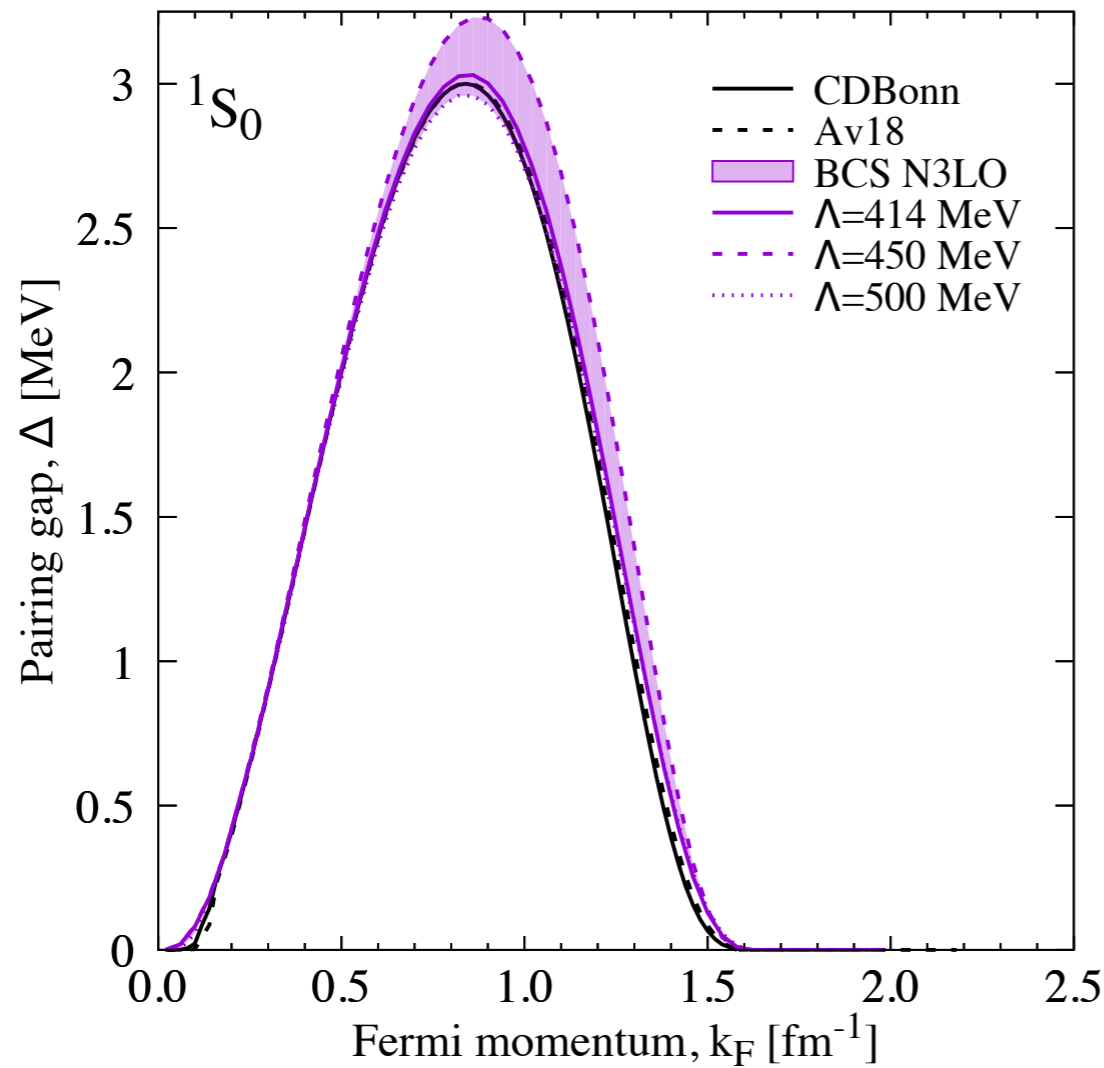
1. Neutron star motivation
- 2. Infinite matter BCS**
3. Beyond-BCS with SCGF methods



BCS equation

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\chi_{k'}^2 + |\Delta_{k'}^{L'}|^2}} \Delta_{k'}^{L'} + \chi_k = \varepsilon_k - \mu$$

- Single-particle spectrum choice: $\varepsilon_k = \frac{k^2}{2m} + U(k) - \mu$
- Angular gap dependence: $|\Delta_k|^2 = \sum_L |\Delta_k^L|^2$

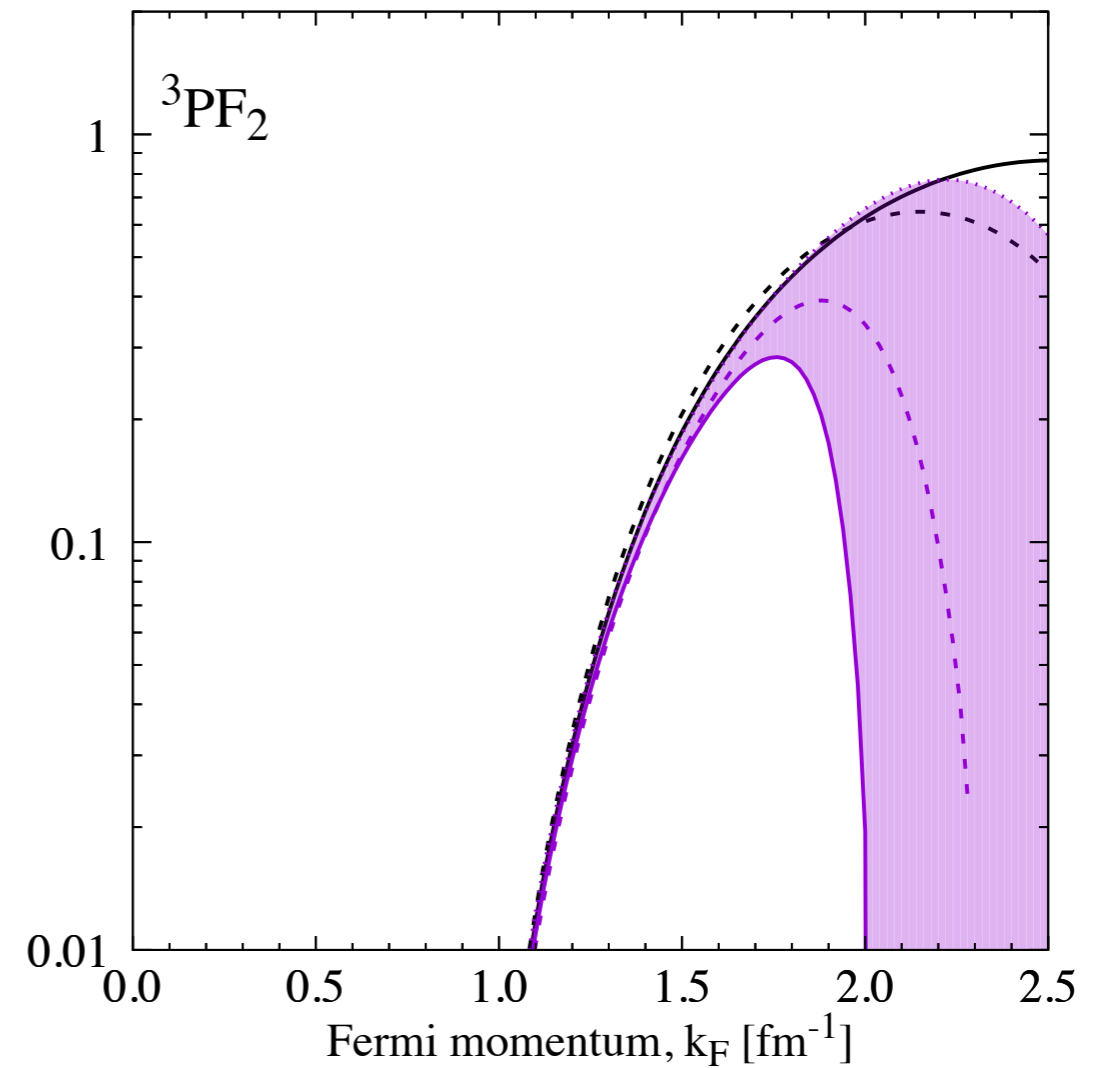
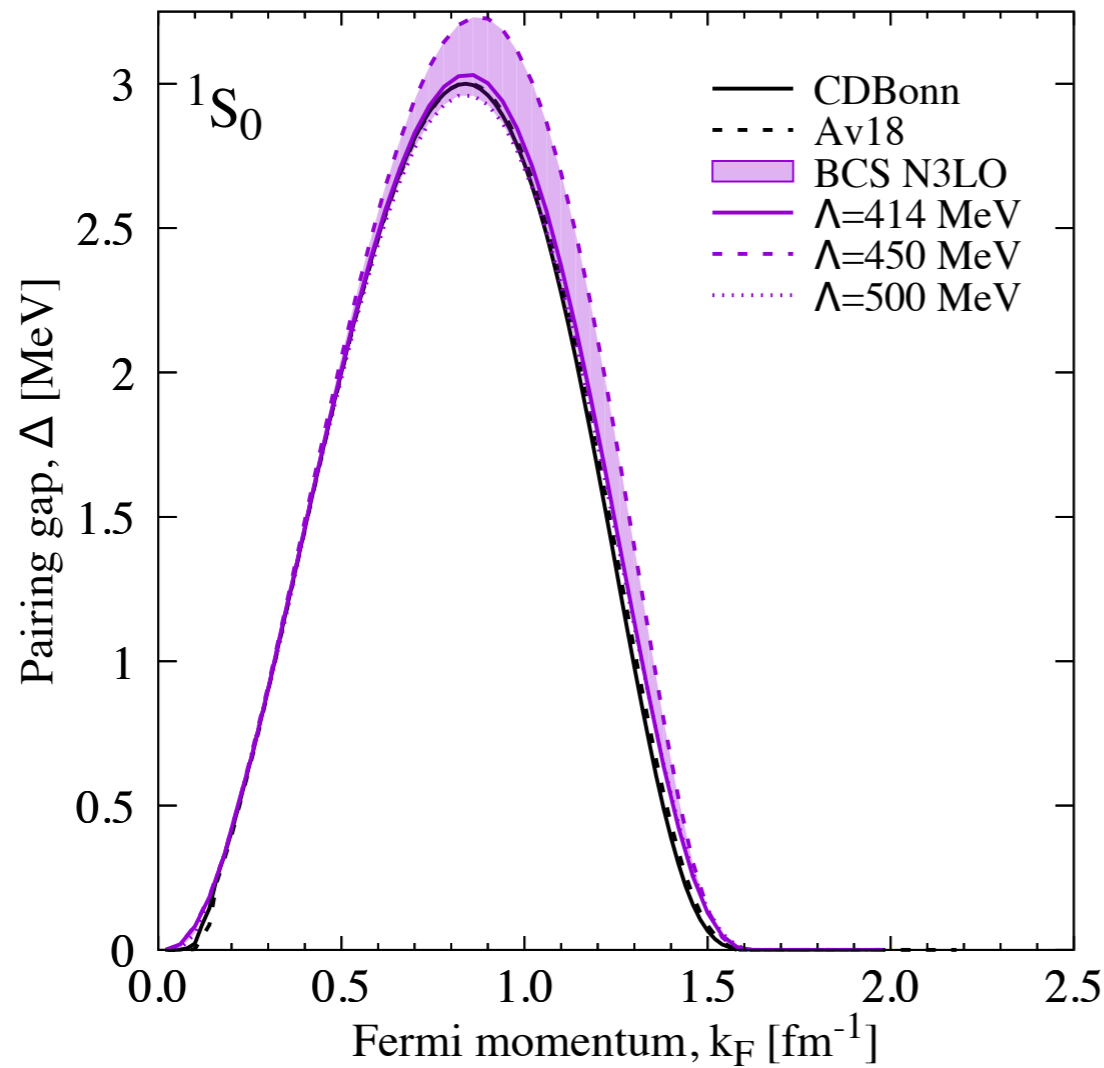


BCS equation

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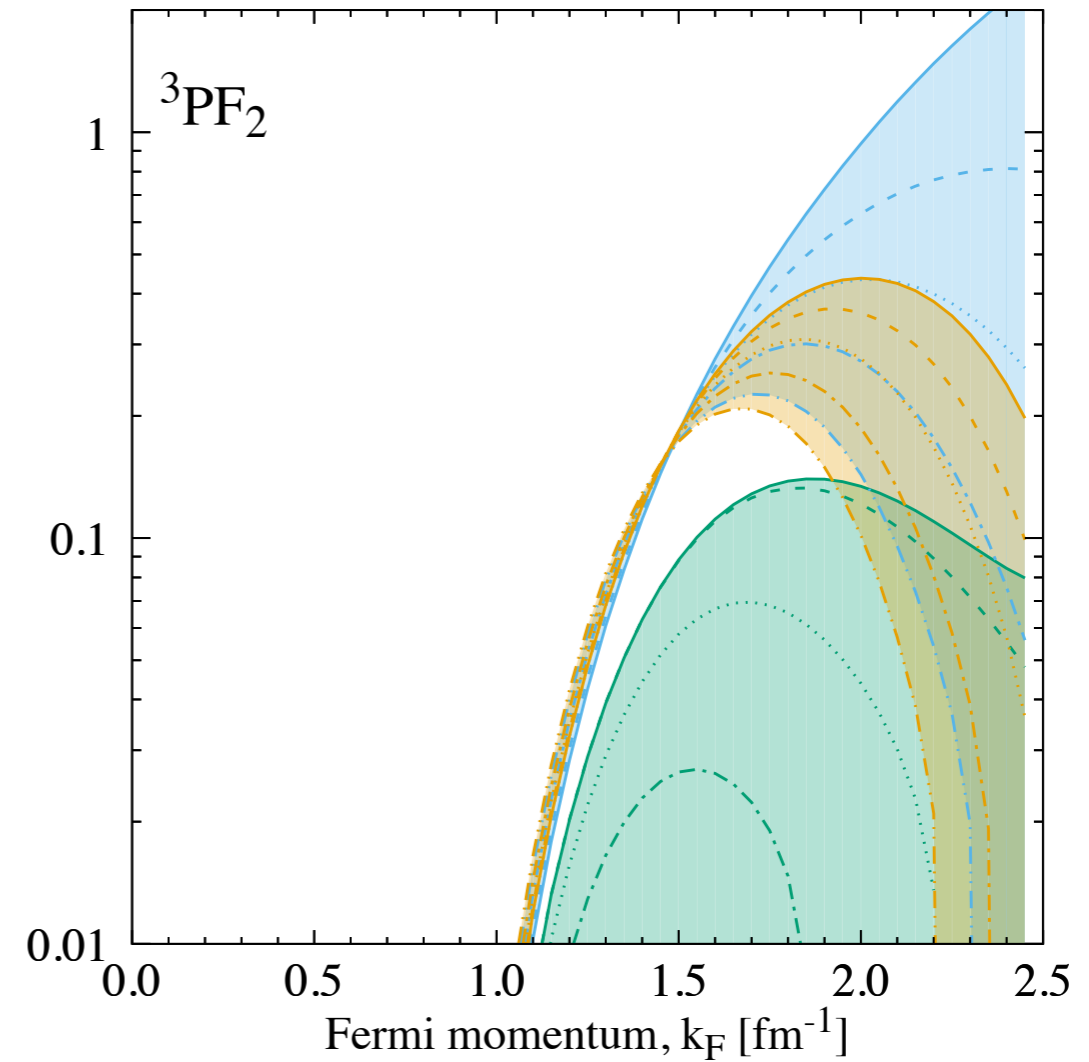
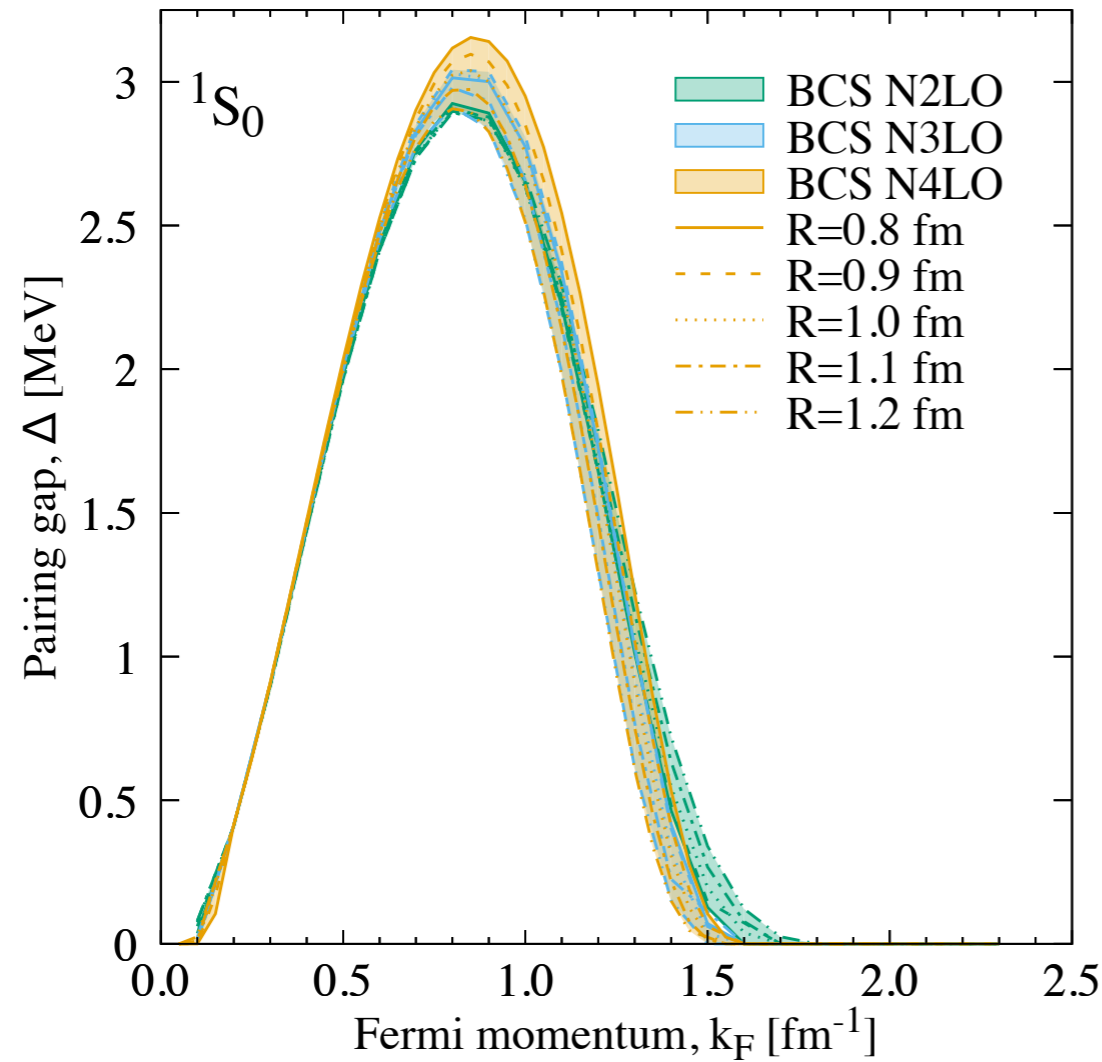
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BCS equation

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\chi_{k'}^2 + |\Delta_{k'}^{L'}|^2}} \Delta_{k'}^{L'} + \chi_k = \varepsilon_k - \mu$$

- Single-particle spectrum choice: $\varepsilon_k = \frac{k^2}{2m} + \cancel{U(k)} - \mu$
- Angular gap dependence: $|\Delta_k|^2 = \sum_L |\Delta_k^L|^2 \approx |\Delta_k^L|^2$

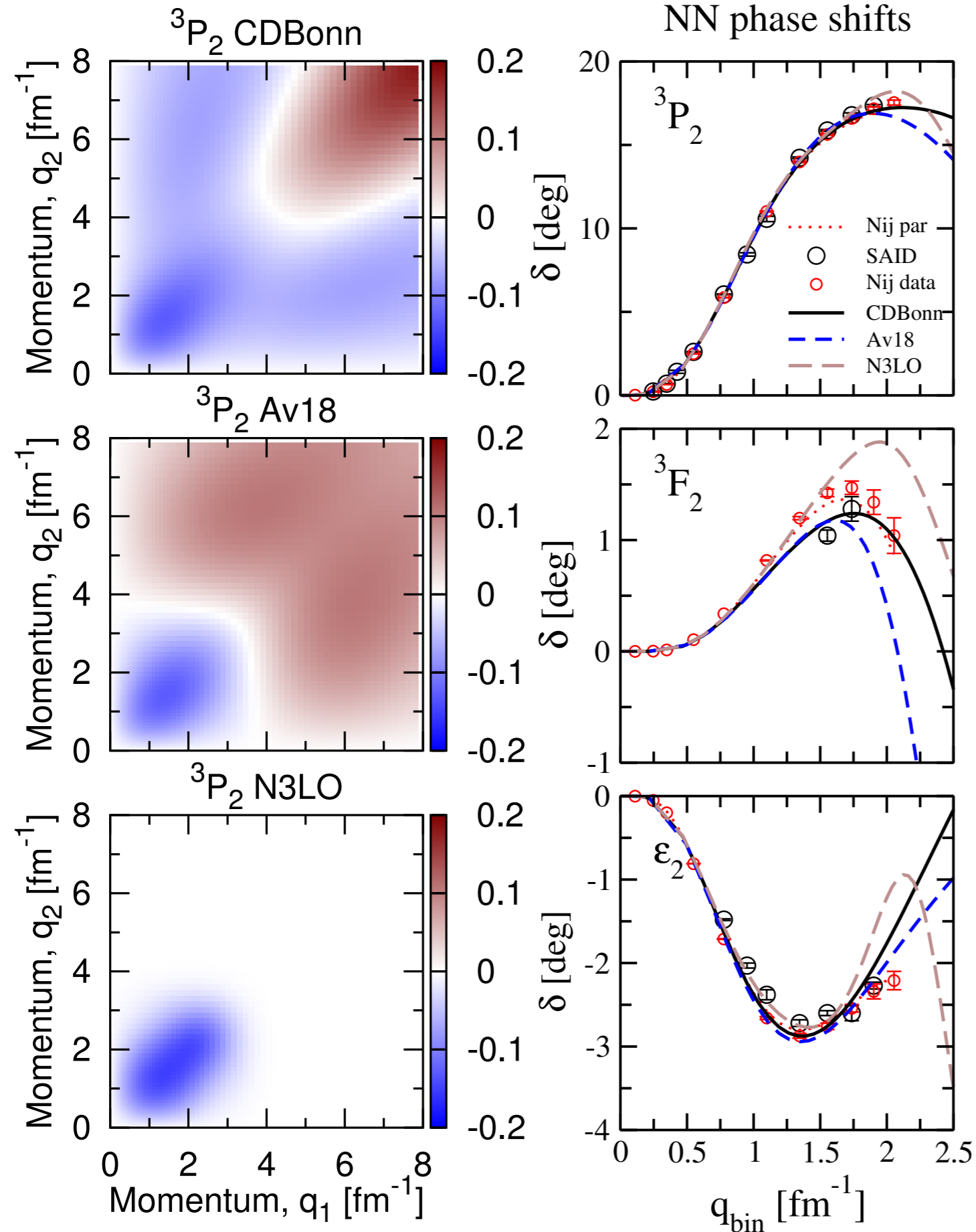


BCS equation

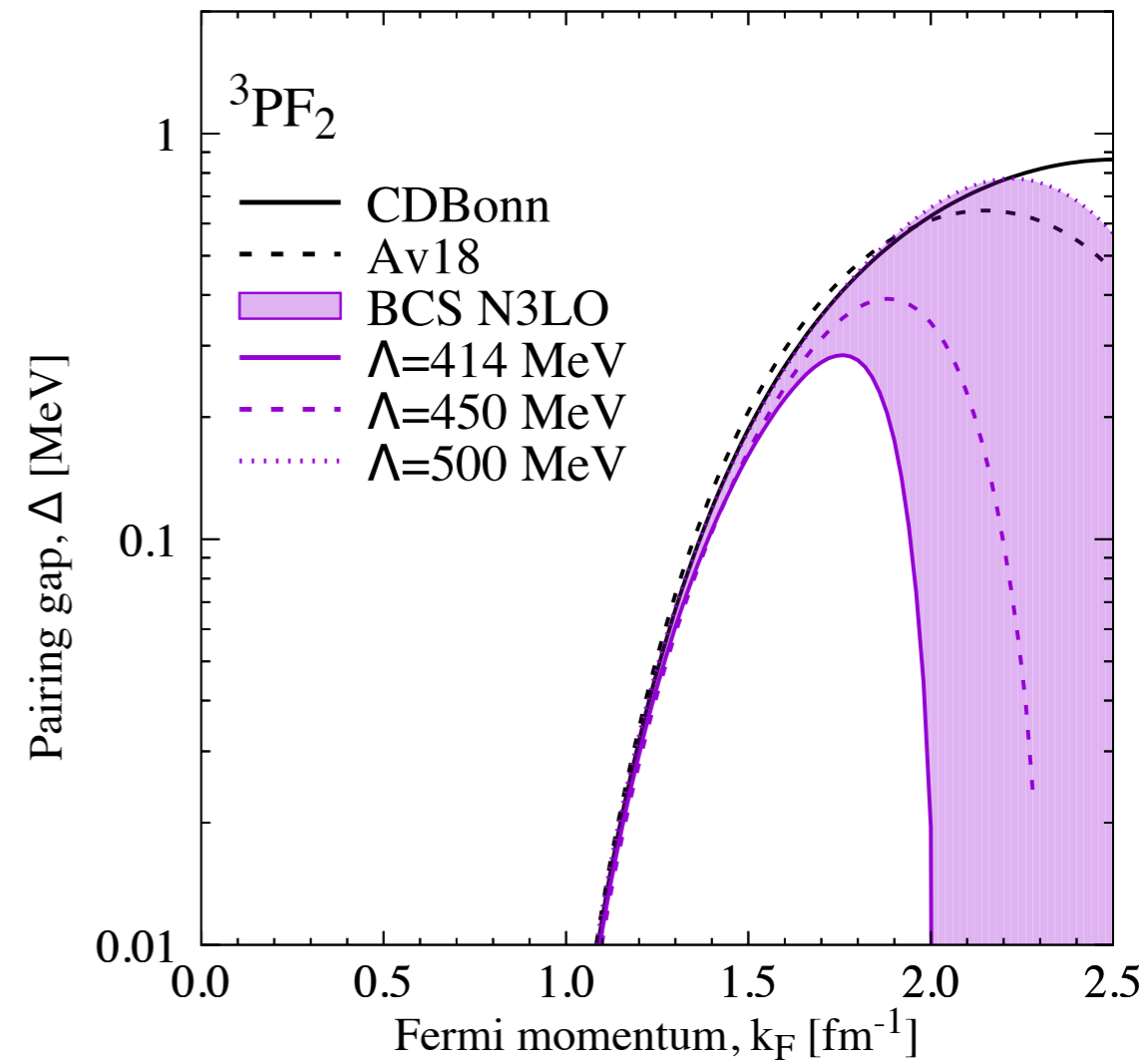
$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\chi_{k'}^2 + |\Delta_{k'}^{L'}|^2}} \Delta_{k'}^{L'} + \chi_k = \varepsilon_k - \mu$$

- Single-particle spectrum choice: $\varepsilon_k = \frac{k^2}{2m} + \cancel{U(k)} - \mu$
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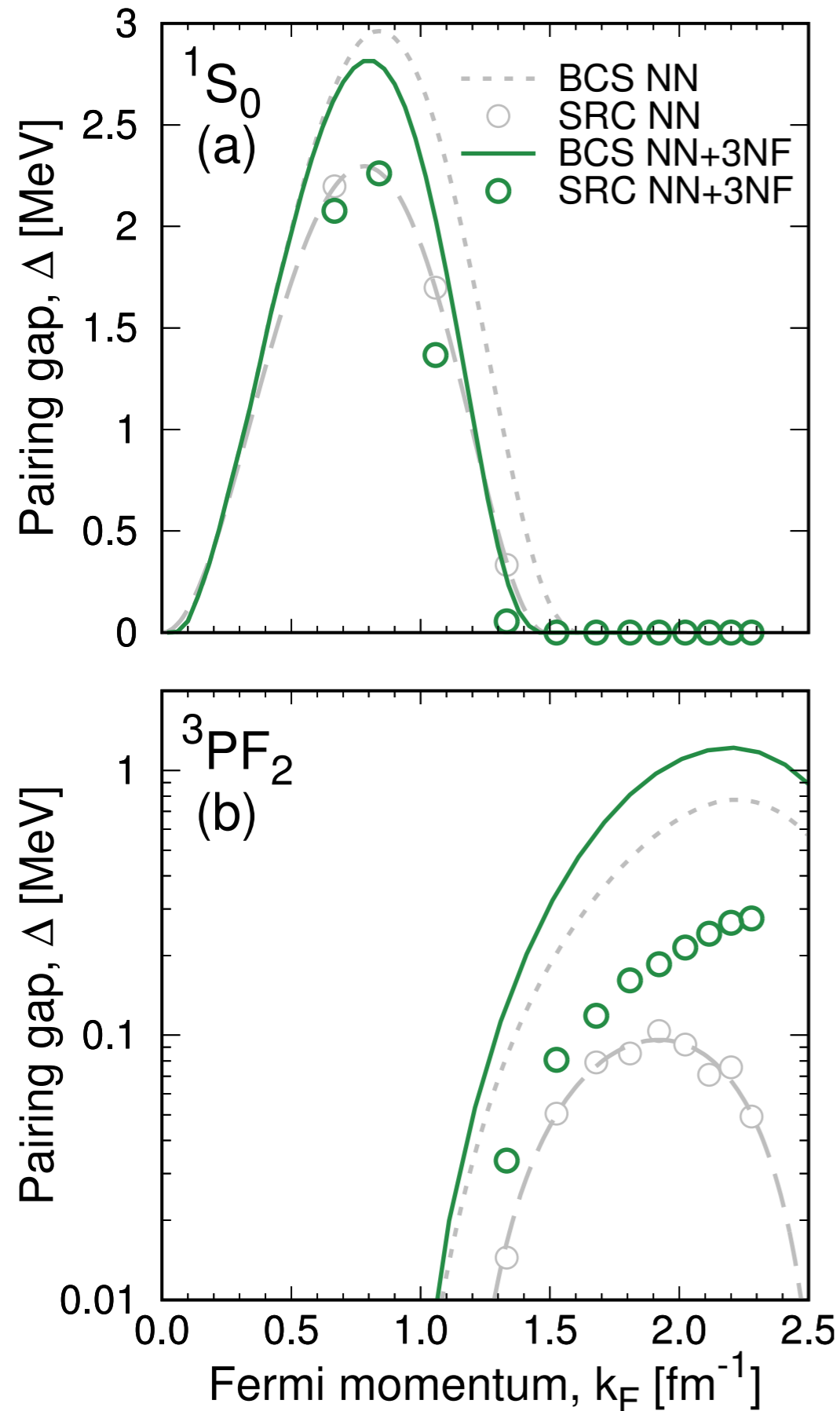
3PF_2 pairing: phase shift equivalence



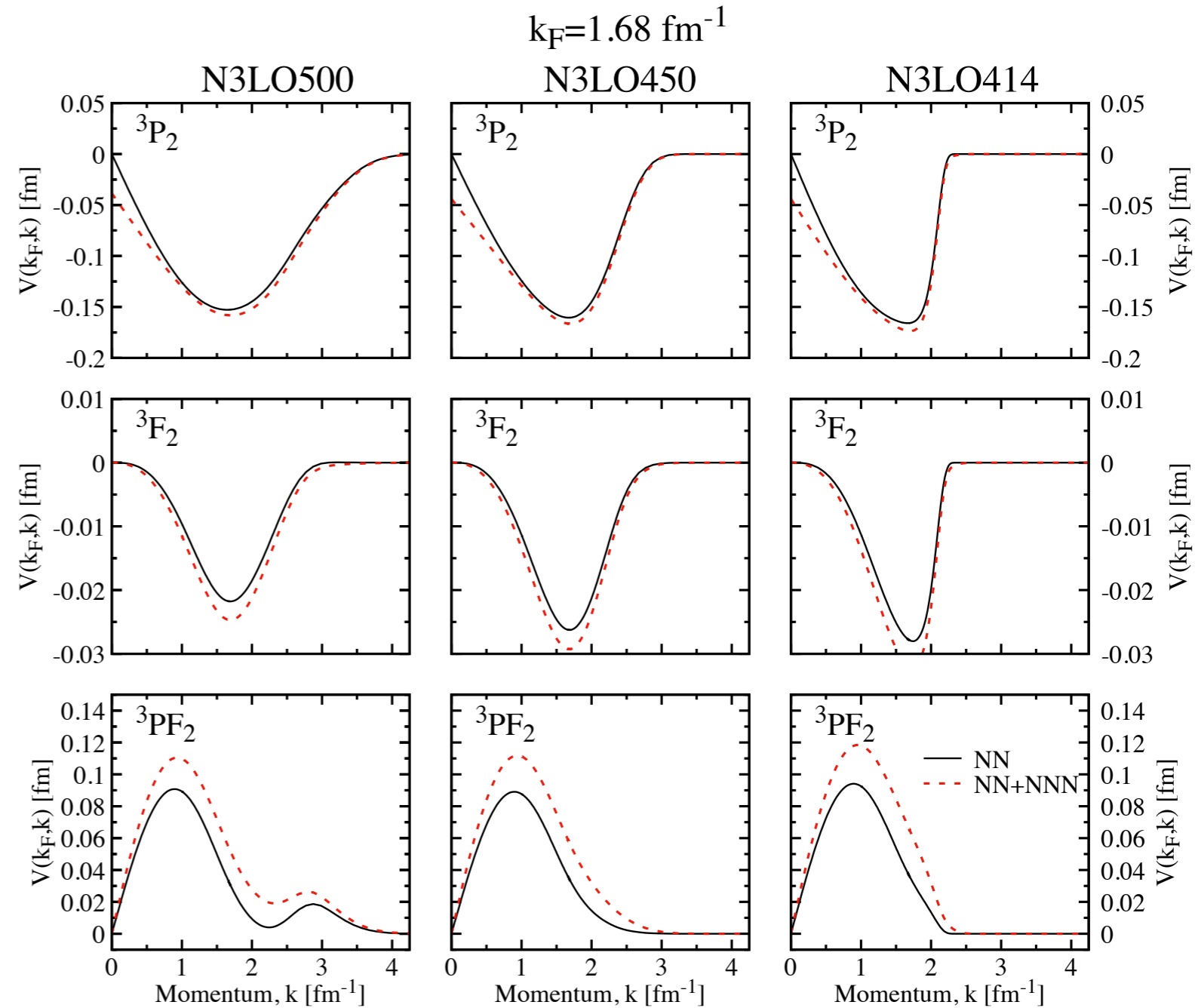
Neutron matter BCS gaps



3BF effect: estimate



Effective two-body force \Rightarrow NN forces



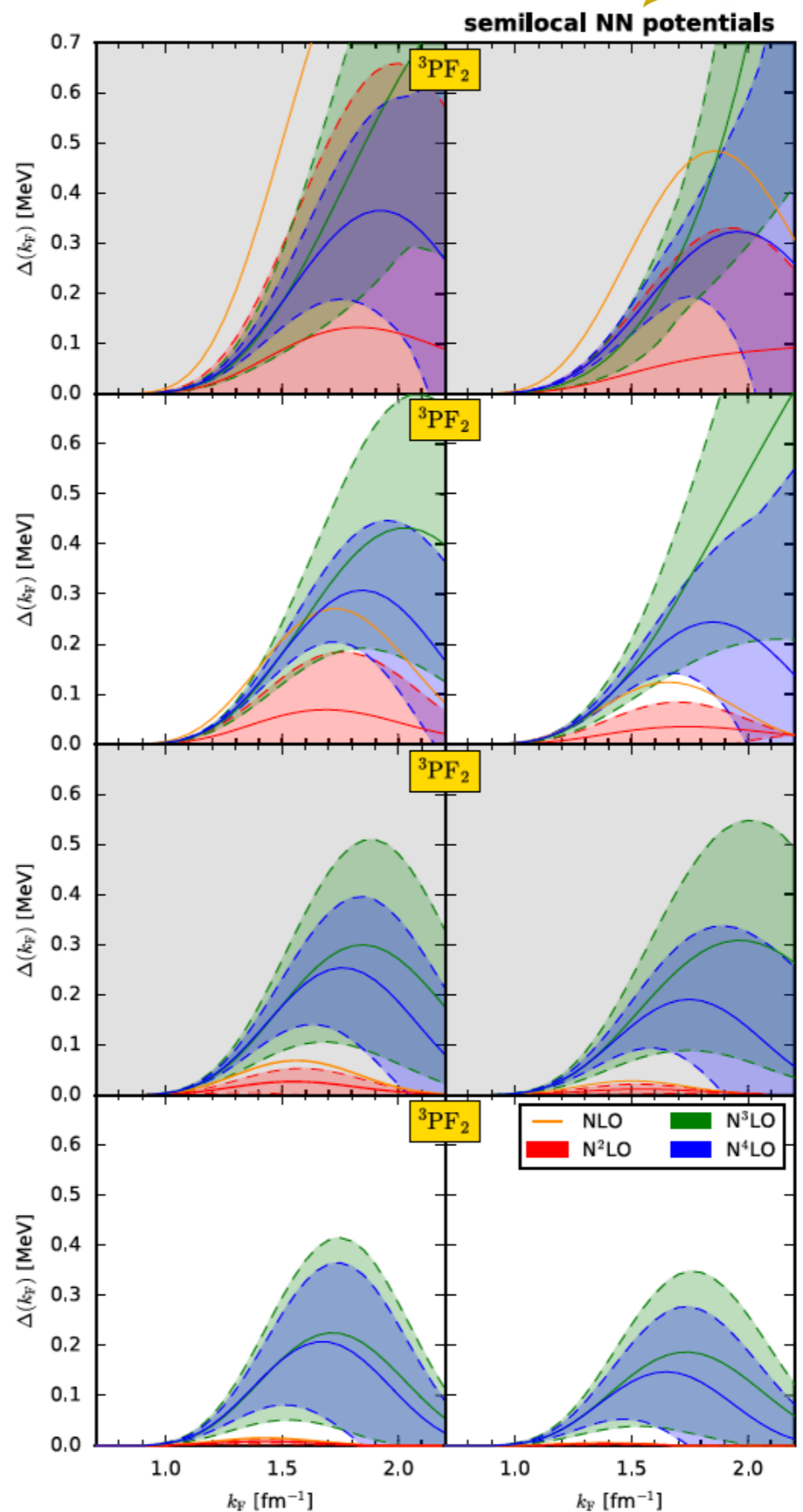
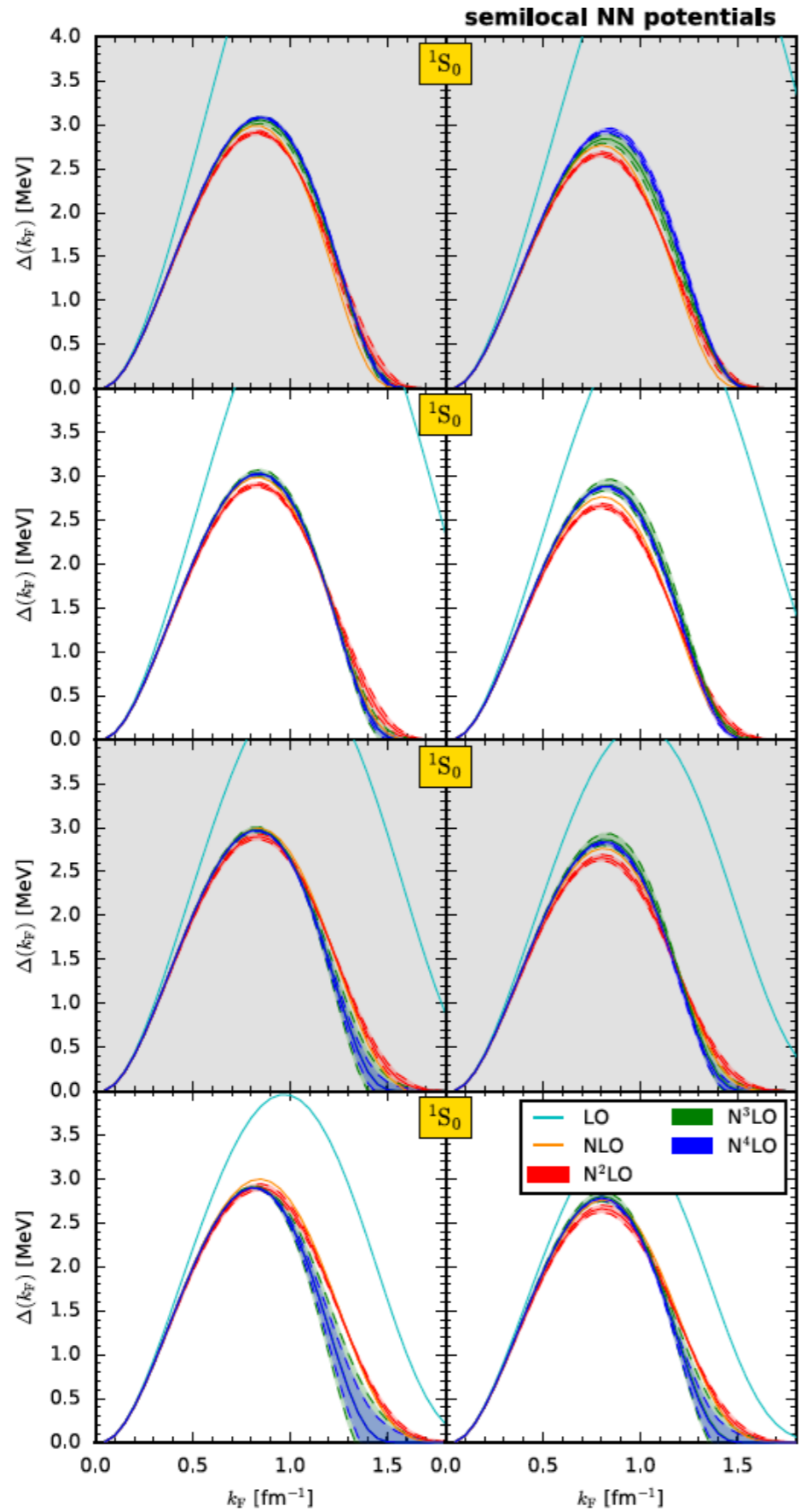
- Singlet gap: 3NF **reduce** closure
- Triplet gap: 3NF **increase** gap

Uncertainties at the BCS + HF level

“Hard”



“Soft”



R=0.9 fm

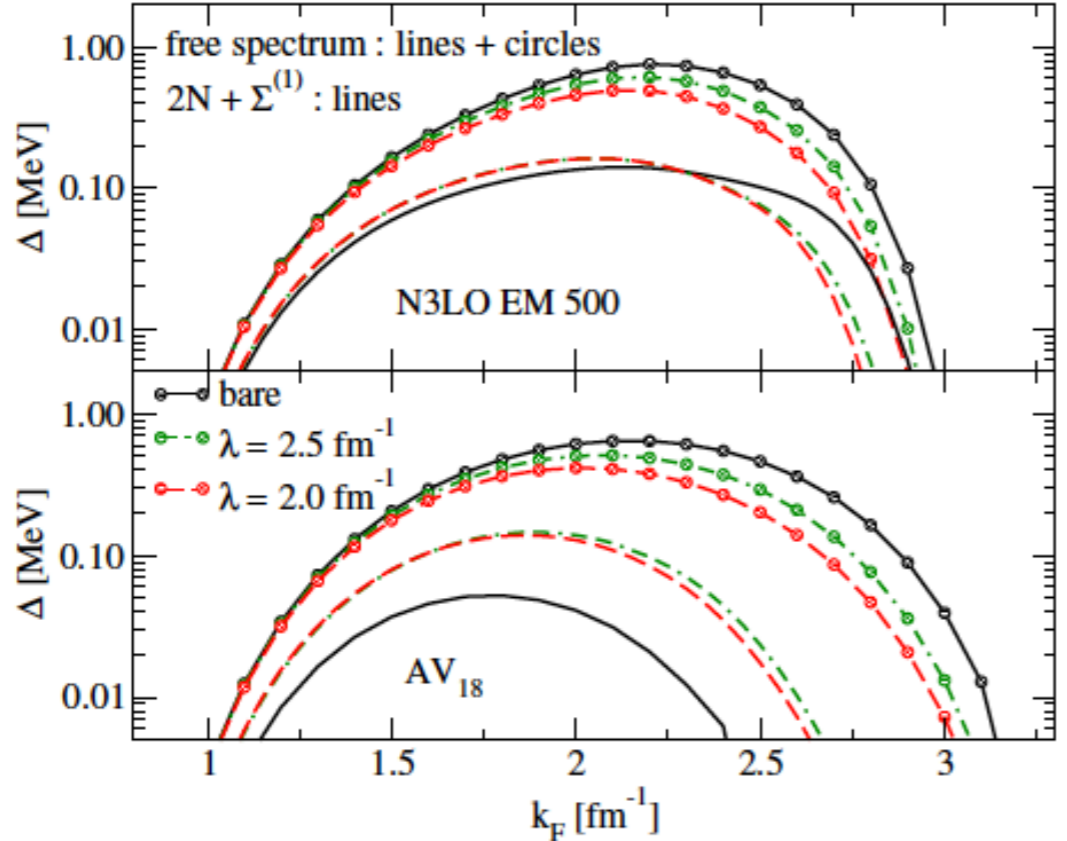
R=1.0 fm

R=1.1 fm

R=1.2 fm

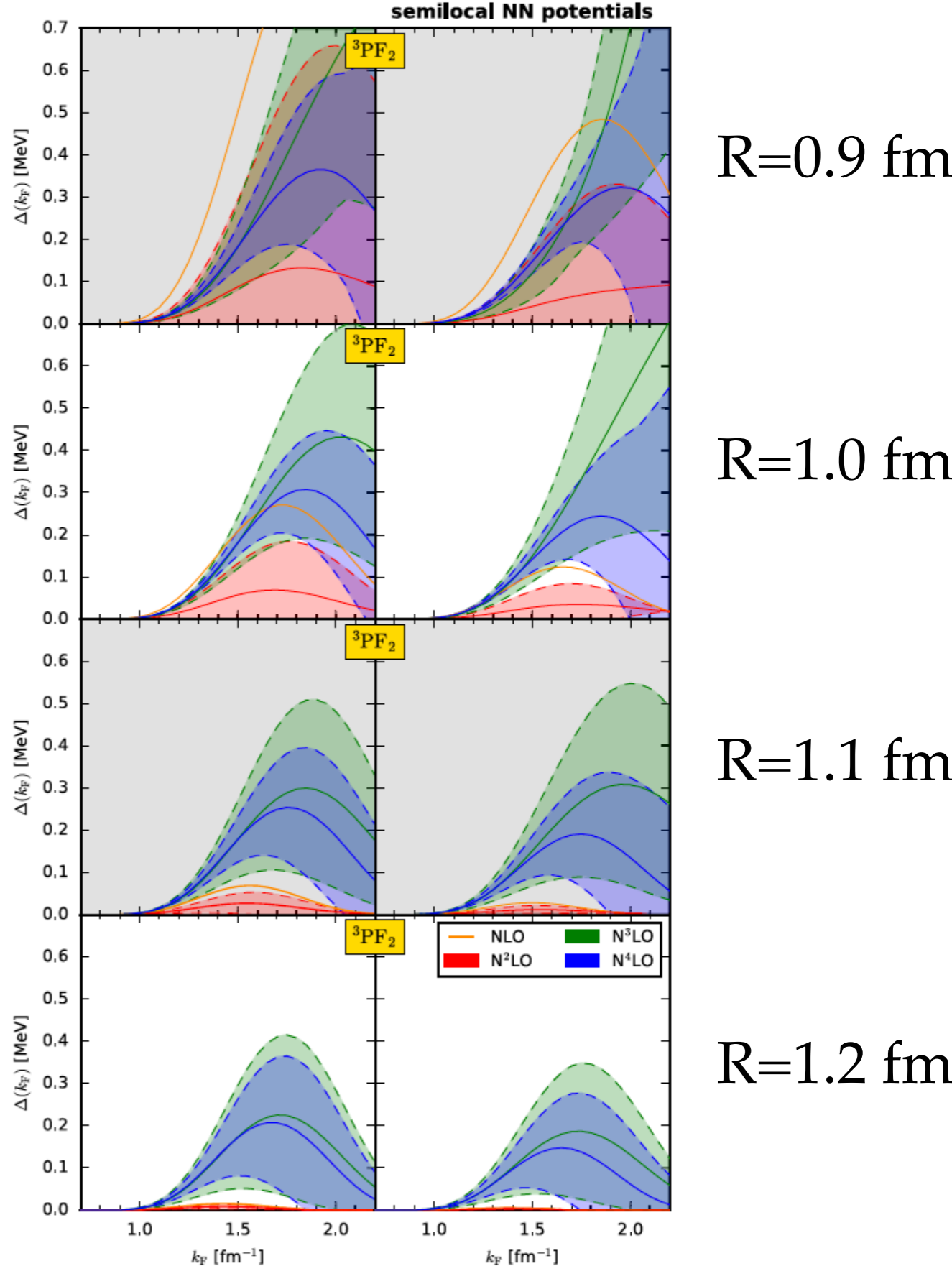
Uncertainties at the BCS + HF level

“Hard”



Srinivias & Ramanan, *PRC* **94** 064303 (2016)

“Soft”



1. Neutron star motivation
2. Infinite matter BCS
- 3. Beyond-BCS with SCGF methods**

(B) $iF(1, 2) = \langle T\{\psi(1)\psi(2)\} \rangle = \text{---} = \text{---} \circlearrowleft \Sigma \text{---} + \text{---} \circlearrowright \Delta^* \text{---}$

(B') $iF^\dagger(1, 2) = \langle T\{\psi^\dagger(1)\psi^\dagger(2)\} \rangle = \text{---} = \text{---} \circlearrowleft \Sigma \text{---} + \text{---} \circlearrowright \Delta \text{---}$

(C) $iG(1, 2) = \langle T\{\psi(1)\psi^\dagger(2)\} \rangle = \text{---} = \text{---} + \text{---} \circlearrowleft \Sigma \text{---} + \text{---} \circlearrowright \Delta \text{---}$

(D) $\text{---} \circlearrowleft \Sigma \text{---} = \text{---} \text{K} \text{---}$

(E) $\text{---} \circlearrowright \Delta \text{---} = \text{---} \text{K} \text{---}$

Normal
state

Superfluid
 $\Delta(k_F)$

?

- **BCS** is lowest order in Gorkov Green's function expansion
- T-matrix can be extended to paired systems

Gorkov gap equation

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\xi_{k'}} \Delta_{k'}^{L'} + \frac{1}{2\xi_k} = \int_{\omega} \int_{\omega'} \frac{1 - f(\omega) - f(\omega')}{\omega + \omega'} A(k, \omega) A_s(k, \omega')$$

BCS+SRC gap equation

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\bar{\chi}_{k'}^2 + |\Delta_{k'}^{L'}|^2}} \Delta_{k'}^{L'} + \frac{1}{2\bar{\chi}_k} = \int_{\omega} \int_{\omega'} \frac{1 - f(\omega) - f(\omega')}{\omega + \omega'} A(k, \omega) A_s(k, \omega')$$

BCS+Z-factor equation

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{Z_k \langle k | V_{nn}^{LL'} | k' \rangle Z_{k'}}{2\sqrt{\chi_{k'}^2 + |\Delta_{k'}^{L'}|^2}} \Delta_{k'}^{L'} \quad \begin{aligned} &+ \chi_k = \varepsilon_k - \mu \\ &+ \varepsilon_k = \frac{k^2}{2m} + U(k) \end{aligned}$$

BCS gap equation

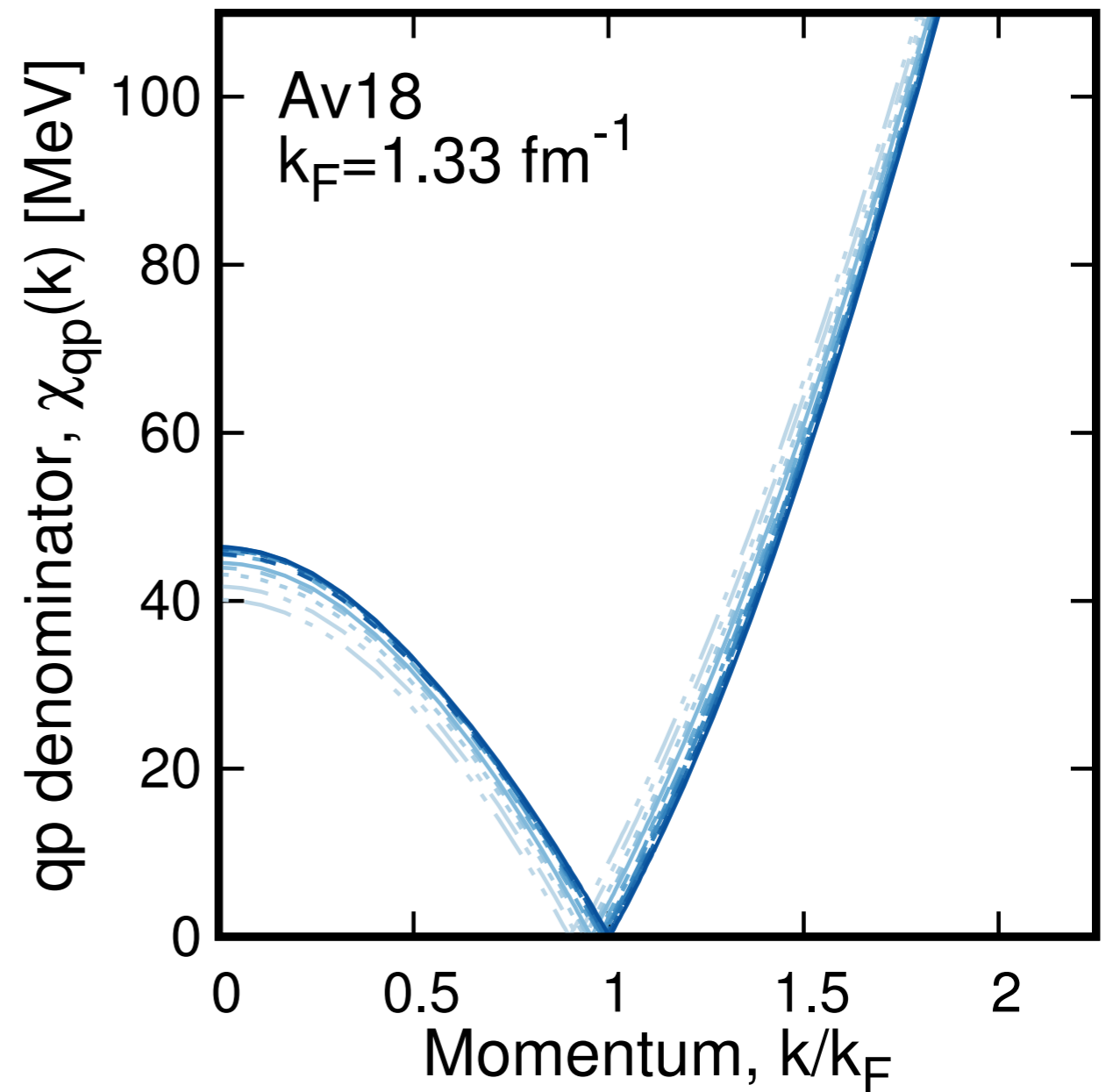
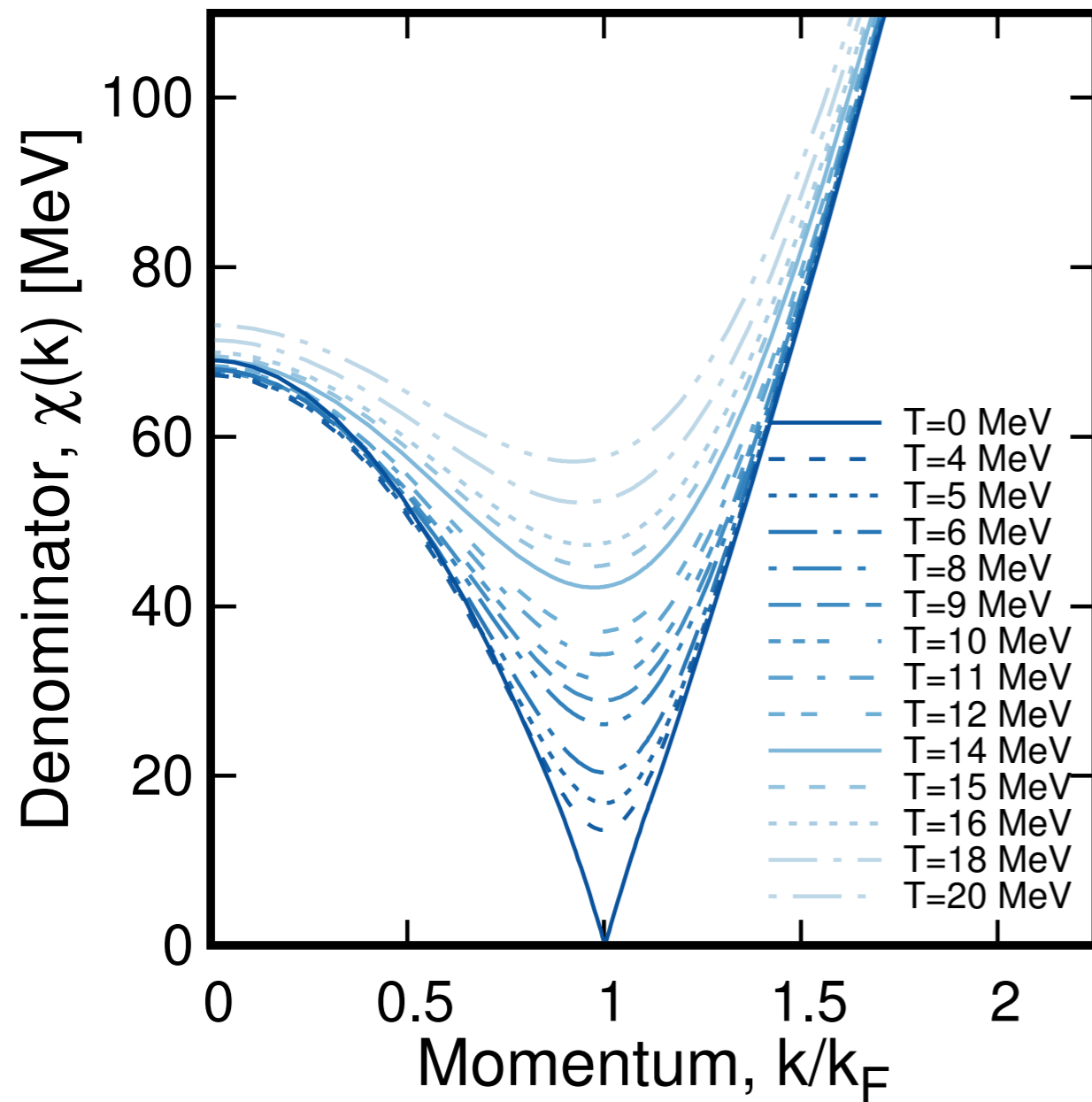
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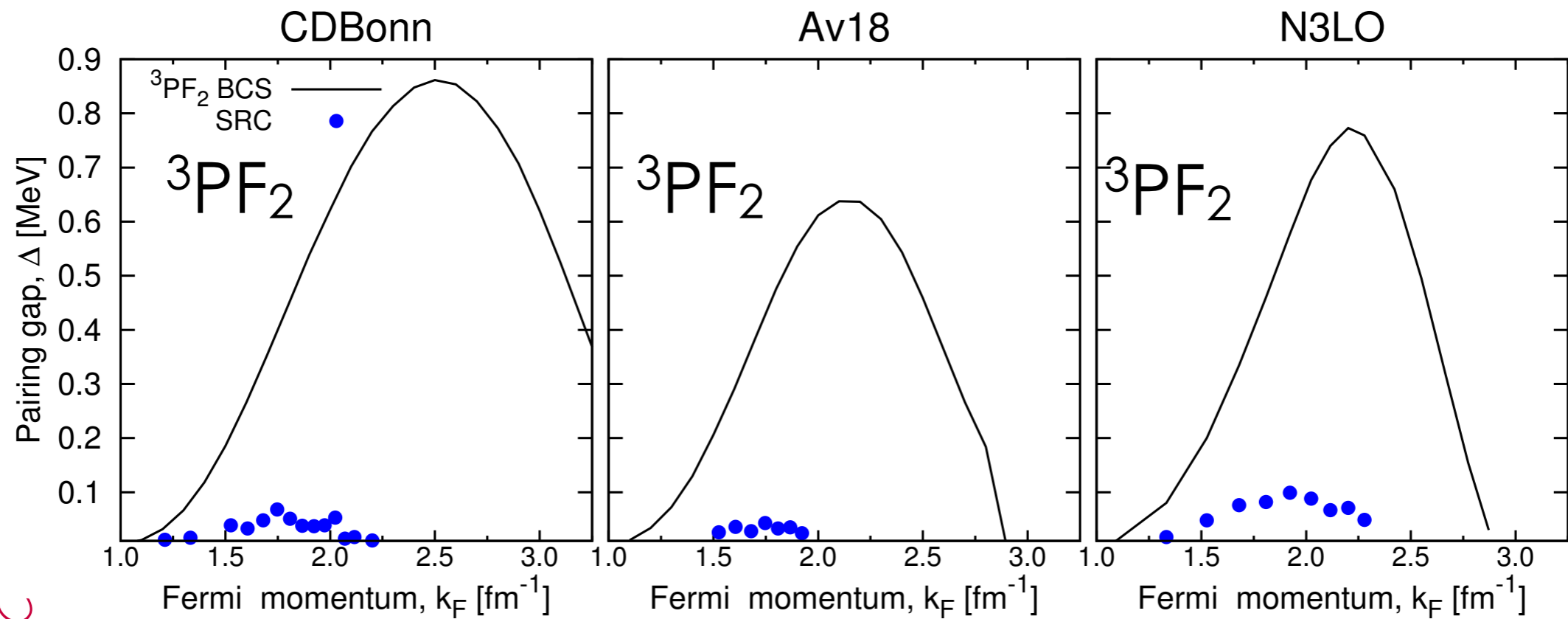
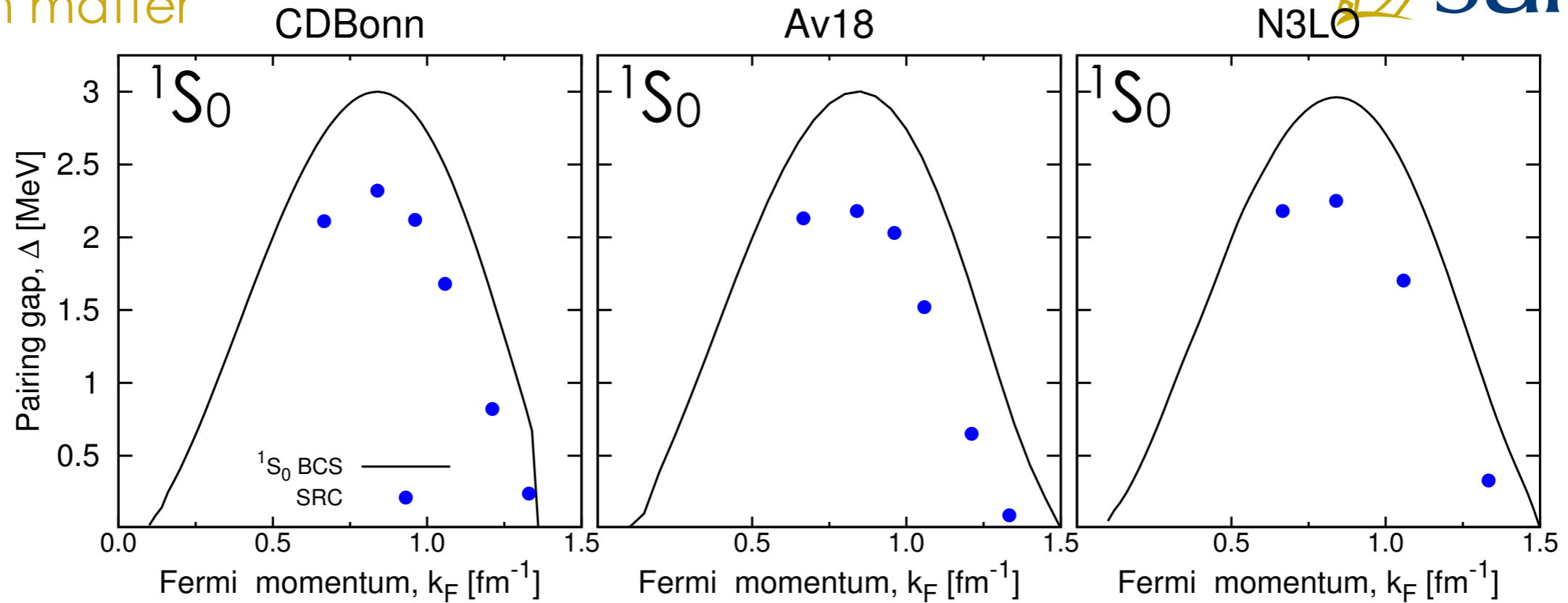
Full off-shell

$$\frac{1}{2\bar{\chi}_k} = \int_{\omega} \int_{\omega'} \frac{1 - f(\omega) - f(\omega')}{\omega + \omega'} A(k, \omega) A(k, \omega')$$

Quasi-particle \leftrightarrow BCS

$$\frac{1}{2\bar{\chi}_k} = \frac{1}{2|\varepsilon_k - \mu|}$$



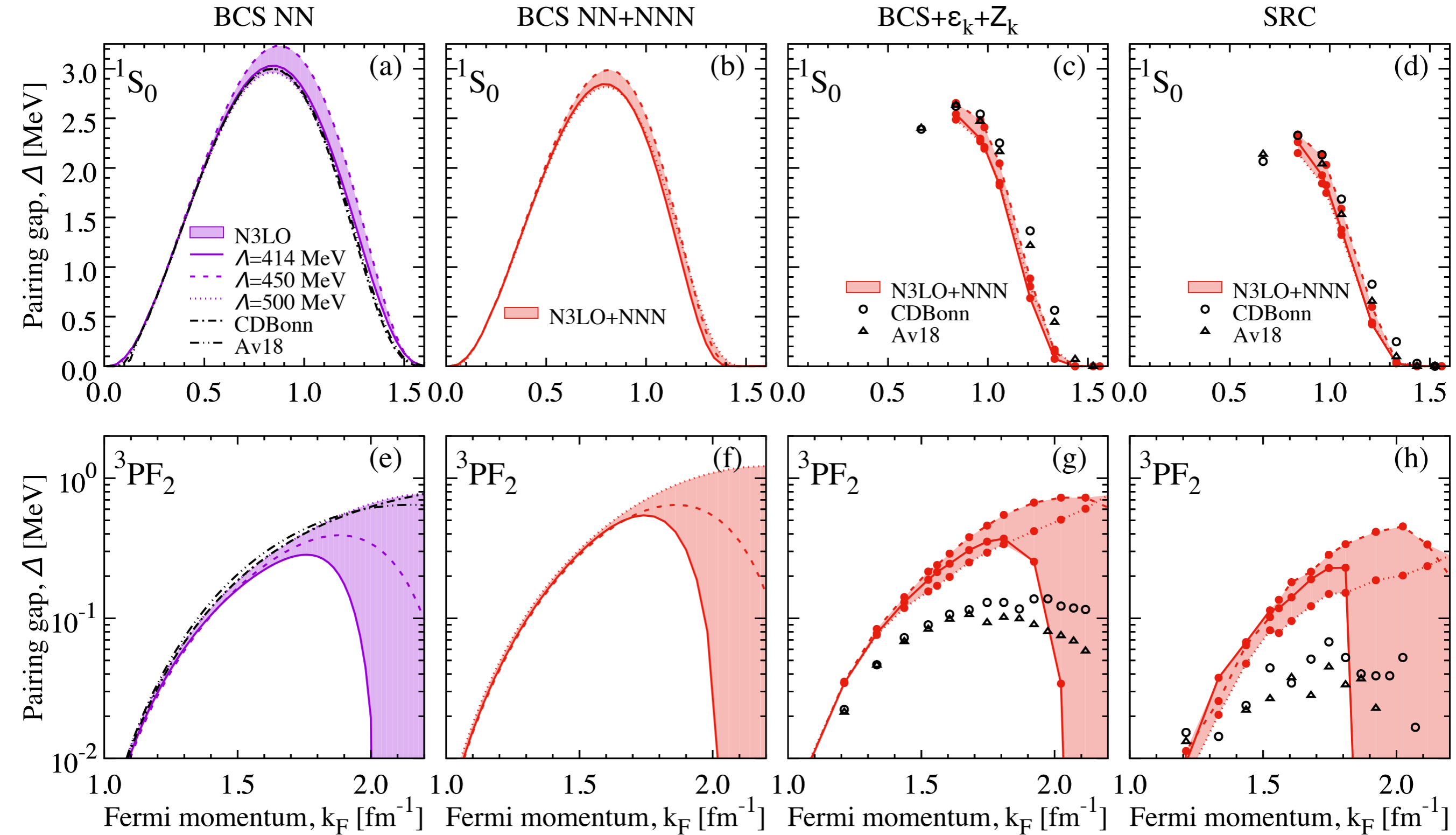


BCS+SRC

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\chi_{k'}^2 + |\Delta_{k'}^{L'}|^2}} \Delta_{k'}^{L'} + \frac{1}{2\chi_k} = \int_{\omega} \int_{\omega'} \frac{1 - f(\omega) - f(\omega')}{\omega + \omega'} A(k, \omega) A_{\times}(k, \omega')$$

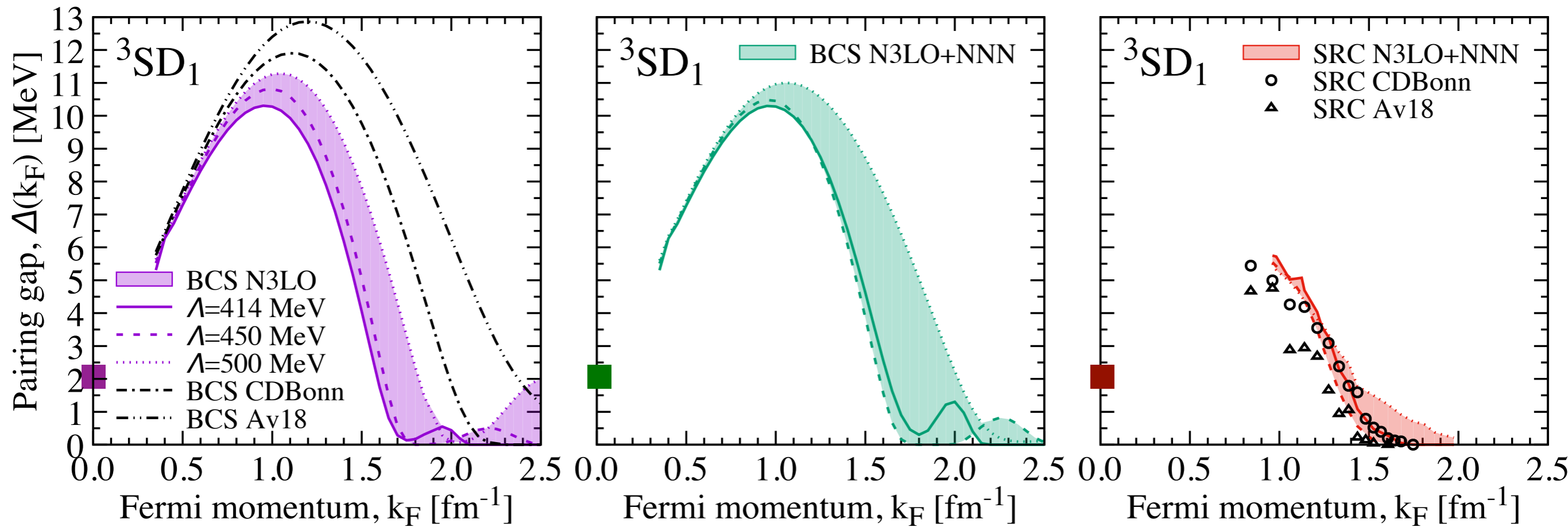
Beyond BCS 101: SRC

Neutron matter



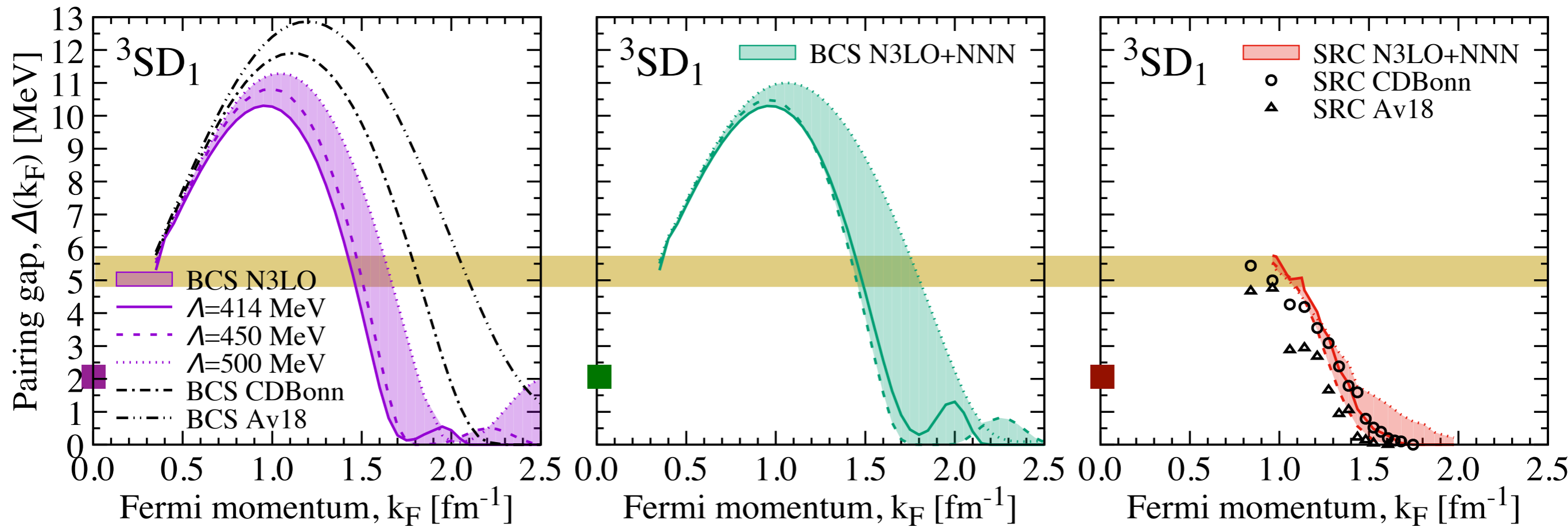
BCS+SRC

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\chi_{k'}^2 + |\Delta_{k'}^{L'}|^2}} \Delta_{k'}^{L'} + \frac{1}{2\chi_k} = \int_{\omega} \int_{\omega'} \frac{1 - f(\omega) - f(\omega')}{\omega + \omega'} A(k, \omega) A_{\times}(k, \omega')$$



Muether & Dickhoff, *PRC* **72** 054313 (2005)
 Rios, Polls & Dickhoff, arXiv:1707.04140
 Maurizio, Holt & Finelli, *PRC* **90**, 044003 (2014)
 U. Lombardo's talk

- **Massive** gaps 3SD_1 channel but...
- **No evidence** of strong np nuclear pairing
- 3NF do **not** alter picture **significantly**
- Short-range correlations **deplete** gap



Muether & Dickhoff, *PRC* **72** 054313 (2005)
 Rios, Polls & Dickhoff, arXiv:1707.04140
 Maurizio, Holt & Finelli, *PRC* **90**, 044003 (2014)
 U. Lombardo's talk

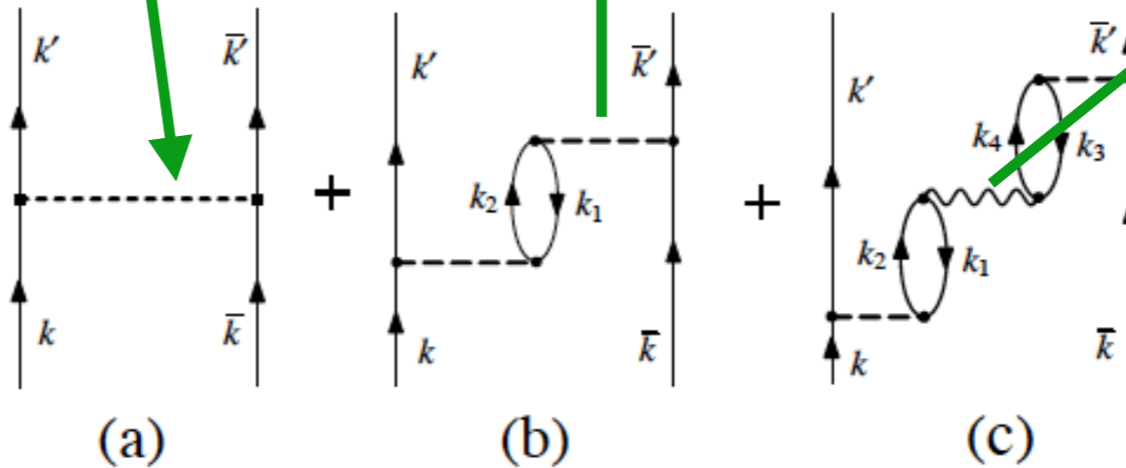
- **Massive** gaps 3SD_1 channel but...
- **No evidence** of strong np nuclear pairing
- 3NF do **not** alter picture **significantly**
- Short-range correlations **deplete** gap

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\chi_{k'}^2 + |\Delta_{k'}^{L'}|^2}} \Delta_{k'}^{L'}$$

ph recoupled
G-matrix

Effective Landau
parameters

$\mathcal{V}_{\text{pair}} =$

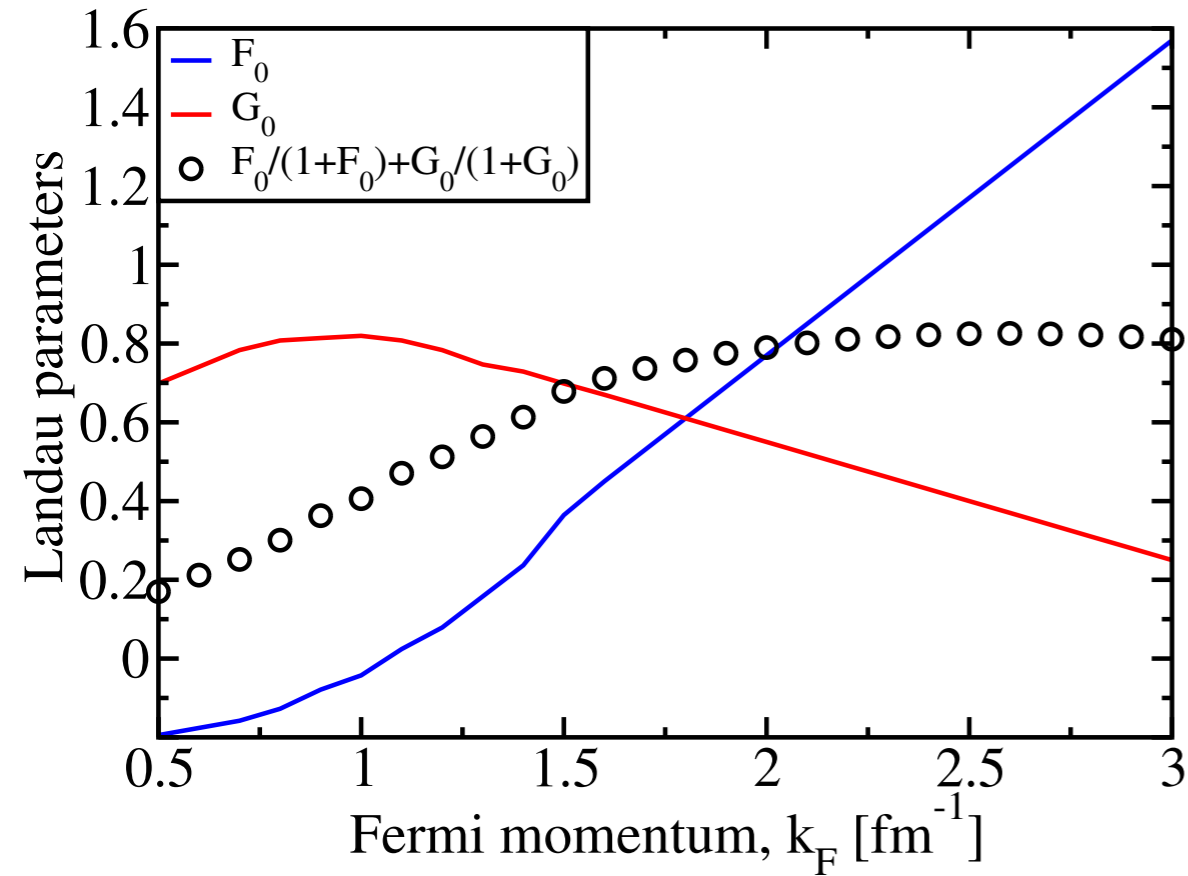


$$\langle 1\bar{1} | \mathcal{V} | 1\bar{1} \rangle = \frac{1}{4} \sum_{2,2'} \sum_{S,T} (-)^S (2S+1) \langle 12 | G_{ST}^{\text{ph}} | 1'2' \rangle_A \langle 2'\bar{1} | G_{ST}^{\text{ph}} | 2\bar{1}' \rangle_A \Lambda(22')$$

$$\Lambda_{ST}(q) = \frac{\Lambda_{ST}^0(q)}{1 - \Lambda_{ST}^0(q) \times F_{ST}}$$

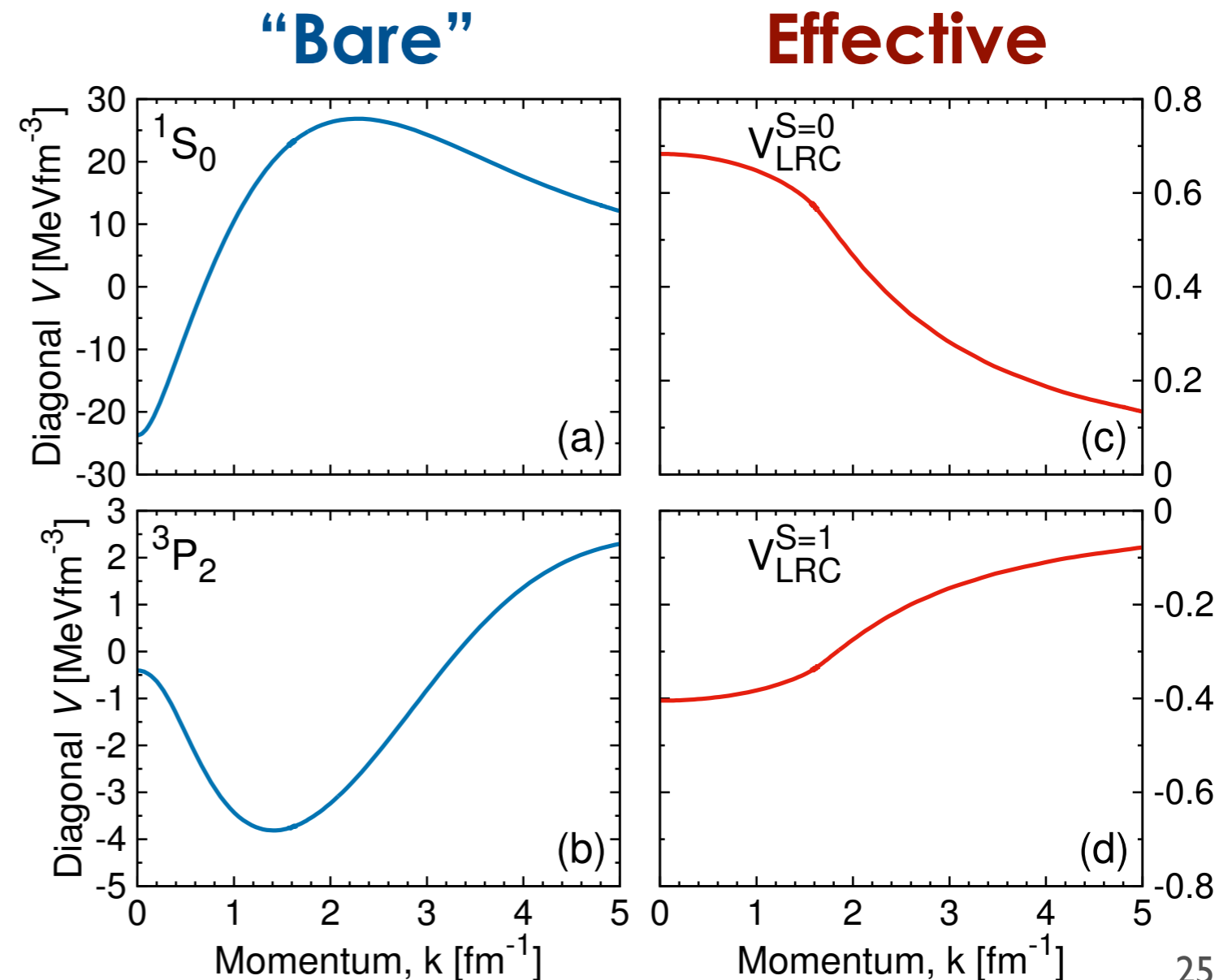
- Bare NN potential only is not the only possible interaction
- Diagram (a): nuclear interaction
- Diagram (b): in-medium interaction, density and spin fluctuations
- Diagram (c): included by Landau parameters

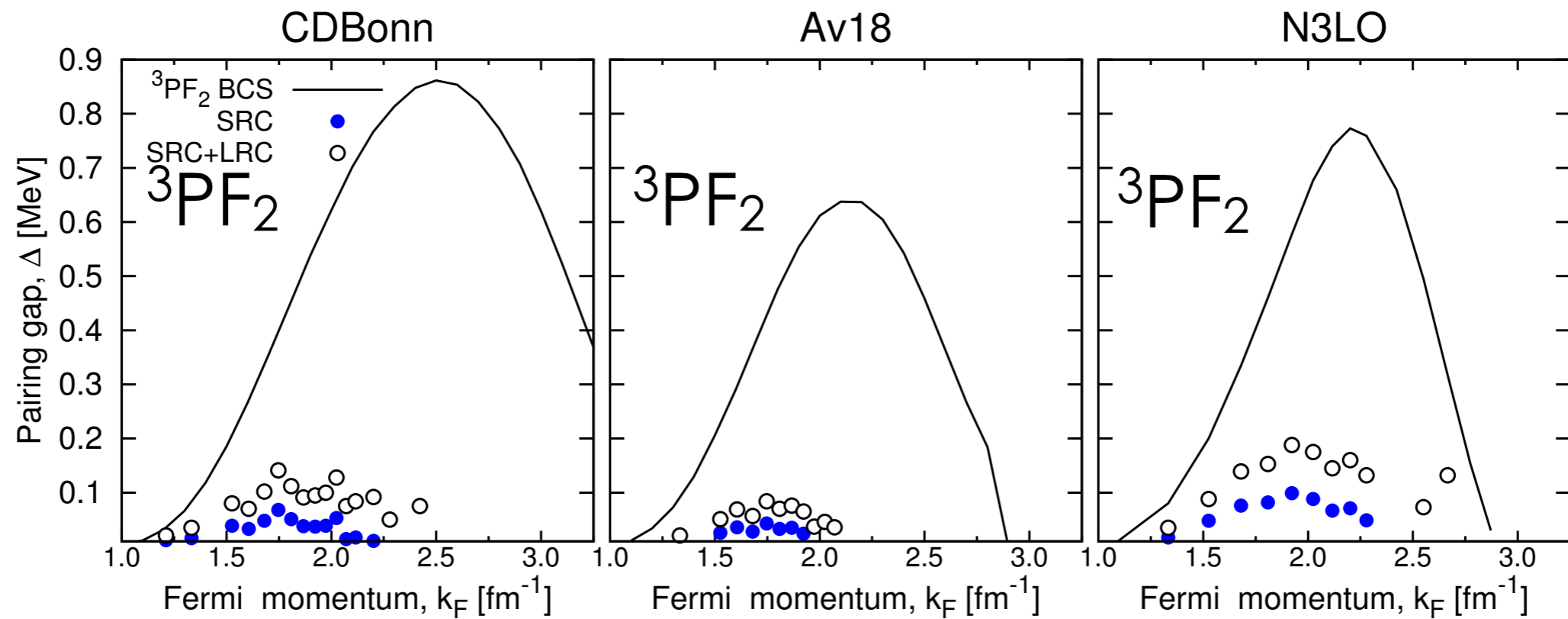
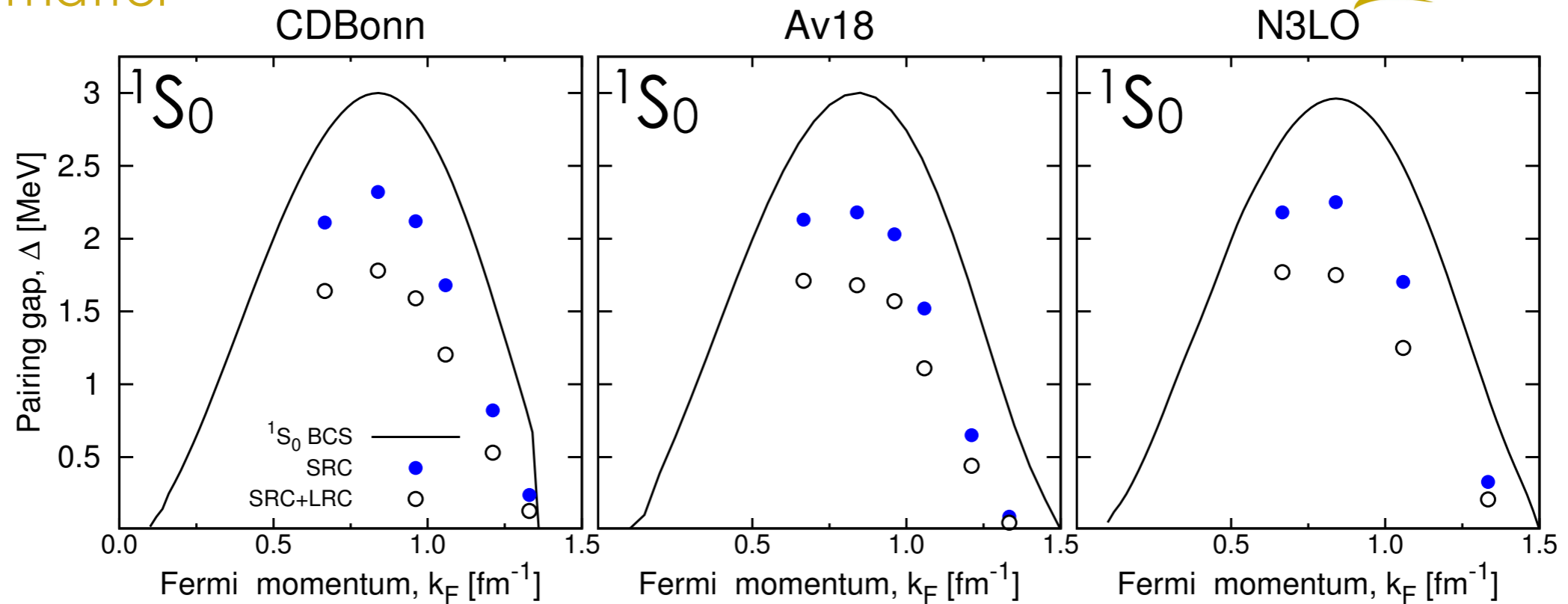
Landau parameters



- **Effective** Landau parameters
- **Not consistent** (yet) with NN force
- LRC in **nuclei** will be **different**

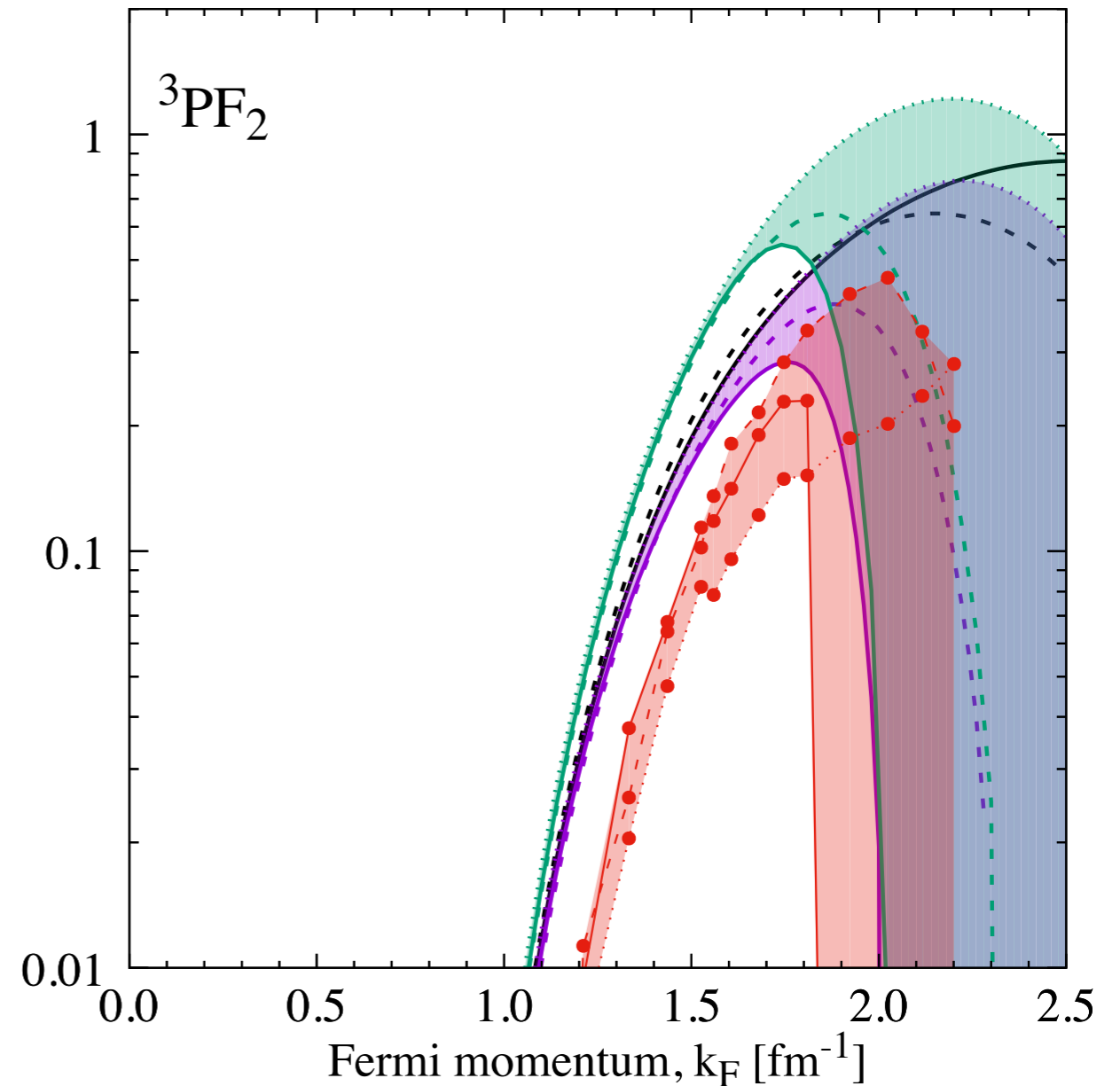
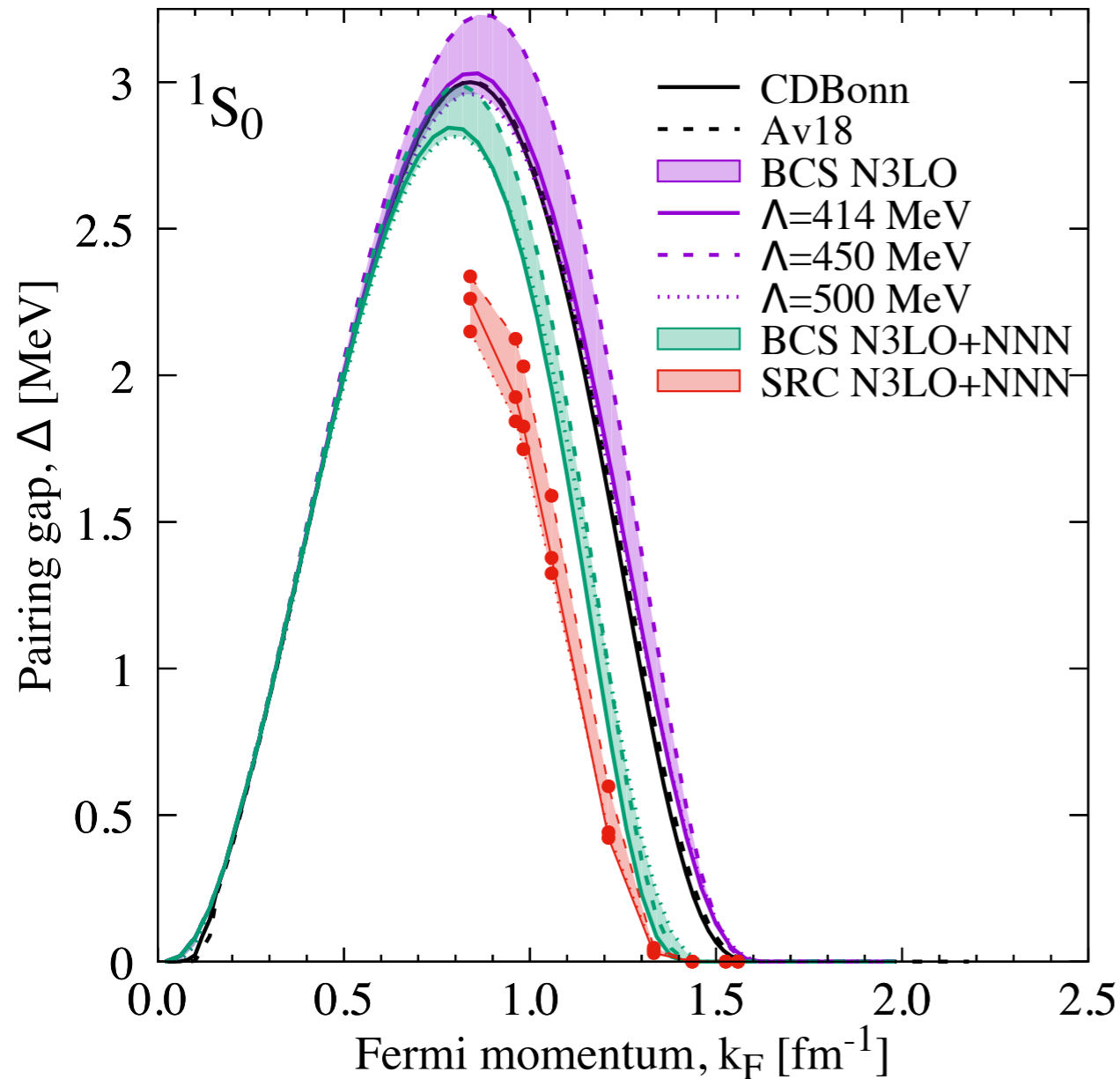
- **Small** correction
- Repulsive in **singlet**
- Attractive in **triple**





- LRC 1S_0 ($^3P_{F_2}$) produces (anti-)screening

3BF effect: uncertainty estimate



- Singlet gap: 3NF **reduce** closure
- Triplet gap: 3NF **increase** gap
- **LRC** effect to be explored

Effective one-body force \Rightarrow spectrum

$$\bullet \text{---} \times = \bullet \text{---} \bullet \text{---} \bullet + \frac{1}{2} \bullet \text{---} \bullet \text{---} \bullet \text{---} \bullet$$

Effective two-body force \Rightarrow NN forces

$$\bullet \text{---} \bullet \text{---} \bullet = \bullet \text{---} \bullet + \bullet \text{---} \bullet \text{---} \bullet$$

RESEARCH ARTICLE

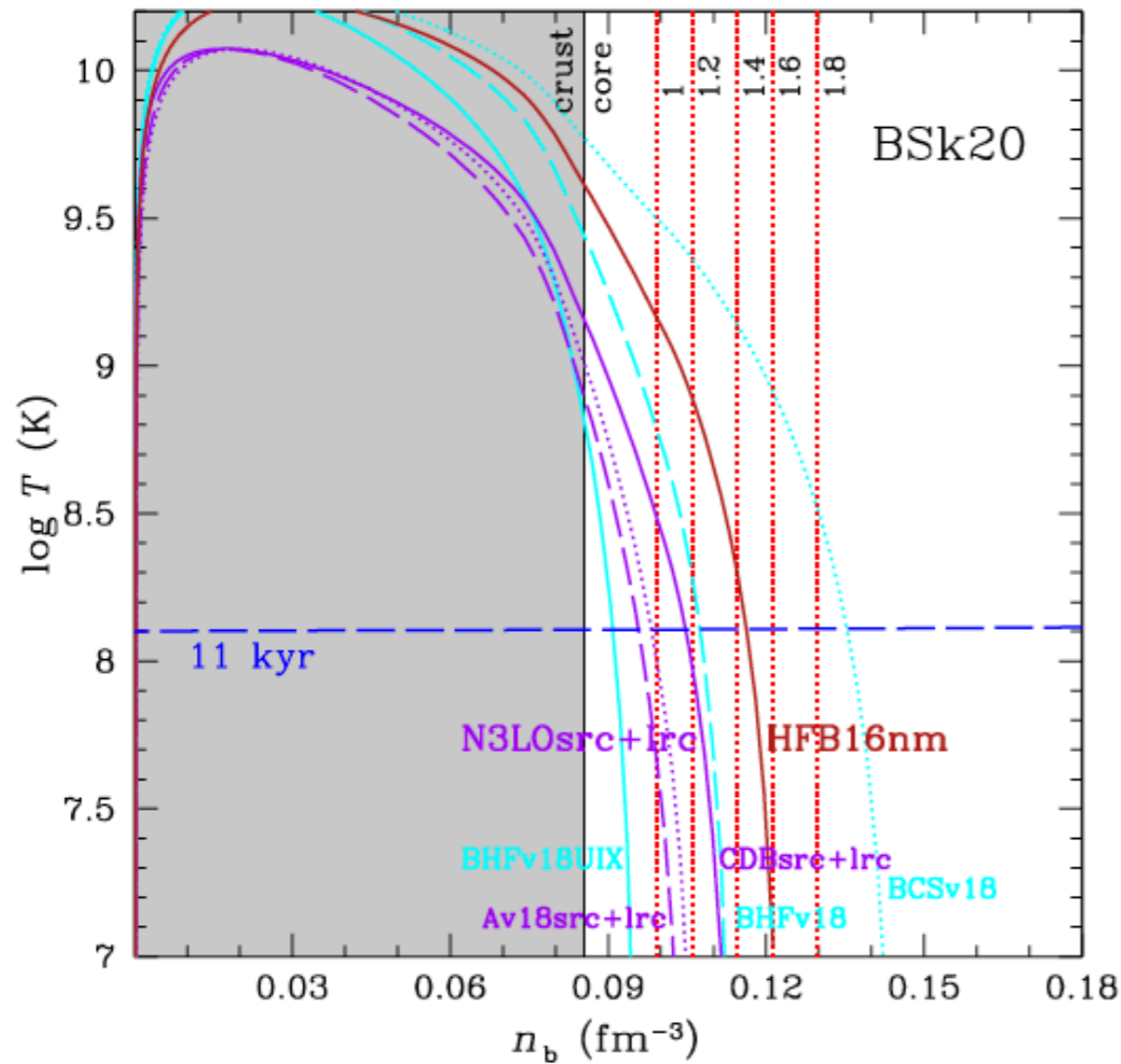


Fig. 2. Temperature dependence of neutron superfluid models as a function of baryon number density for the BSk20 nuclear equation of state. Thick curved lines are the superfluid critical temperature for the (labeled) models from (23). The vertical solid line indicates the separation between the crust (shaded region) and the core. Vertical dotted lines denote the density at which the superfluid moment of inertia (using the SFB superfluid model) is 1.6% of the total stellar moment of inertia for neutron stars of different mass (labeled in units of solar mass). The (nearly horizontal) dashed line is the temperature of a $1.4 M_{\text{Sun}}$ neutron star at an age of 11,000 years.

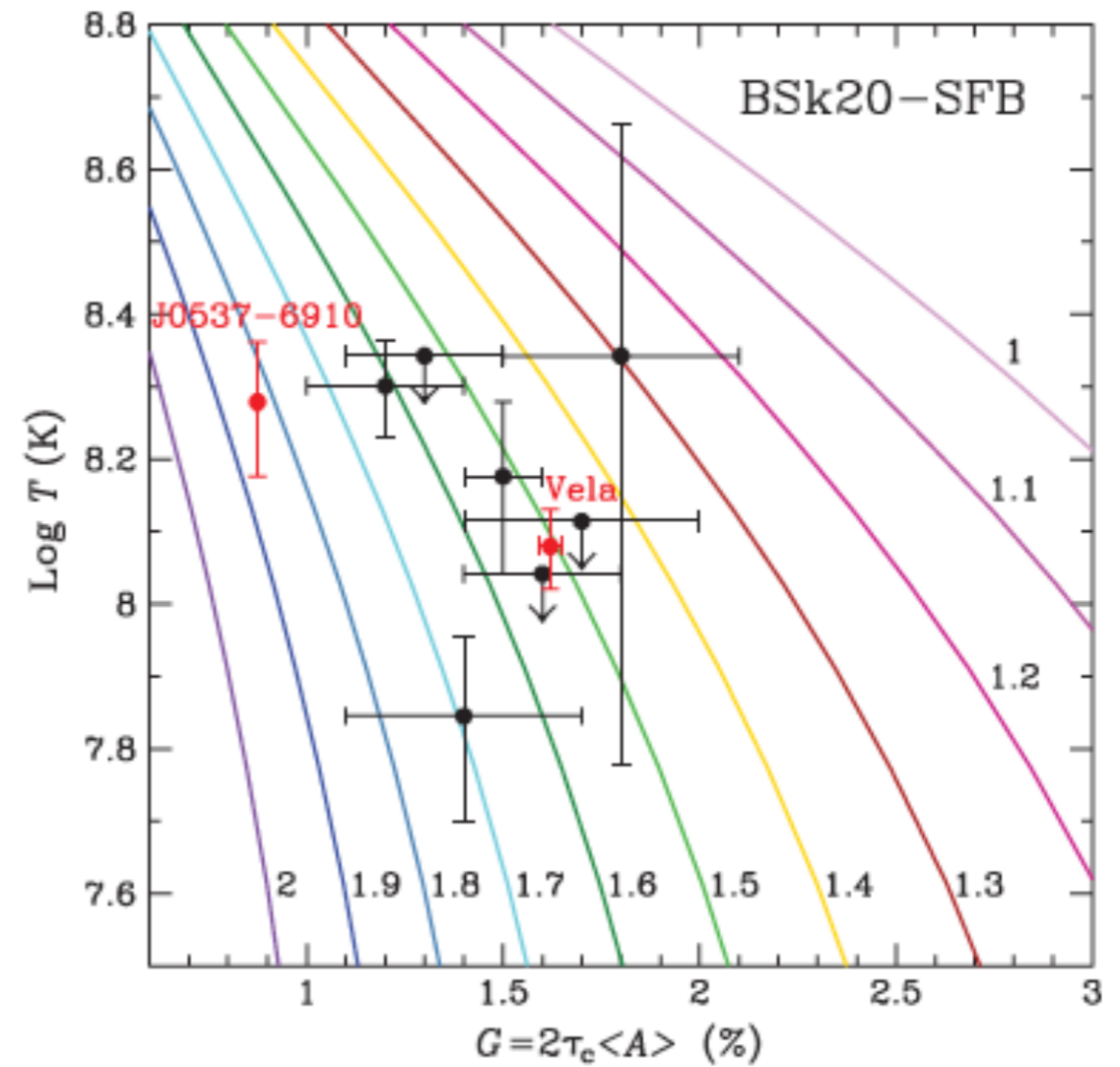


Fig. 3. Neutron star mass from pulsar observables G and interior temperature T . Data points are for pulsars with measured G from glitches and T from an age or surface temperature observation (see Table 1). Lines (labeled by neutron star mass, in units of solar mass) are the theoretical prediction for G and T using the BSk20 nuclear equation of state and SFB neutron superfluid models.

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- Ab initio nuclear **theory** to treat beyond-BCS **correlations**
- **Different** NN forces provide robust predictions
- Challenges ahead
 - Full self-consistent **Gorkov**
 - **Consistent** treatment of **LRC**
 - **Pairing** in isospin **asymmetric** matter

Postdoc position opening imminently